

Sunyaev-Zeldovich Effect Survey Results from Bolocam

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KICP Friday Seminar

April 4, 2008

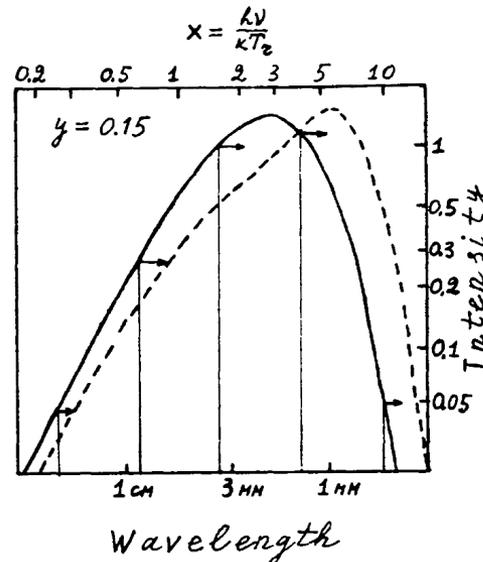
Overview

- Review of the SZ effect
- Applications of the SZ effect
- Bolocam instrument description
- Sky noise removal and analysis techniques
- Constraints on SZ anisotropy
- Upcoming work

The Sunyaev-Zeldovich Effect in Galaxy Clusters

- Thermal SZE is the Compton up-scattering of CMB photons by hot electrons in the intracluster plasma
- $\Delta T_{CMB}/T_{CMB}$ depends only on cluster y (line-of-sight integral of $n_e T_e$). Both ΔT_{CMB} and T_{CMB} are redshifted as photons propagate from clusters, so ratio is independent of distance.

- Thermal SZE causes nonthermal change in spectrum



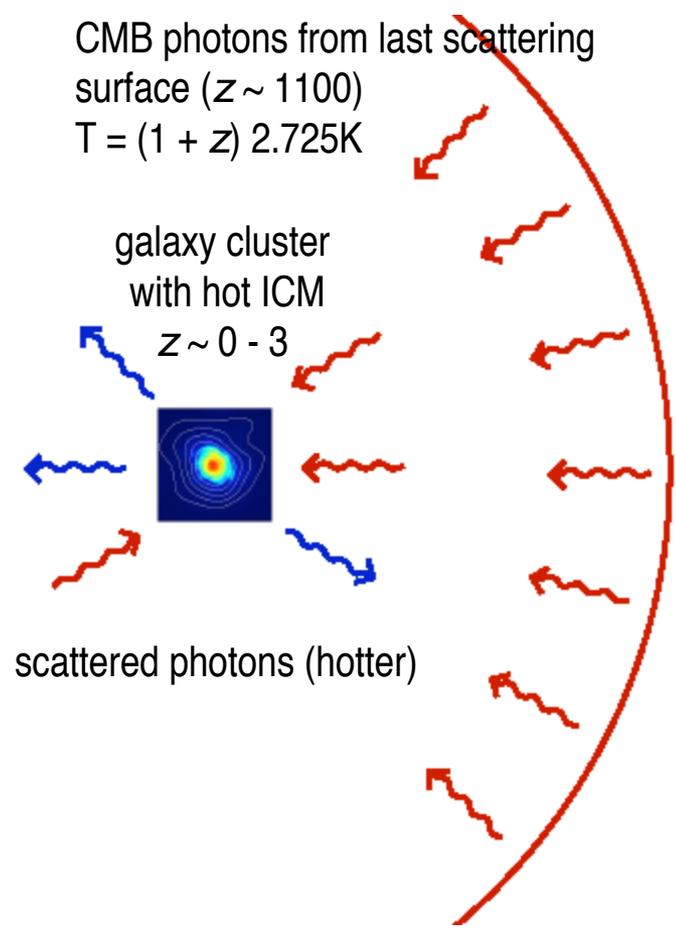
observer
 $z = 0$



CMB photons from last scattering surface ($z \sim 1100$)
 $T = (1 + z) 2.725K$

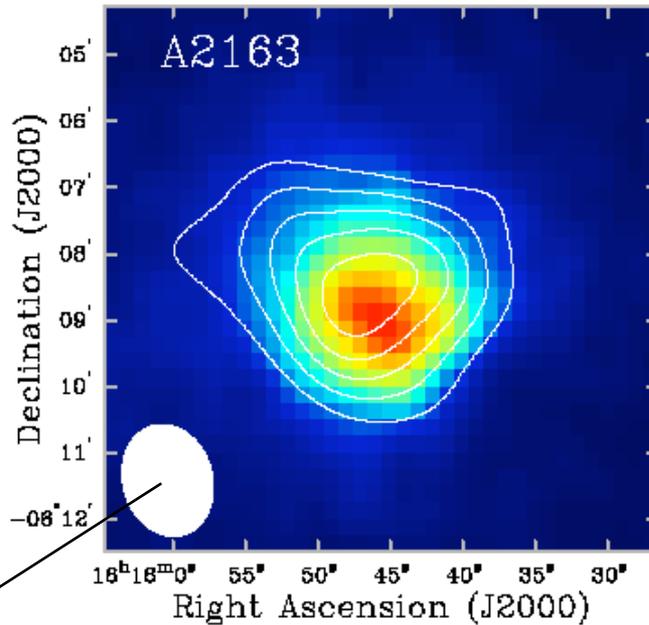
galaxy cluster with hot ICM
 $z \sim 0 - 3$

scattered photons (hotter)

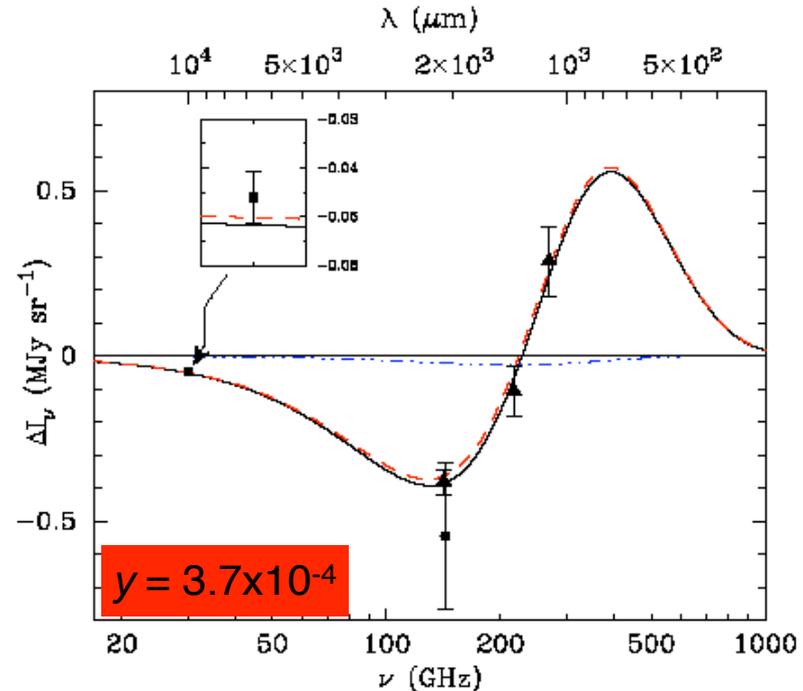


The Sunyaev-Zeldovich Effect in Galaxy Clusters

BIMA SZ + ROSAT X-ray



SuZIE, DIABOLO, and BIMA spectral points



- Beautiful images of SZ from Chicago group using OVRO/BIMA interferometers at 30 GHz
- Spectrum confirmed by measurements from RJ tail through null
- To date, only seen in pointed observations of massive clusters

Applications of the SZ Effect

- Cluster astrophysics
 - measures pressure
 - scaling relations
- Cosmology
 - Hubble constant (geometric effect, with X-ray)
 - Baryon fraction (now measured better by CMB)
 - Evolution of cluster abundance as a probe of dark energy

Studying Clusters with the SZ Effect

Clusters are complicated objects!

SZ measures pressure, in contrast to other observables

gas density

galaxies

$SZ \sim n T$
(gas pressure)

X-ray
 $\sim n^2 \sqrt{T}$



10 Mpc/h comoving

6

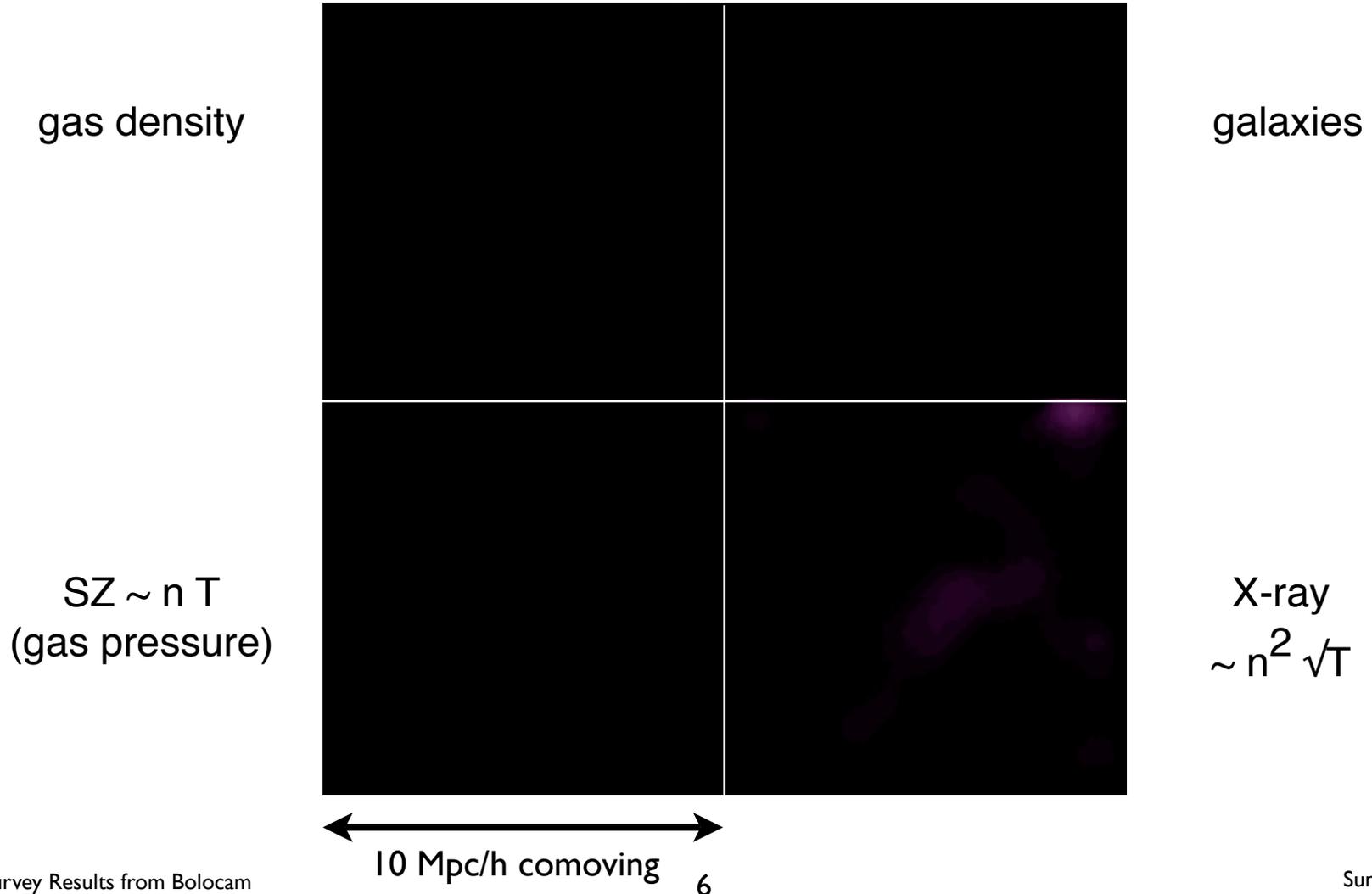
<http://astron.berkeley.edu/~mwhite/clusterform.html>

Martin White, Berkeley

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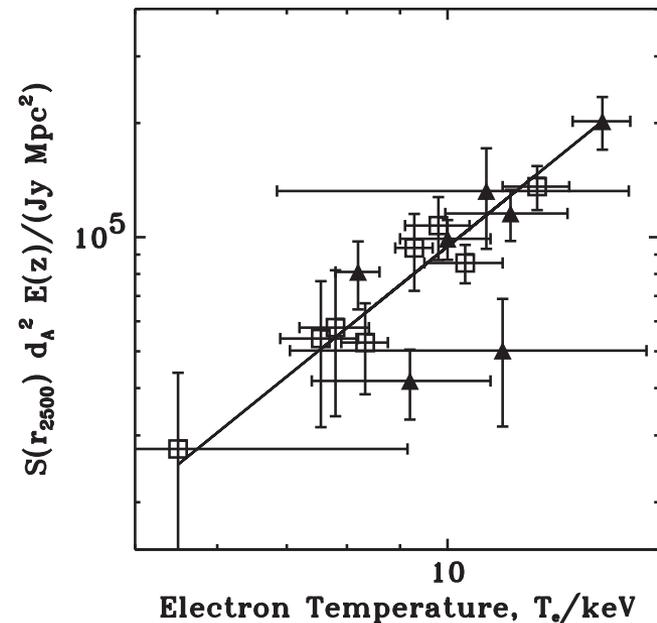
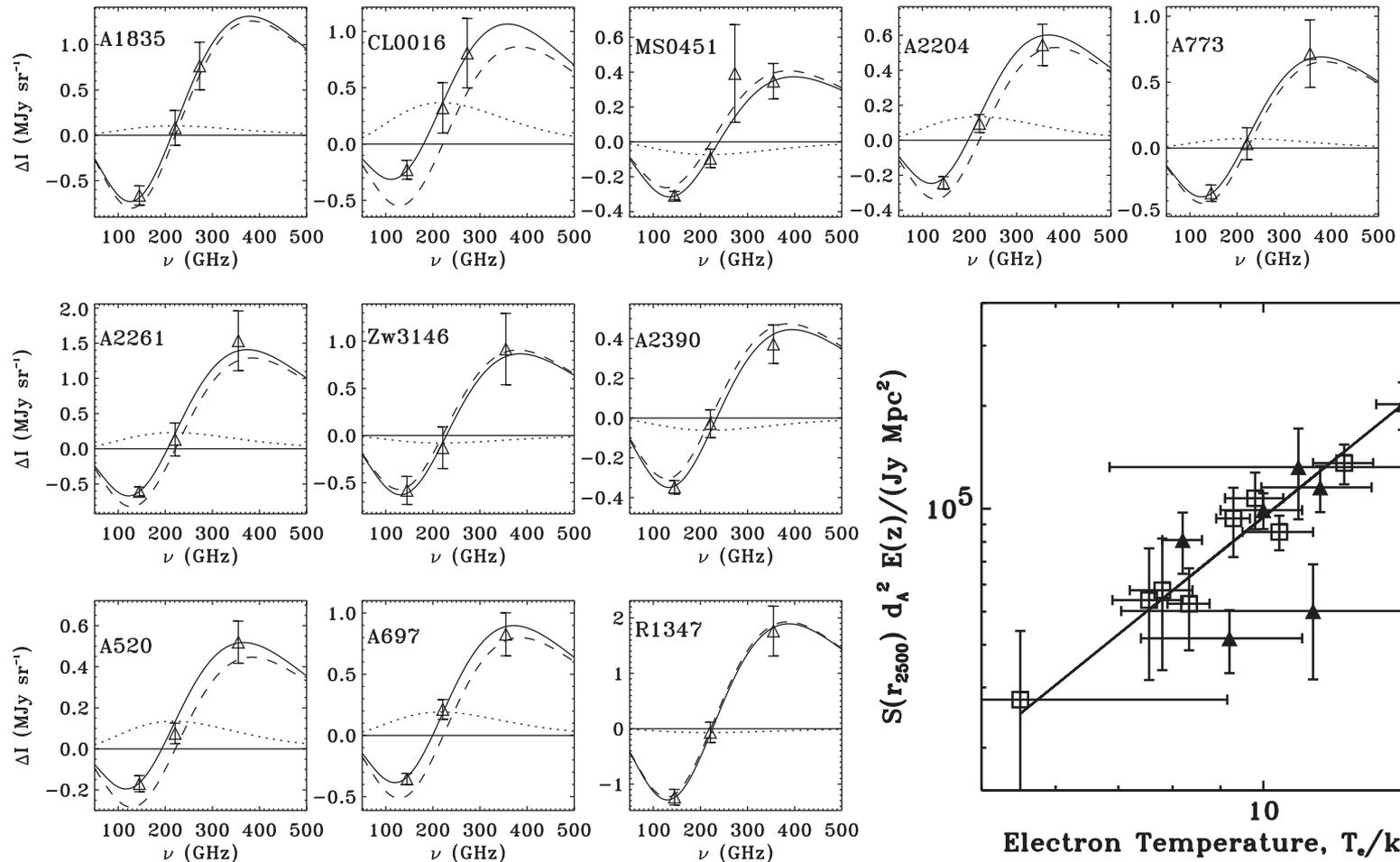
Martin White, Berkeley

Studying Clusters with the SZ Effect: Scaling Relations

- SuZIE (S. Church, Stanford)

Benson *et al*, *Apj* 617:829 (2004)

- published 11 clusters at 150/220/275(350) GHz, observed SZ flux- T_x scaling relation, but not an imaging experiment

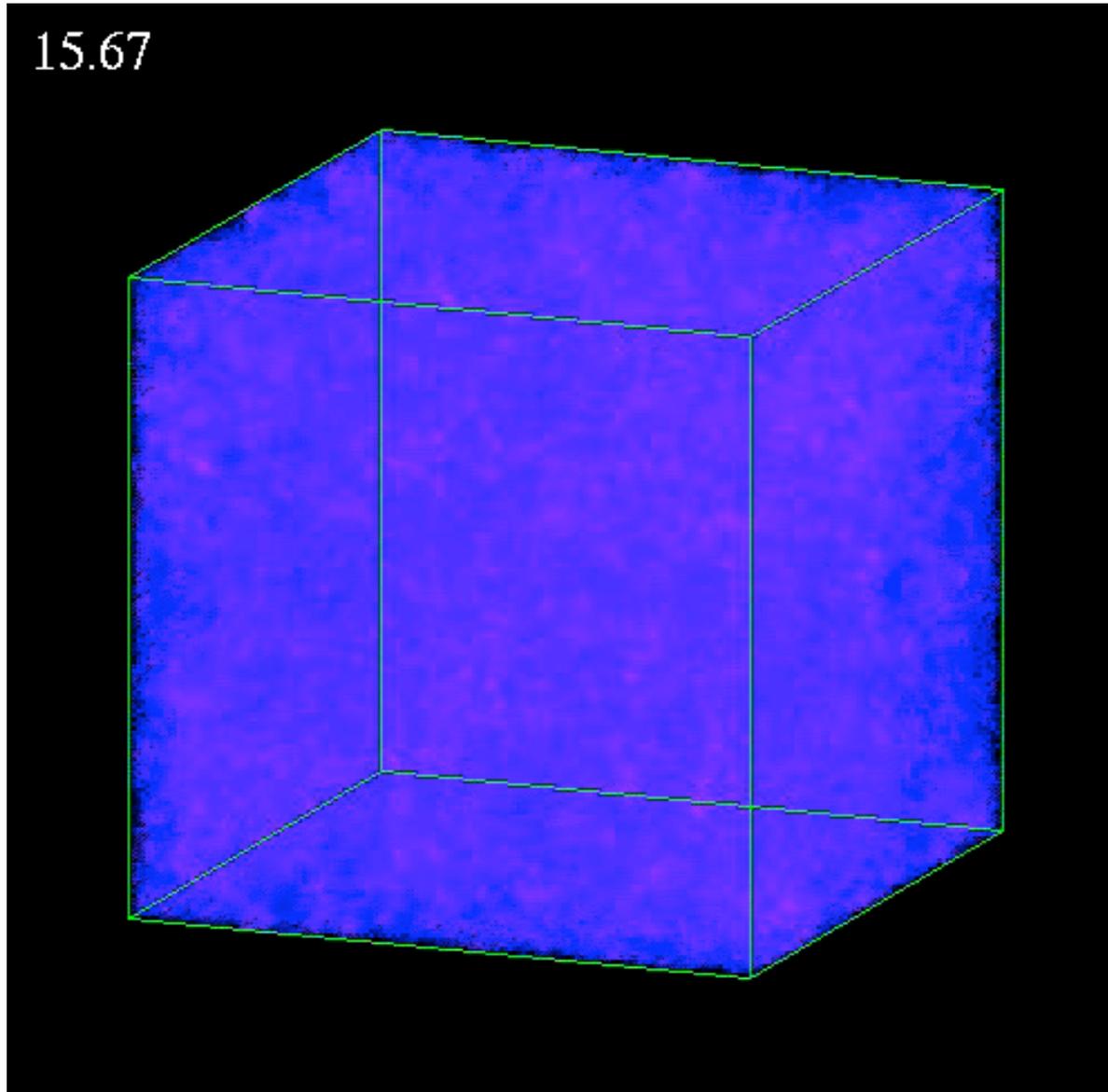


Probing Dark Energy via the Growth of Structure

Virgo Consortium

<http://www.icc.dur.ac.uk/Outreach/Movies.html>

Probing Dark Energy via the Growth of Structure

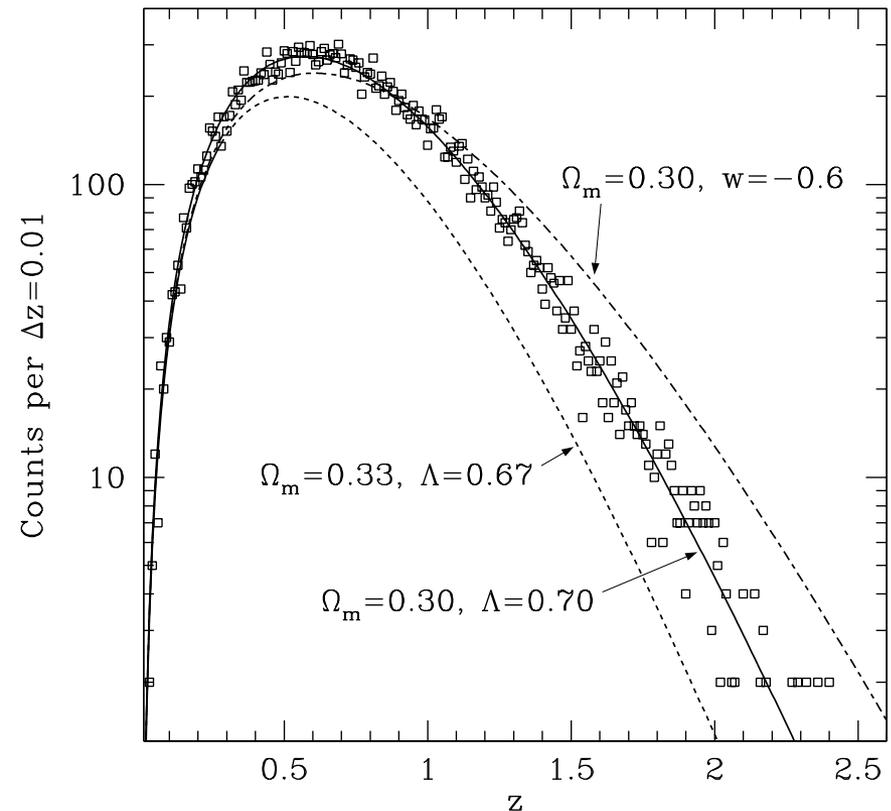


Virgo Consortium
<http://www.icc.dur.ac.uk/Outreach/Movies.html>

Using Cluster Abundance for Cosmology

- Very sensitive to normalization of power spectrum, and thus to growth function, because clusters are statistically rare excursions
- Clusters form recently ($z < 2$) and so abundance influenced by recent dark-energy domination
- Has historically been a robust predictor of low matter density

Number of clusters per redshift bin above $3.5 \times 10^{14} M_{\text{Sun}}$ in 4000 deg^2



G. Holder

“Unbiased” Cluster Detection via the SZE

- “Unbiased” = mass-limited
- Effect is intrinsically redshift-independent: $\Delta T/T$ depends only on cluster properties, ΔT and T experience same redshift
- Standard argument: Integrated signal provides largely z-independent mass limit (Barbosa *et al*, Holder *et al*, etc.)

$$S_{tot} = \frac{2k_B^2 \nu^2 g(x) \sigma_T T_{CMB}}{m_e c^4 d_A(z)^2} \langle T_e \rangle_n \frac{\overset{\text{cluster mass}}{M_{200} f_{ICM}}}{\mu_e m_p}$$

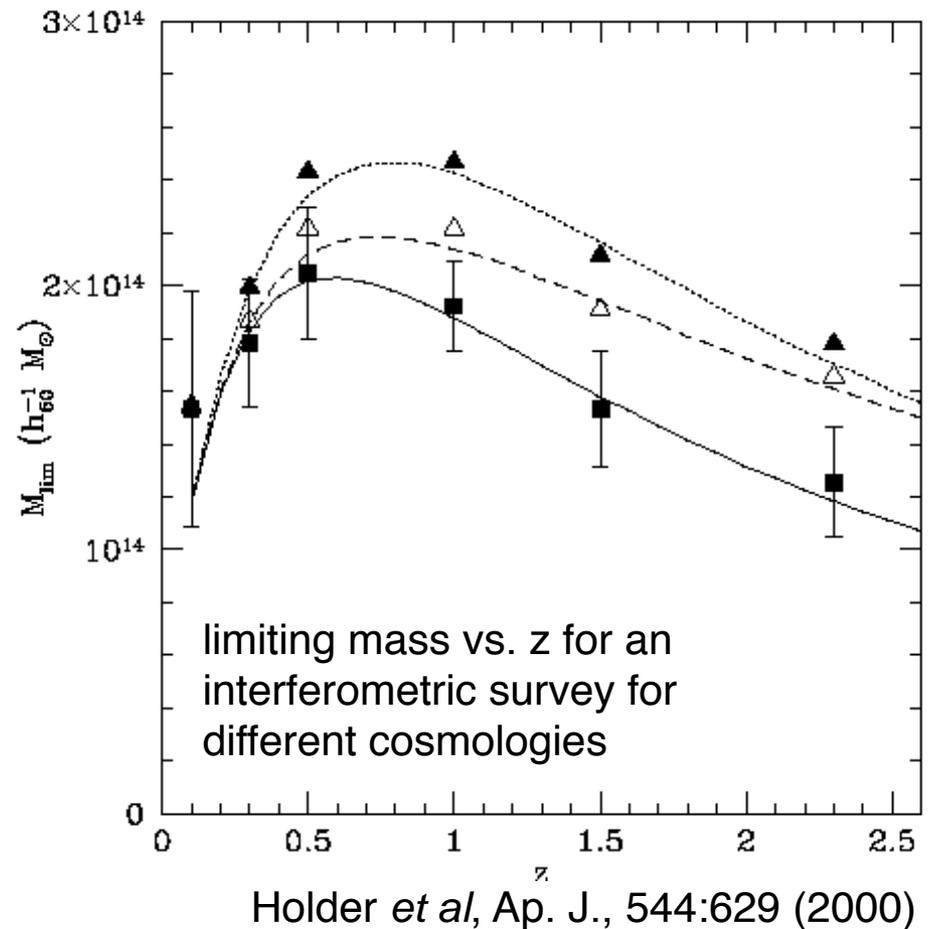
- Integrate $n_e T_e$ over cluster face

weak z-dependence of ang. diam. distance

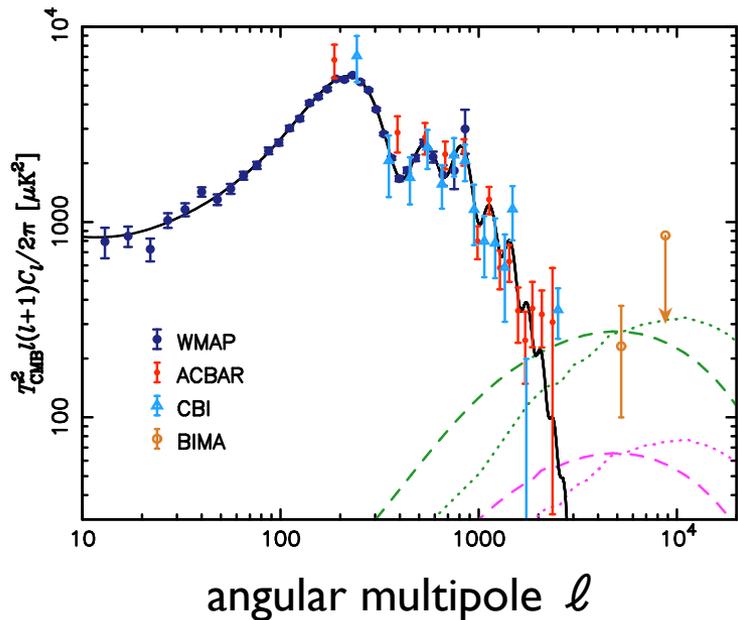
- d_A^2 factor tends to reduce flux as z increases ($1/r^2$ law)
- But for a given mass, a cluster at high redshift has smaller R and hence higher T
- These two effects approximately cancel

“Unbiased” Cluster Detection via the SZE

- Holder, Mohr, et al (2000) modeled the mass limit of an interferometric SZE survey using simulations of cluster growth
- Simulations bear out expectation of weak z-dependence of mass limit
- v. different selection function from optical/x-ray surveys
- For any survey, careful modeling will be required to determine this precisely, understand uncertainties

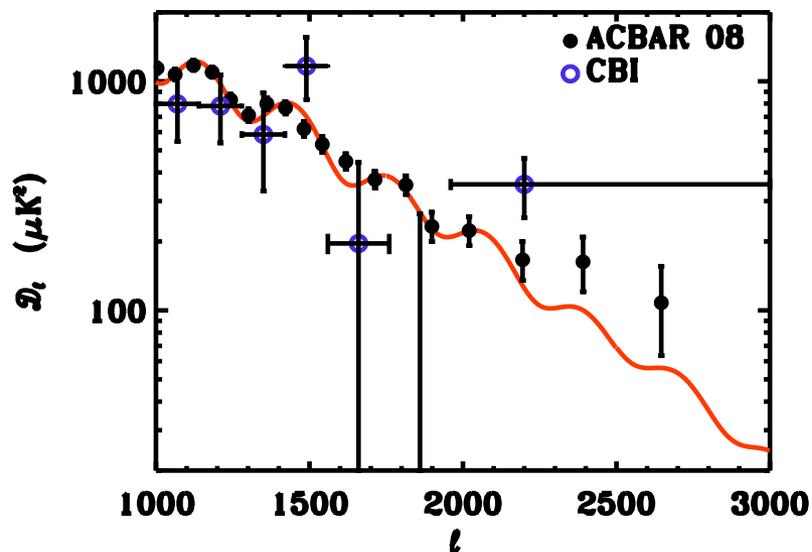


For the Near Term: High- ℓ SZ Anisotropy



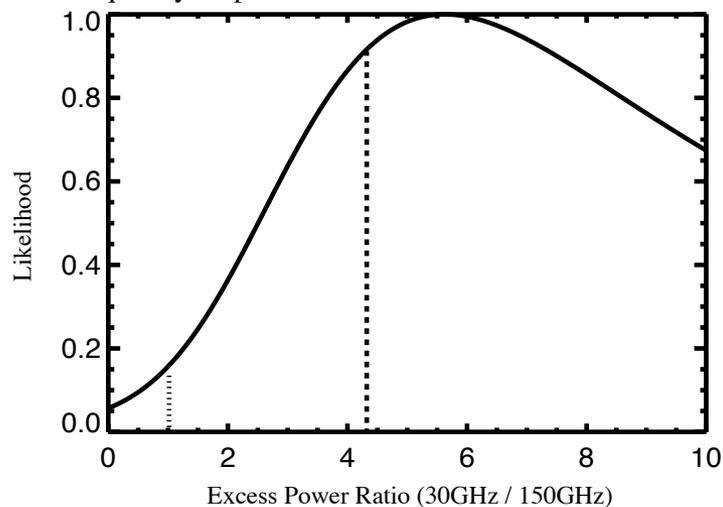
Readhead et al 2004,
modified

- In the absence of cluster detections and redshifts, use anisotropy power spectrum to describe effect
- Tentative detection at high- ℓ by CBI, ACBAR, BIMA,



SZ Survey Results from Bolocam

Frequency Dependence of the ACBAR and CBI Excess

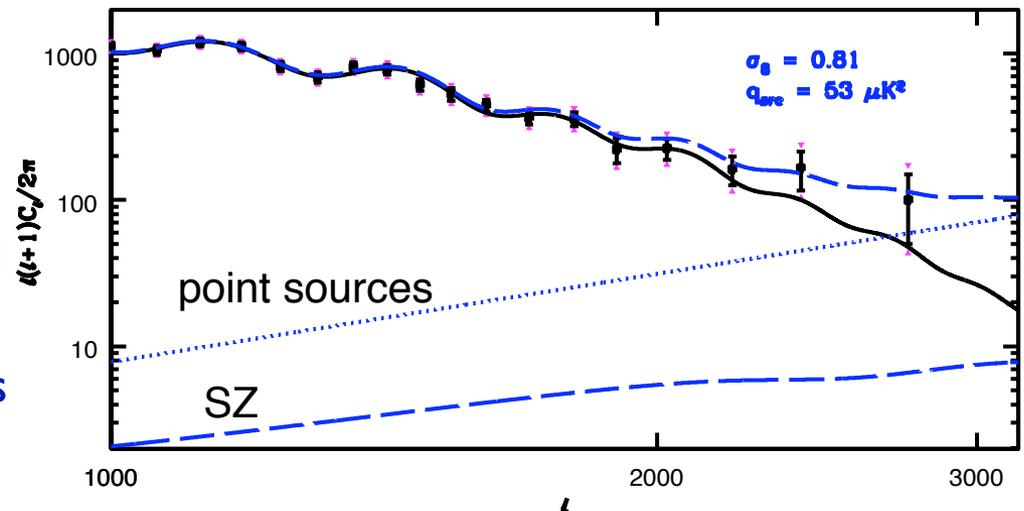


Reichardt et al 2008,
modified

The High- ℓ Excess

- The high- ℓ excess seen by ACBAR, CBI, and BIMA is not entirely consistent with a SZ anisotropy explanation
 - SZ anisotropy expected to scale as $\sigma_8 (\Omega_b h)^2$;
constraint on high- ℓ excess yields constraint on σ_8
 - CMB primary anisotropy + LSS also yields constraint on σ_8
 - ACBAR + WMAP3 primary PS + LSS $\rightarrow \sigma_8 = 0.81-0.85 \pm 0.03$
 - ACBAR + CBI excess interpreted as SZ $\rightarrow \sigma_8 = 0.95 \pm 0.04$
 - Dawson et al (2006) BIMA point: $220 \pm 130 \mu\text{K}_{\text{CMB}}^2$ at 30 GHz
 $\rightarrow 55 \pm 33 \mu\text{K}_{\text{CMB}}^2$ at 150 GHz vs. $< 10 \mu\text{K}_{\text{CMB}}^2$ for $\sigma_8 = 0.80$
 - ACBAR + WMAP3 can be reasonably interpreted as $\sigma_8 \sim 0.80$ SZ + unidentified point sources
 - Need better data!

Reichardt et al 2008, modified



SZ Anisotropy: RJ Interferometers

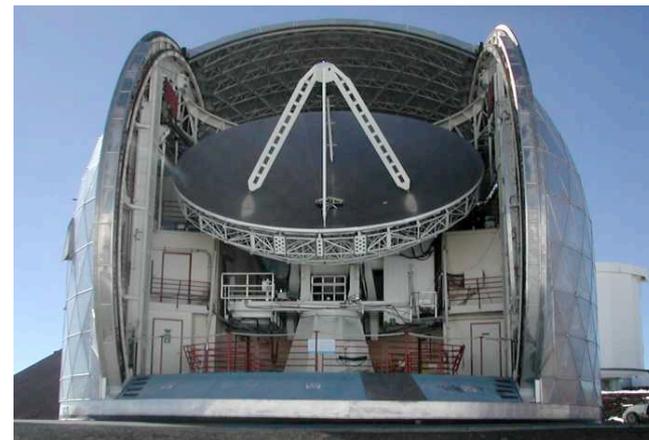
- Experiments:
 - Sunyaev-Zeldovich Array: Carlstrom et al at CARMA site, 8 x 3.5 m dishes at 26-36 GHz and 85-115 GHz + CARMA
 - Arcminute Microkelvin Imager: MRAO, MRAO site, 10 x 3.7 m + 8 x 13 m, 12-18 GHz



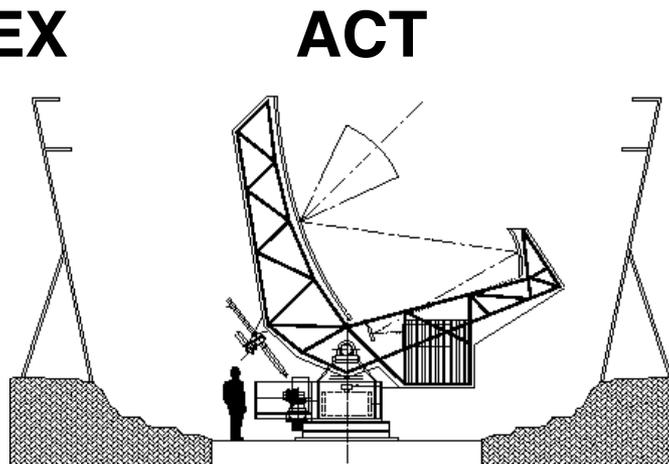
SZ Anisotropy: MM-Wave Arrays

- mm-wave experiments (in order of existence and site quality)
 - Bolocam: 120 pixels at 150 GHz on 10.4 m CSO, Mauna Kea
 - APEX: 300 pixels at 150 GHz on 12 m ALMA prototype, ALMA site
 - ACT: 1000 pixels each at 150, 220, 275 GHz on 6-m off-axis az-scanning dish, Cerro Toco
 - SPT: 1000 pixels distributed across 90, 150, 220 GHz bands on 10-m off-axis dish, South Pole

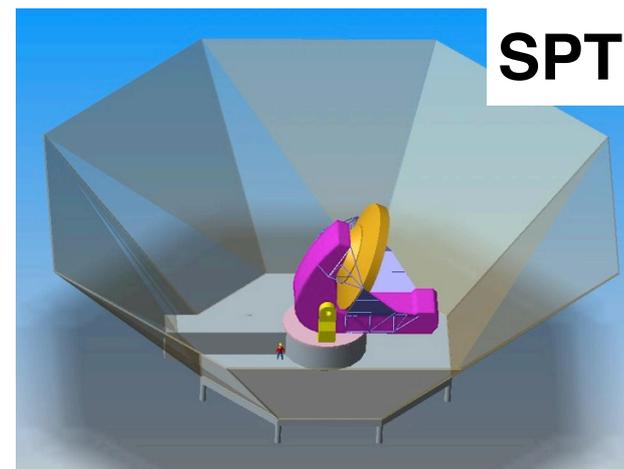
Bolocam/CSO



APEX



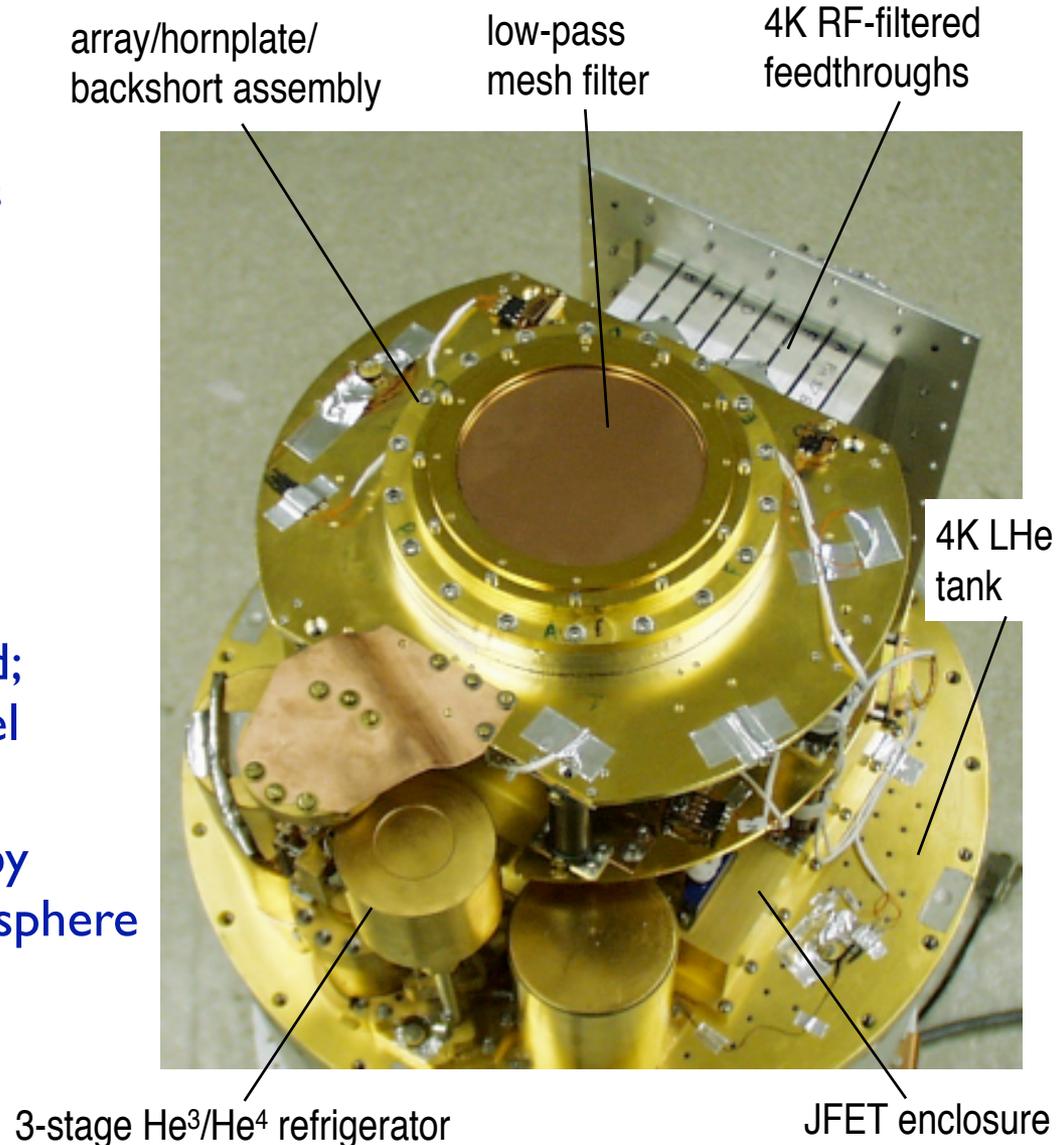
ACT



SPT

Bolocam Overview

- **Observation bands:**
 - 125-165 GHz: thermal SZ
 - 225-300 GHz: dusty sources
 - (217 GHz: kinetic SZ)
- 144-pixel spiderweb bolometer array operated at ~ 250 mK
- **Array architecture:**
 - Bolometers are bgnd-limited; increase sensitivity with pixel count (8' FOV)
 - Sky noise removal enabled by beam overlap through atmosphere
- At Caltech Submm Obs., 10-m on Mauna Kea

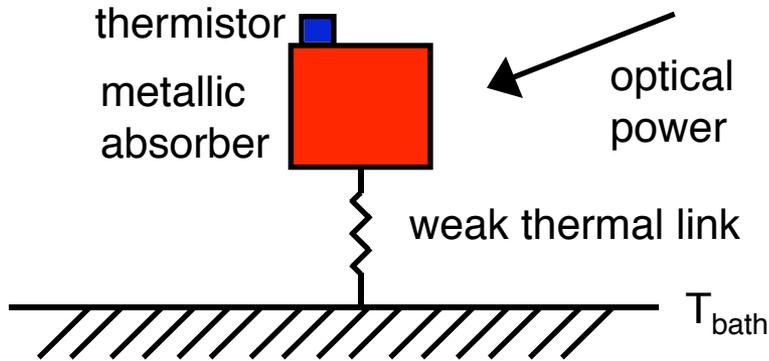


Instrument Team

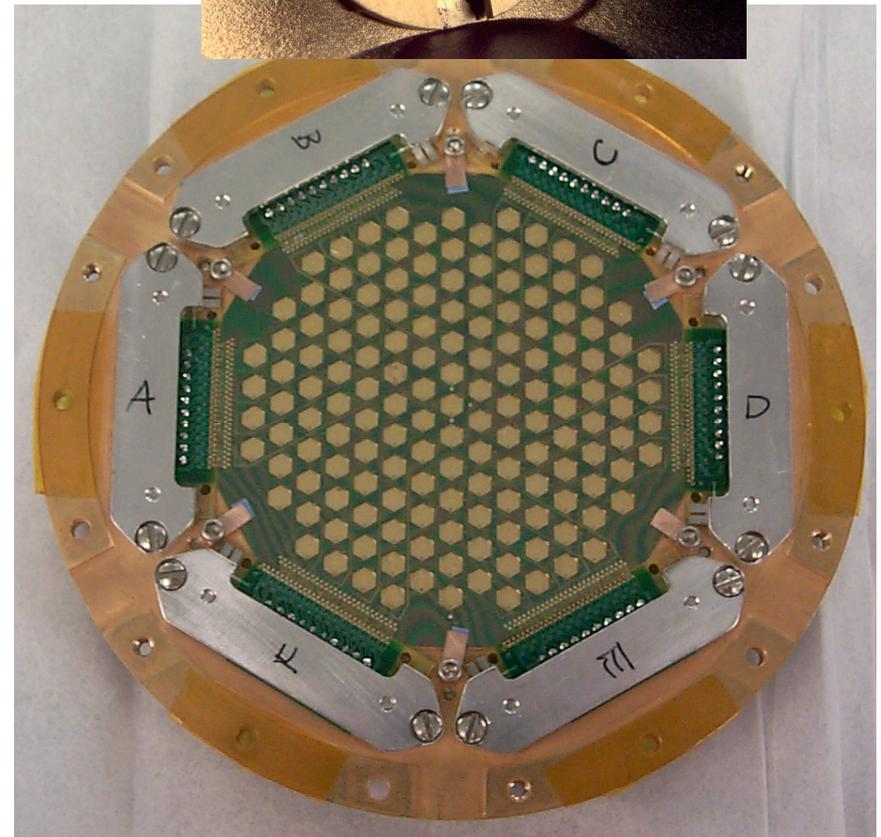
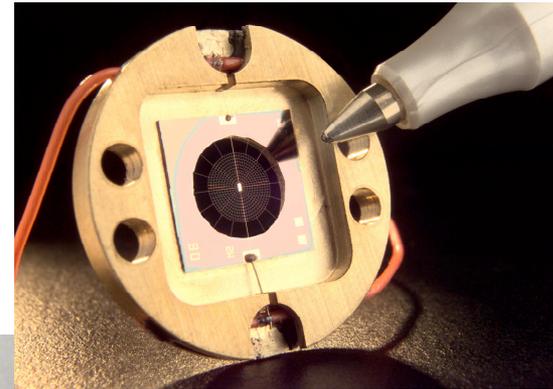
- Caltech
 - (Mihail Amarie), (Samantha Edgington), Sunil Golwala, Andrew Lange, Philippe Rossinot, Jack Sayers
- JPL
 - Jamie Bock, (Alexey Goldin), Hien Nguyen, Fab team at MDL
- University of Colorado, Boulder
 - James Aguirre, Jason Glenn, (Ben Knowles), Glenn Laurent, Phil Maloney, James Schlaerth, (Patrick Stover)
- University of Wales, Cardiff
 - Peter Ade, Douglas Haig, Phil Mauskopf, Rob Tucker

PhD thesis Dec 2007,
has done bulk of analysis work

Bolometer Array

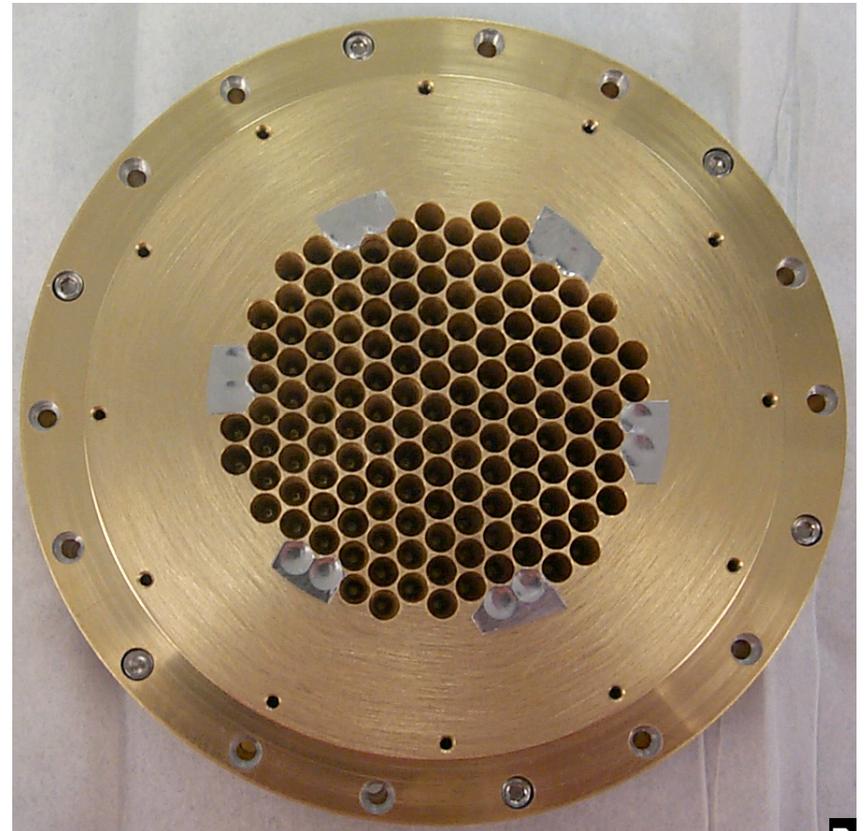


- 144 bolometers on single wafer: J. Bock, JPL/MDL
- 125 Å Au absorber on 1 μm SiN membrane, etched into “spider-web” to minimize C_{Au} , G
- NTD Ge thermistor senses T
- Array production nontrivial



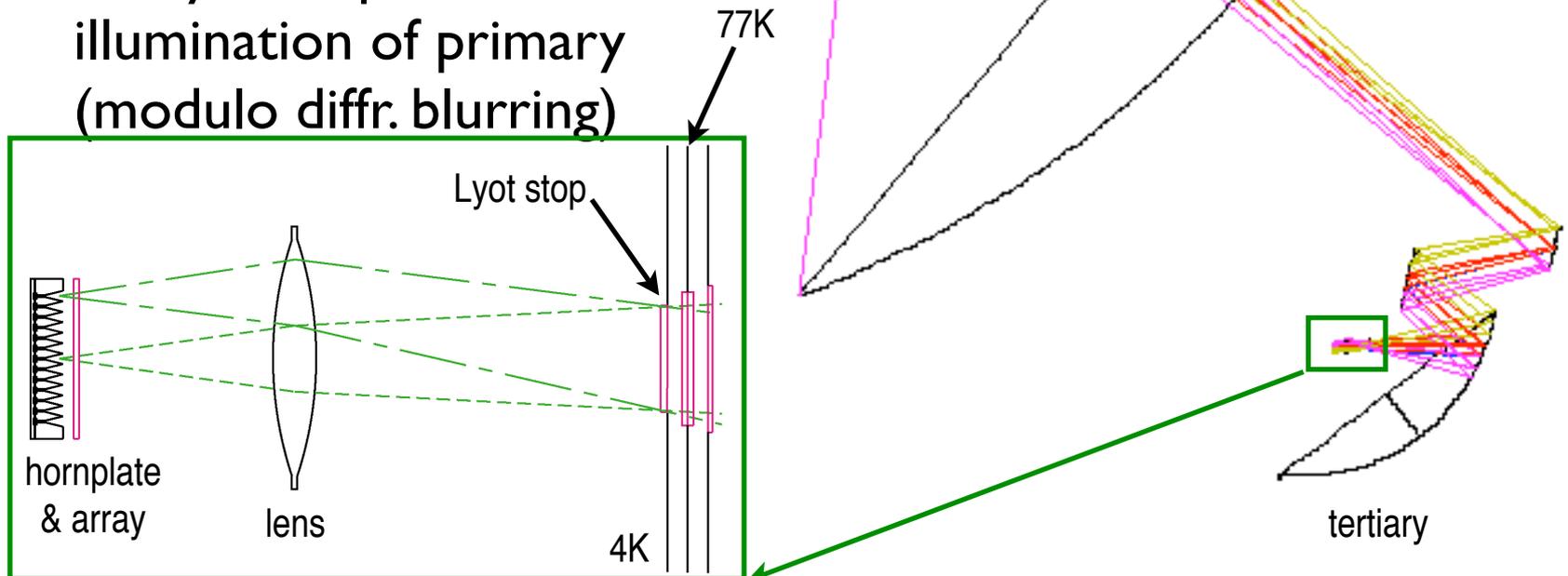
Optical Design

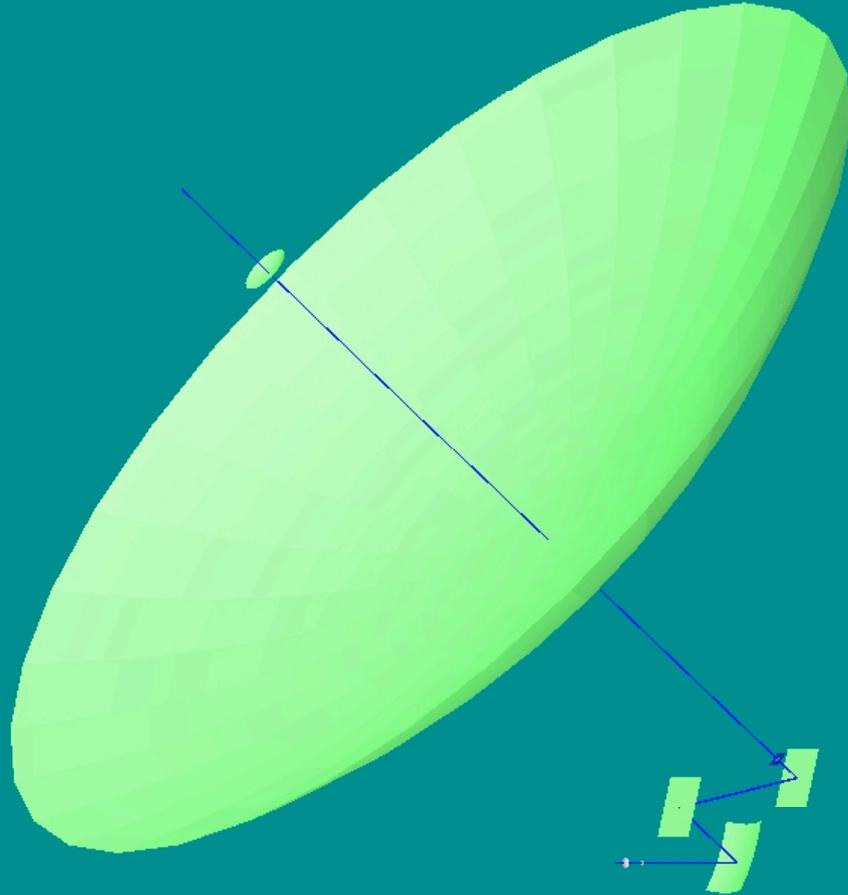
- Smooth-walled conical feedhorns define beams
- Horns coupled to integrating cavities via 2λ length of single-mode waveguide (defines lower edge of BP)
- Integrating cavities house bolos, yield $> 90\%$ efficiency and $< 1\%$ optical crosstalk
- Monolithic construction
 - single feedhorn plate
 - single backshort plate
- Backshort and hornplate can be exchanged easily
 \Rightarrow “easy” to change bands

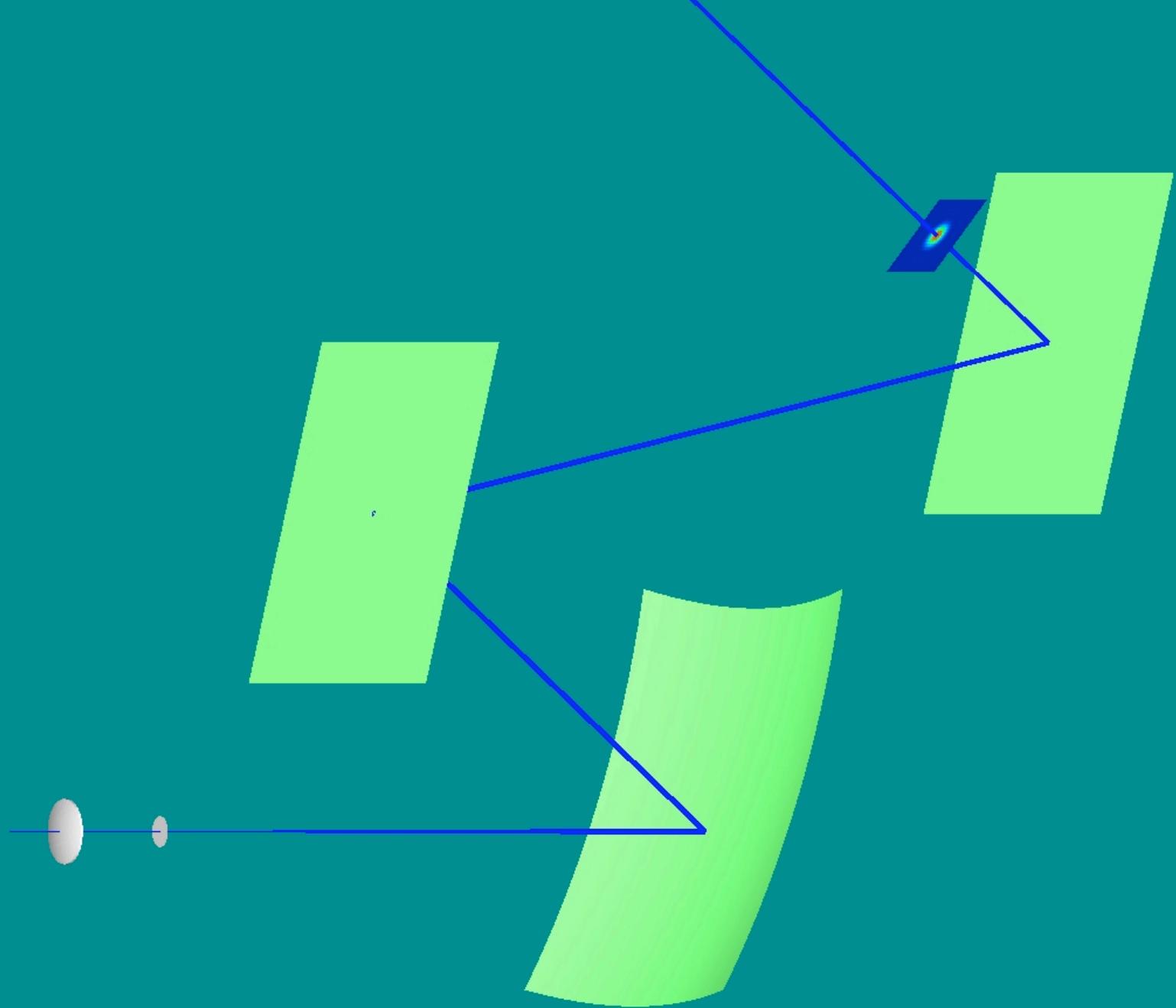


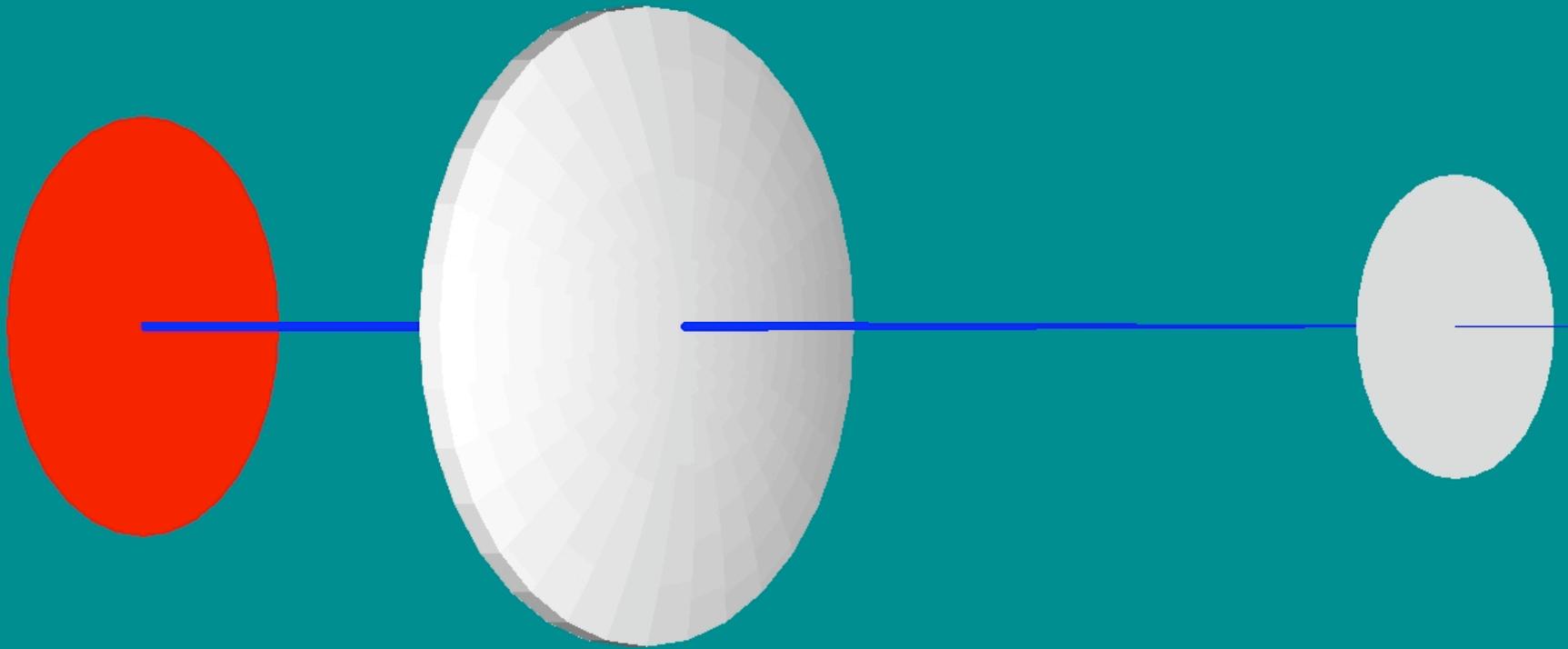
Optical Design

- Ellipsoidal tertiary and cold HDPE lens provide wide FOV (8') at F/3 plate scale
- 5 cold metal-mesh filters define high end of band while minimizing harmonics leaks
- Reflective IR blocker at 77K
- 4K Lyot stop defines illumination of primary (modulo diffr. blurring)



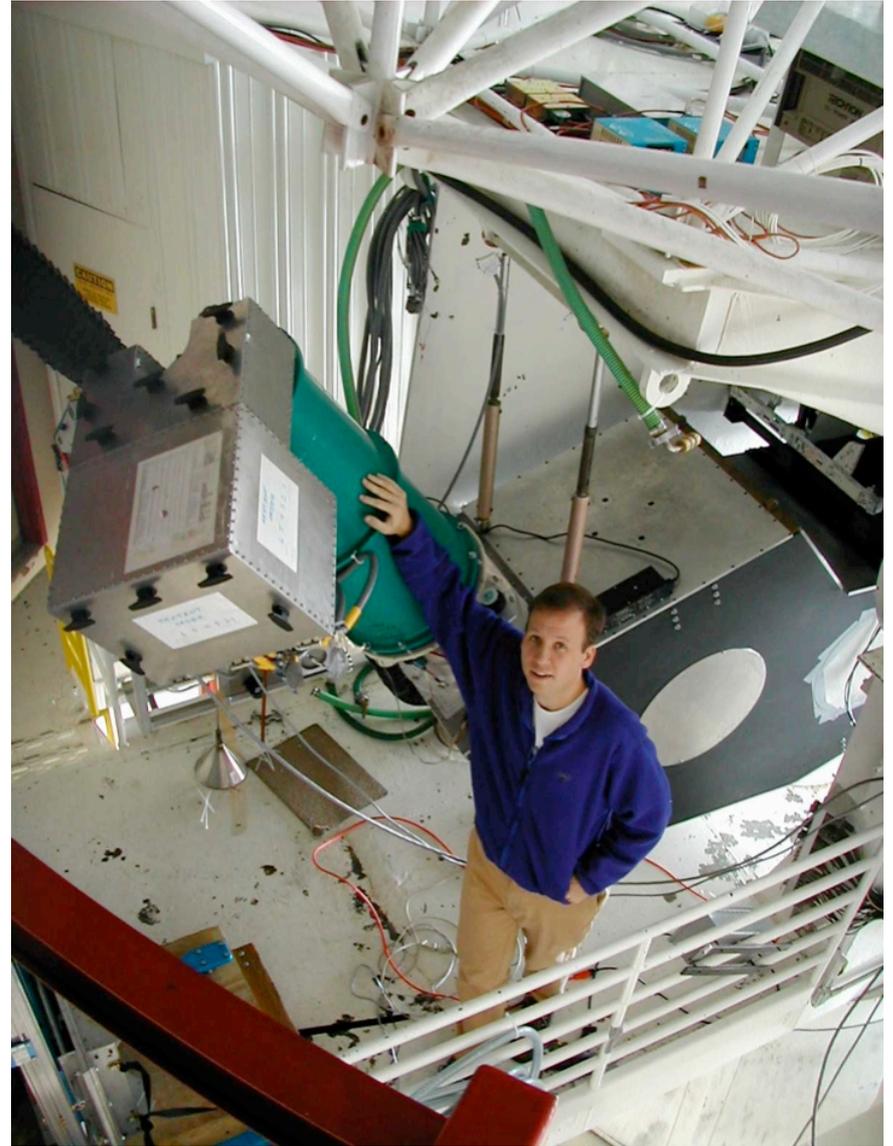
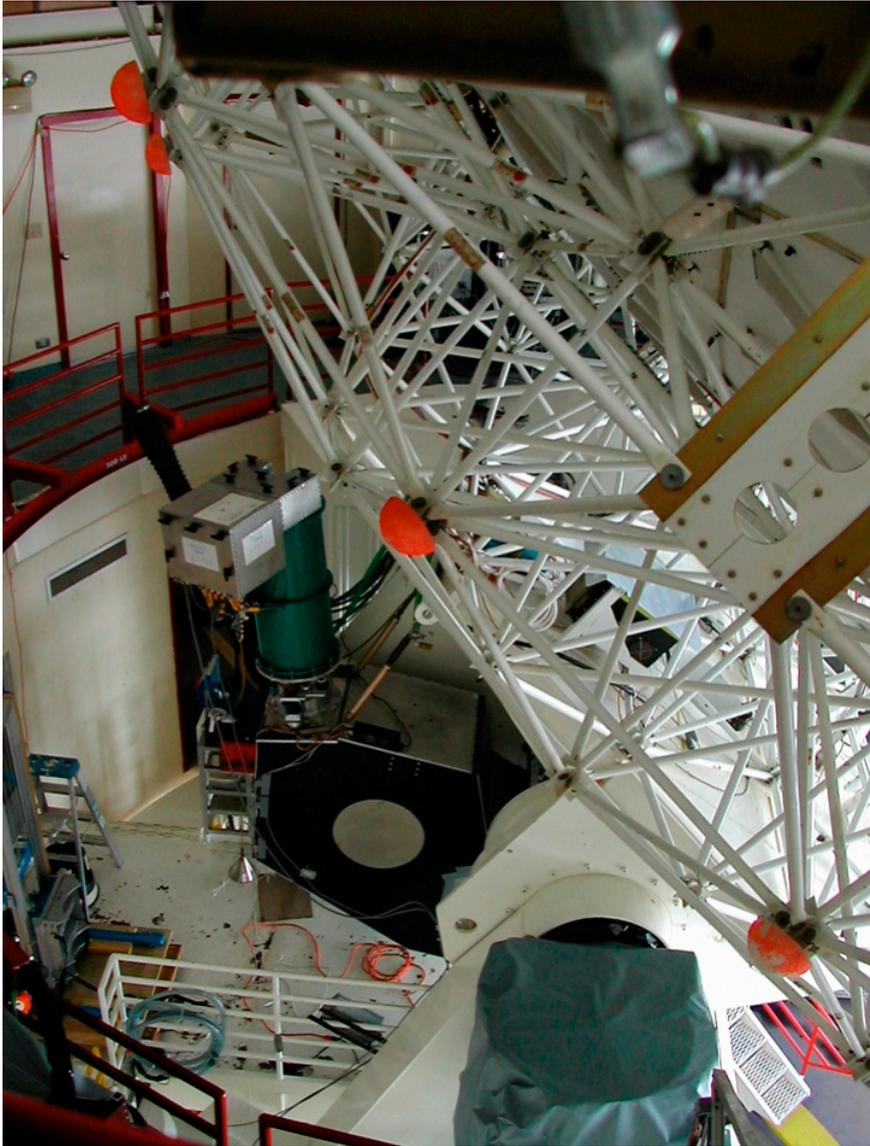












150 GHz Blind Sunyaev-Zeldovich Effect Survey

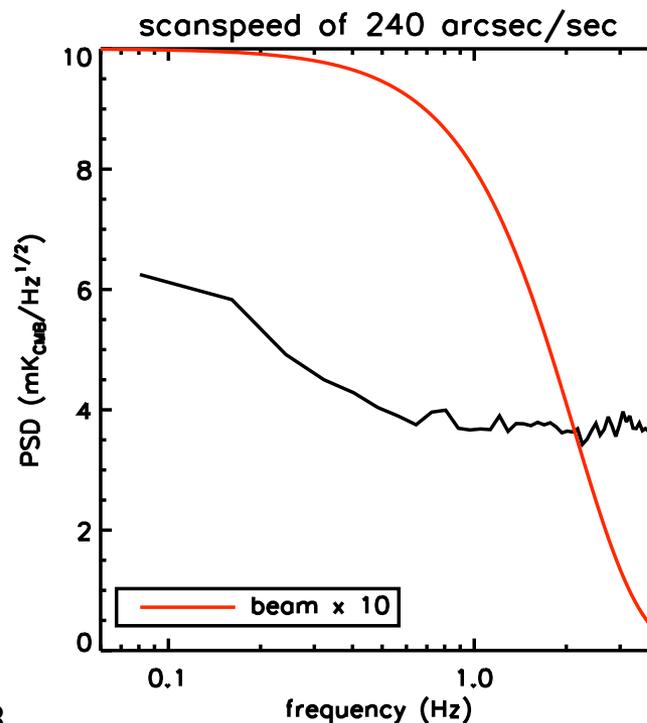
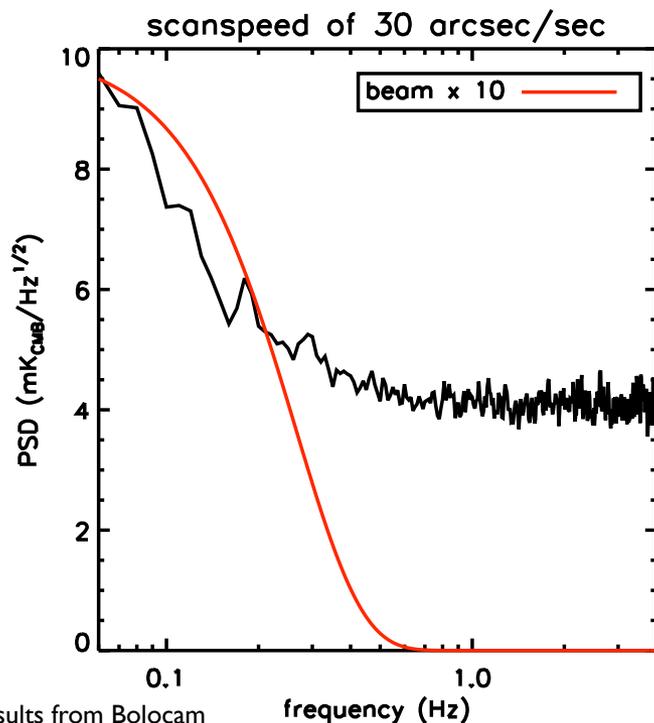
- 2 fields, each 0.5 deg^2
 - Wanted low dust emission, good X-ray and optical coverage in case clusters were found
 - SDSS (aka SXDS): Subaru deep survey field
 - 400 ksec XMM-EPIC integration time
 - OIR coverage by surveys on Subaru, CFHT Legacy, UKIRT, Spitzer SWIRE Legacy survey
 - $12 \mu\text{Jy}$ VLA coverage
 - SCUBA SHADES and BLAST field
 - 1.2 MJy/ster $100 \mu\text{m}$ dust emission, among the lowest in the sky
 - Lynx: not so well complemented
 - 150 ksec XMM-EPIC
 - imaging of small portions containing low-mass clusters
 - 1.3 MJy/ster $100 \mu\text{m}$ emission, also pretty good
- ~ 40 nights of telescope time in fall 2003

150 GHz Blind Sunyaev-Zeldovich Effect Survey

- Observing Strategy

- Spend half the night on each field, 6-8 hrs each per night
- Raster over each field along the RA and dec directions
 - Drift scan would be less prone to scan-synchronous pickup, but sky noise pushes one to active scanning to move signal to higher temporal frequency
 - Active az-only scans produce inefficient coverage pattern due to sky rotation
 - Good belief that array would allow subtraction of elevation dependent signal

after sky noise removal!

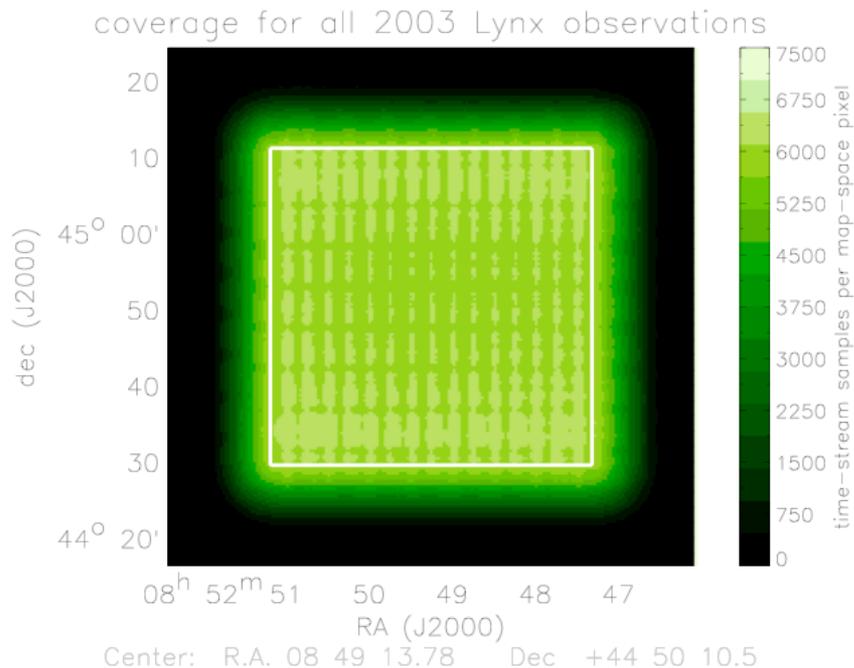
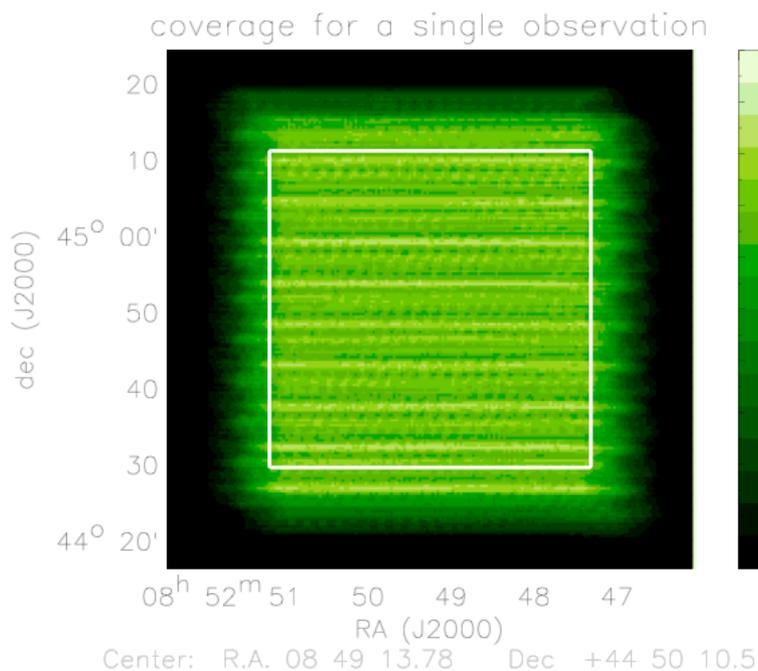


150 GHz Blind Sunyaev-Zeldovich Effect Survey

- Observing Strategy

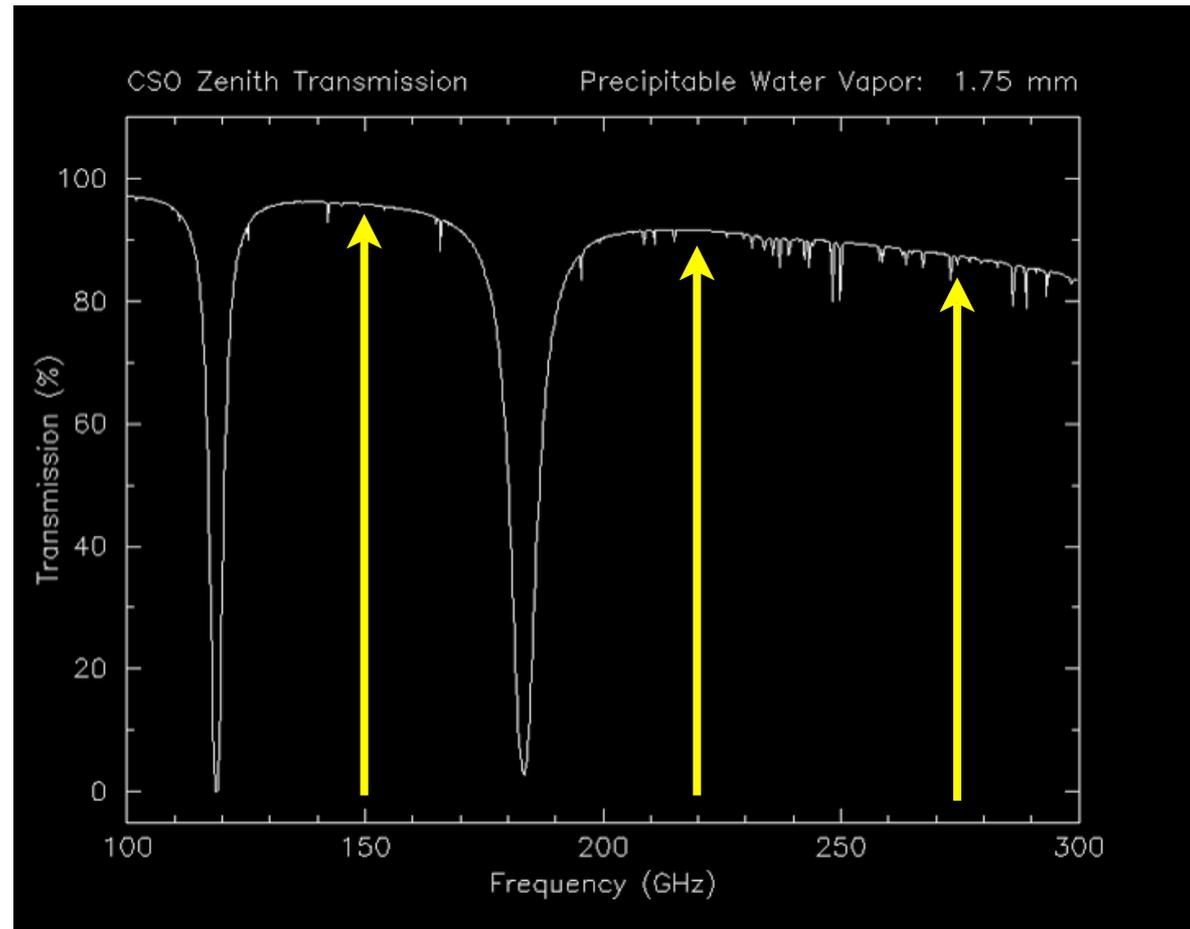
- Data broken up into 8-minute-long “observations”

- Each observation covers the entire field in one scan mode with 8-9% rms coverage variations (4-5% noise variations)
- Alternate RA and dec scans
- 3 sets of offsets perpendicular to scan direction to smooth out coverage
- Final maps have 1.5% coverage variations



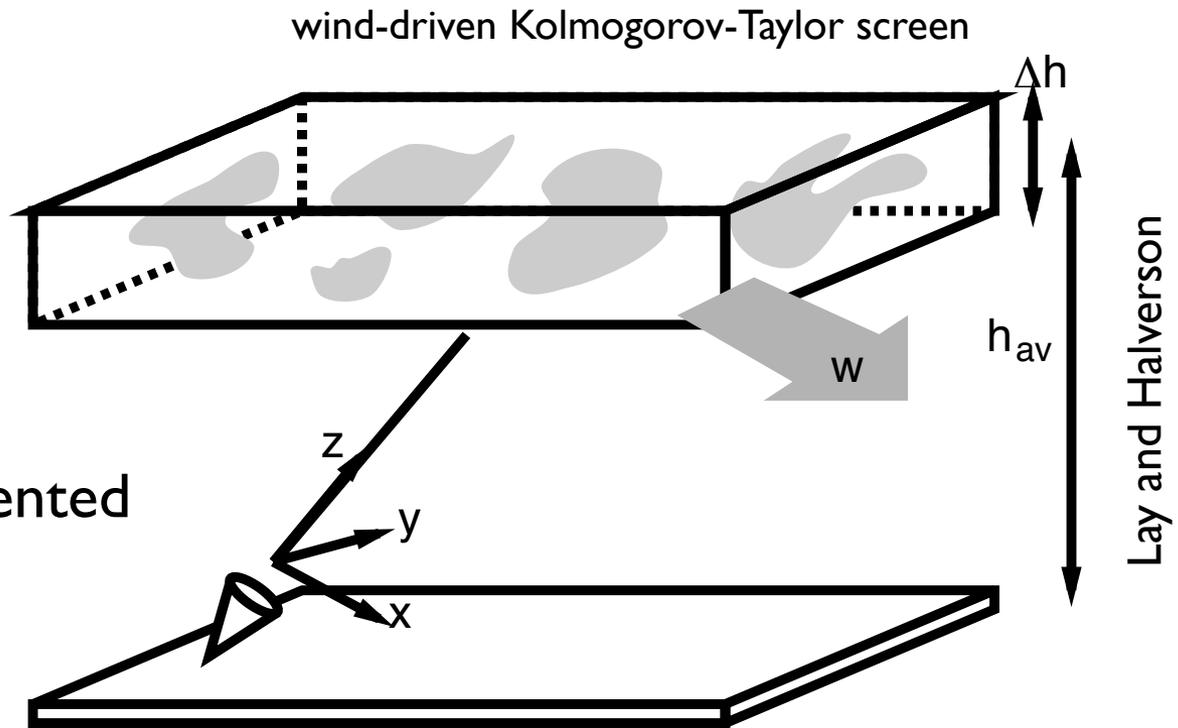
Observing Conditions: Loading and Opacity

- To first order, the atmosphere!
- Median conditions: 1.75 mm of water between the instrument and the CMB!
- Atmospheric optical depth:
 - 150 GHz: $\tau \sim 0.05$
 - 275 GHz: $\tau \sim 0.13$
- Photon Poisson and Bose noise from the emitted power

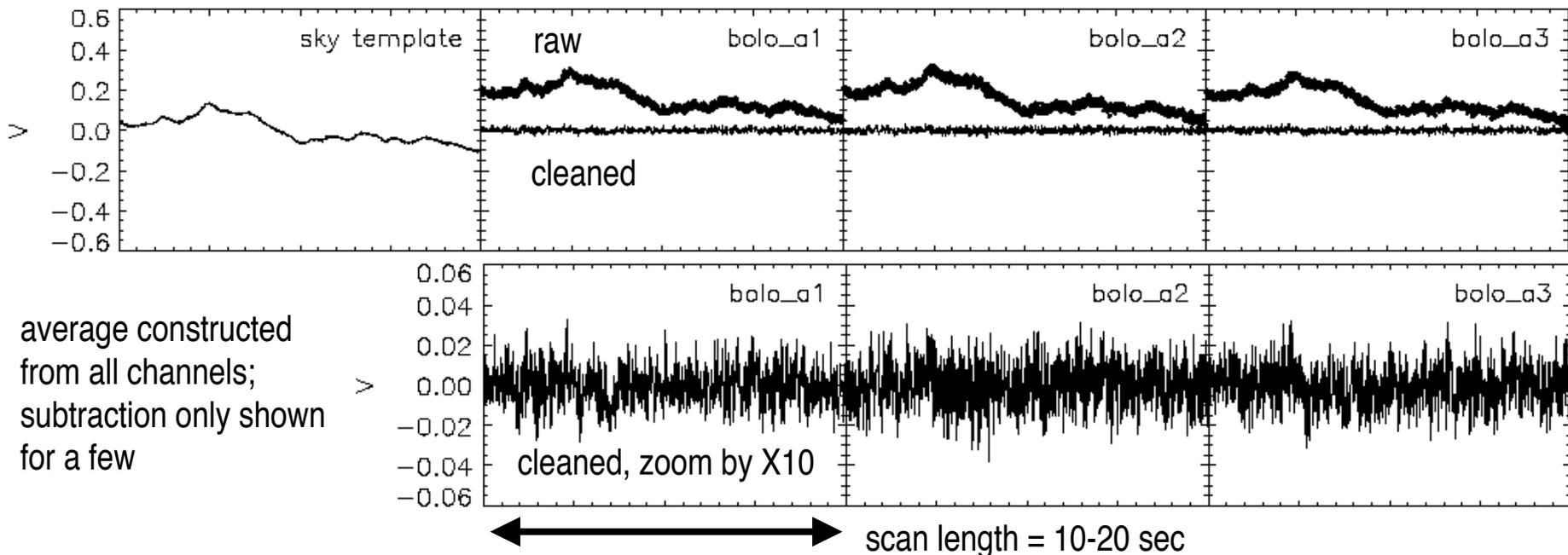


Observing Conditions: Sky Noise

- water vapor w/scale height of ~ 2 km
 - near condensation point, so clumpy
 - strong dipole moment \Rightarrow rotation couples well to mm-waves
- liquid water: same modes, but much less efficient, constrained by inter-molecule forces
- ice: rotation is prevented
- Water vapor present as turbulent screen entrained in wind



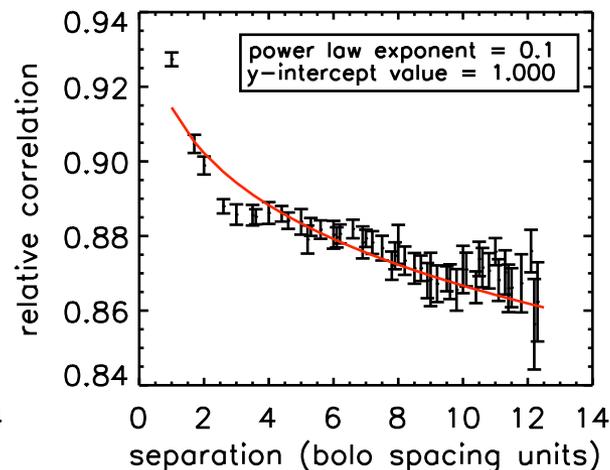
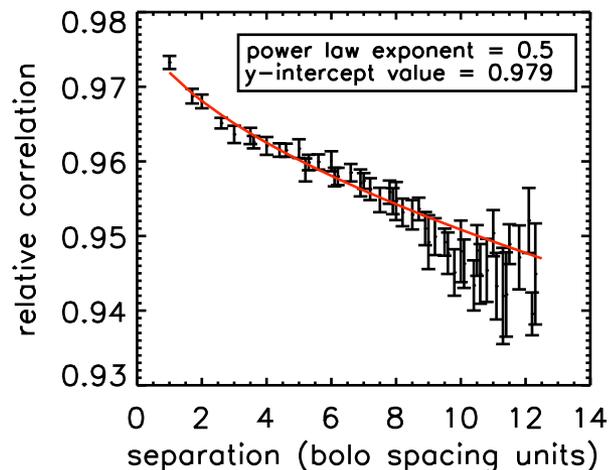
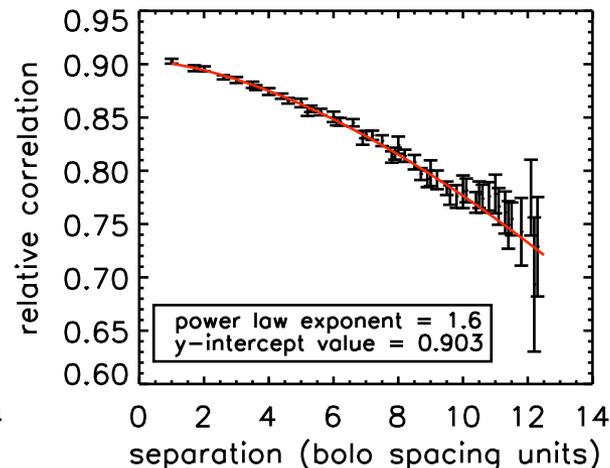
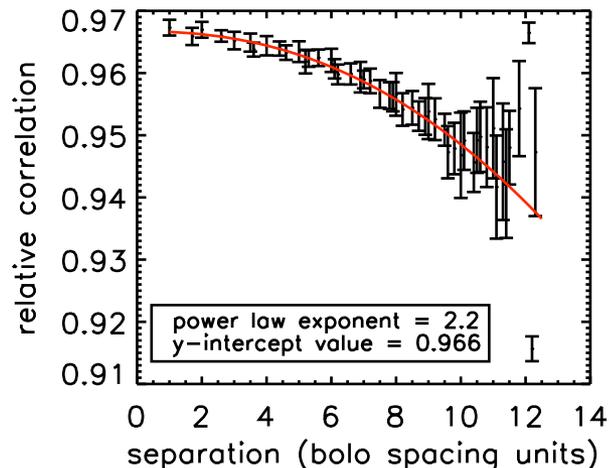
Observing Conditions: Sky Noise



- Sky noise = fluctuations in emission from water vapor in atmosphere due to wind-driven turbulent screen
- Overlap of beams through atmosphere ensures it is mostly common signal
- A simple average removal takes out >90% of sky noise

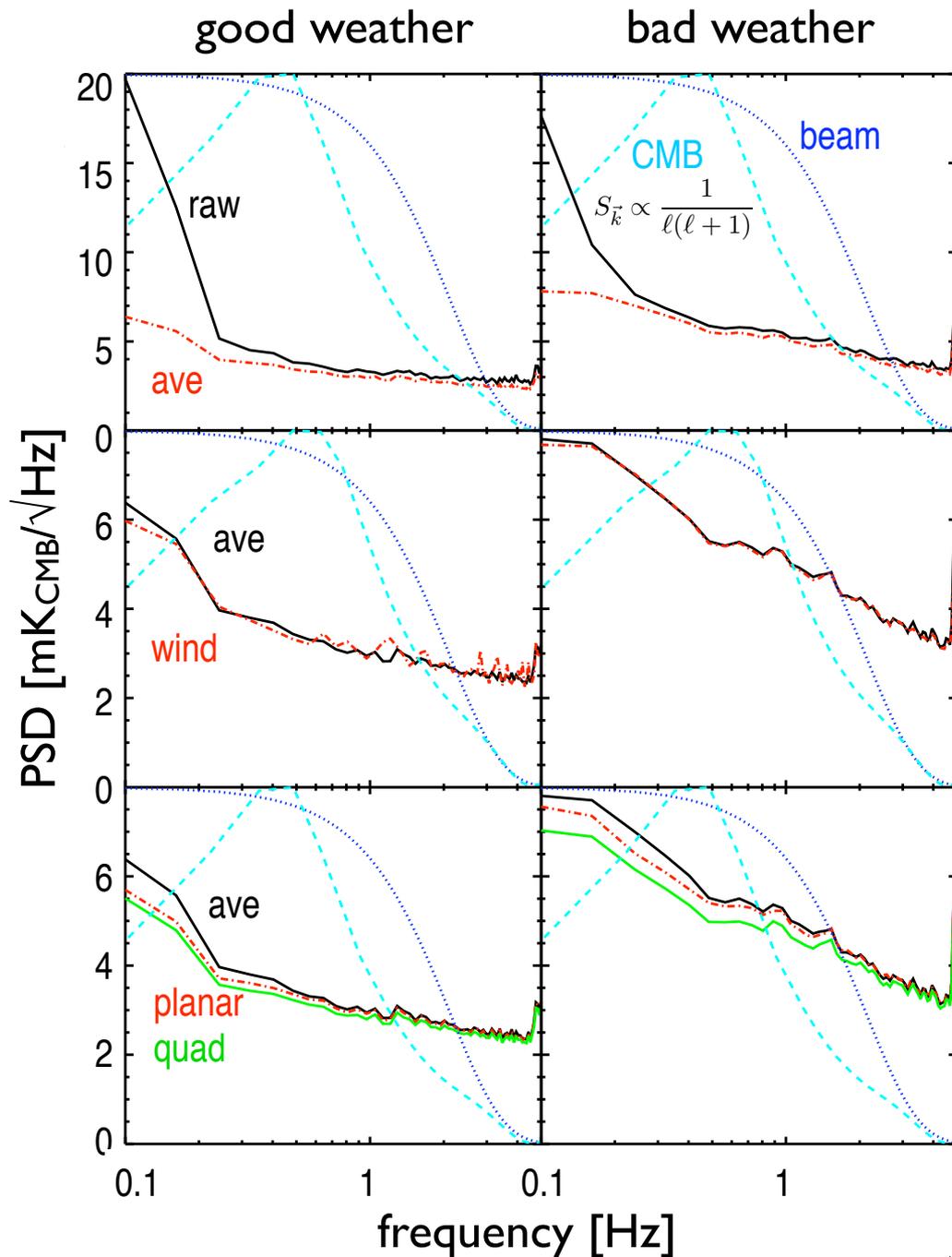
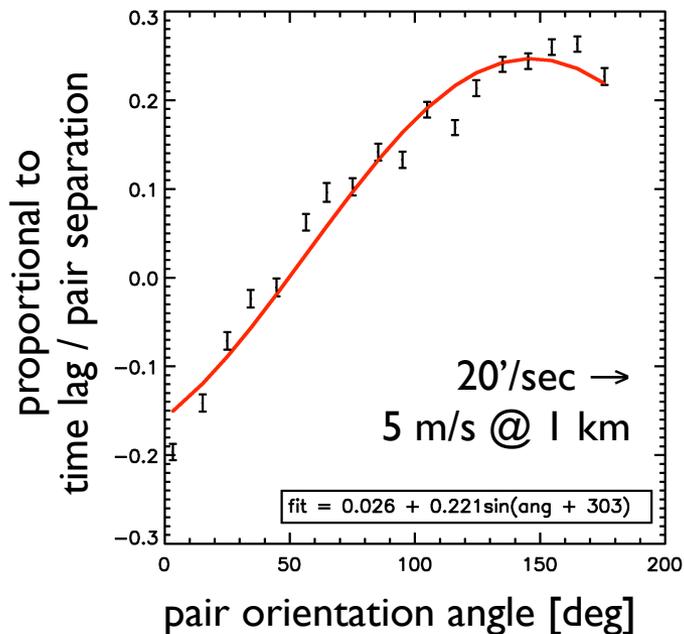
Autocorrelation Function of Sky Noise

- See expected power-law autocorrelation function of sky noise as a function of pixel separation (structure function)
- Correlation length varies; large corr. length \rightarrow good sky subtraction
- Excess correlation visible at small separations, worst when sky noise is poor. Consistent with spread of Airy function.



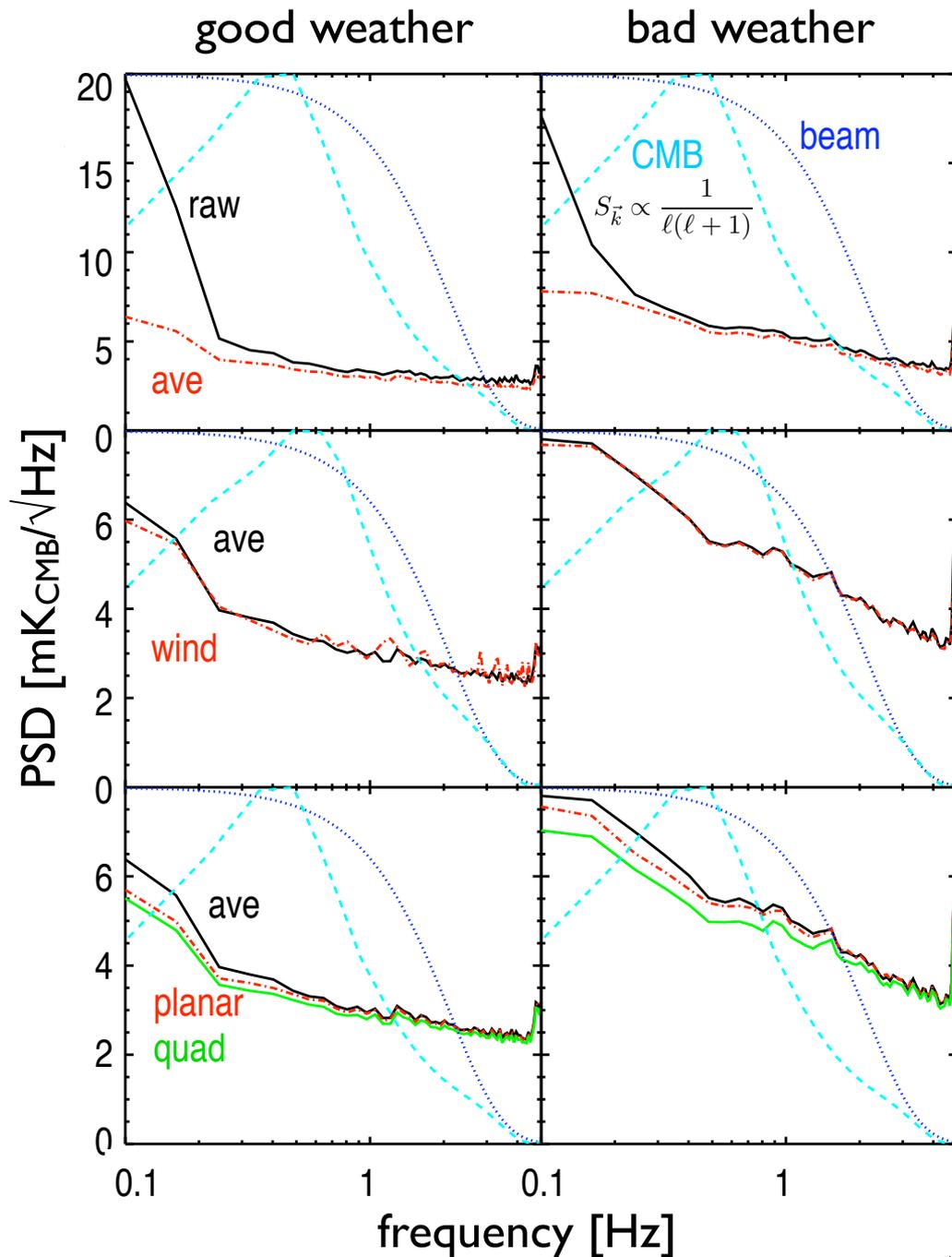
Sky Noise

- Average removal leaves significant noise above fundamental photon + instrument noise
- First, attempt to model as wind-driven screen: get sensible wind speeds, but no improvement



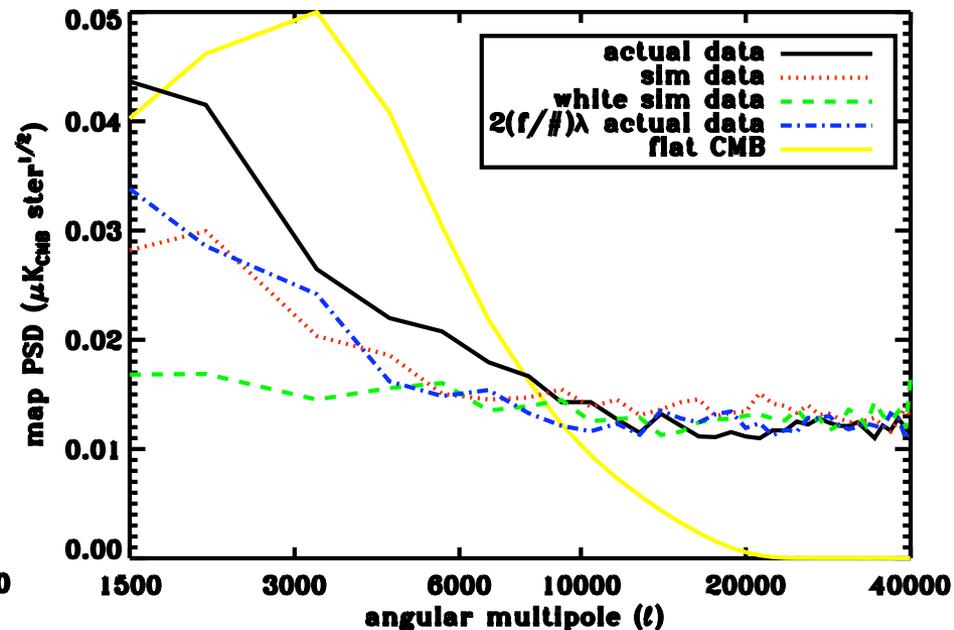
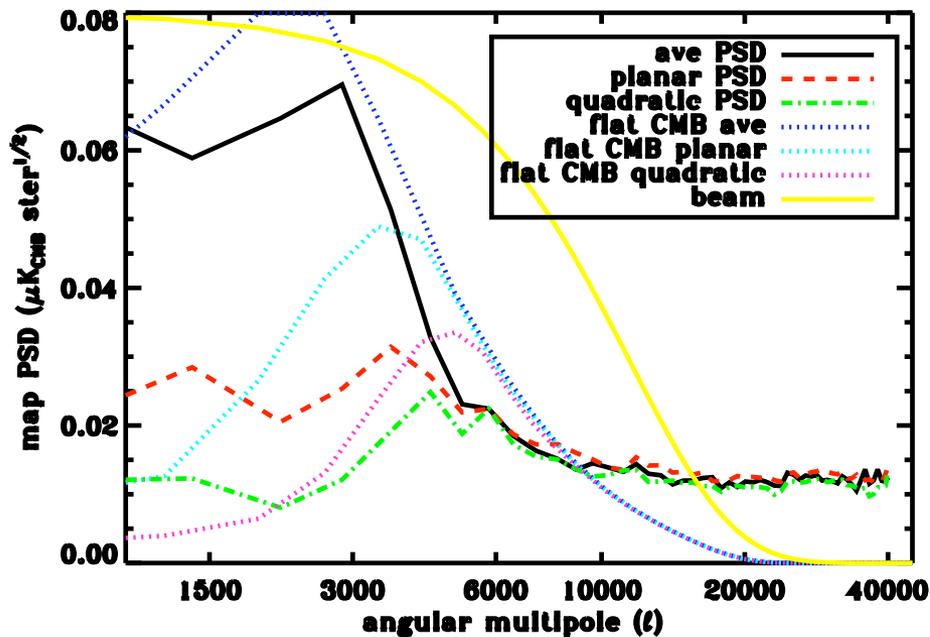
Sky Noise

- Think a bit harder:
 - typical wind speed: 10 m/s @ 1km = 35'/sec
 - telescope scan speed = 4'/sec \ll wind speed so neglect telescope motion
 - noise is below 0.5 Hz; $T = 2$ sec, $w = 35'$ /sec get $\theta = wT > 1$ deg $\gg 8'$ FOV
- \Rightarrow on scale of array, see only polynomial-like portion of mode
- fit for average, plane, or quadratic across FOV



Map-Space PSDs

- Subtraction methods similar in timestream, differ in map space
 - (Naive mapmaker; see below for more sophisticated version)
 - Residual correlations manifest as low- ℓ noise
 - More aggressive methods reduce residual correlations among bolometers



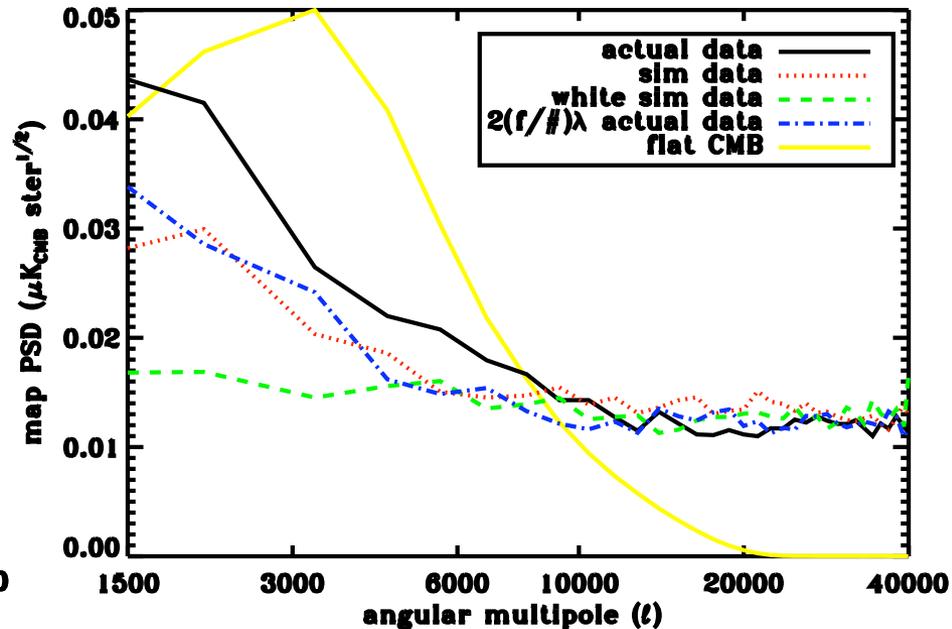
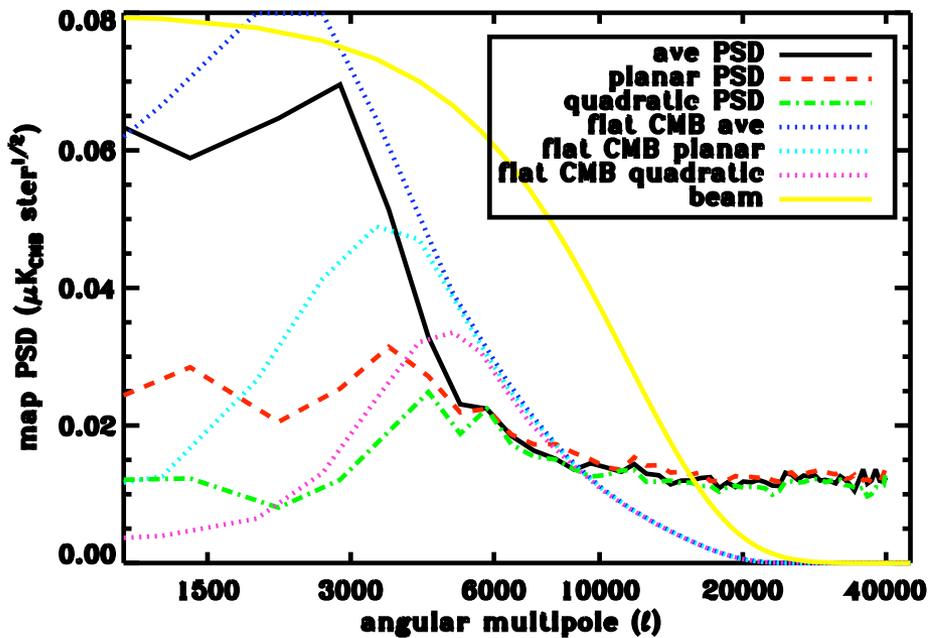
(transfer functions not deconvolved)

Map-Space PSDs

data type	PSD spectrum	PS amplitude uncertainty
actual/36 spaced detectors	data	$550 \mu\text{K}_{\text{CMB}}^2$ ←
actual/115 detectors	data	$270 \mu\text{K}_{\text{CMB}}^2$ ←
sim/115 detectors	data	$170 \mu\text{K}_{\text{CMB}}^2$ ←
sim/115 detectors	instrument, white	$100 \mu\text{K}_{\text{CMB}}^2$

consistent with \sqrt{N}

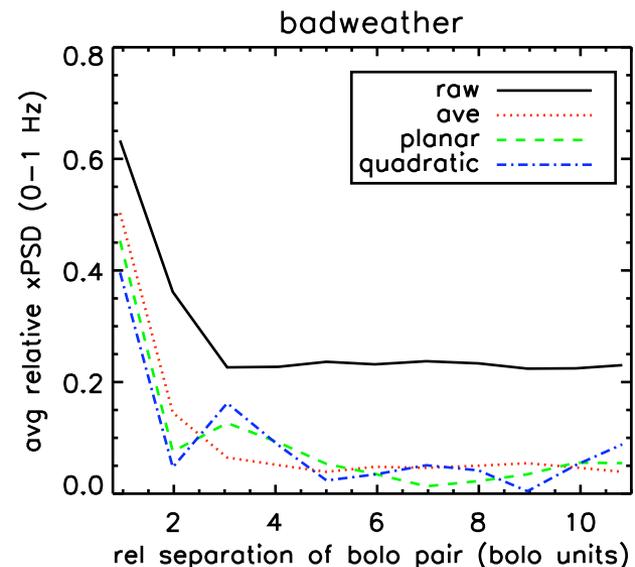
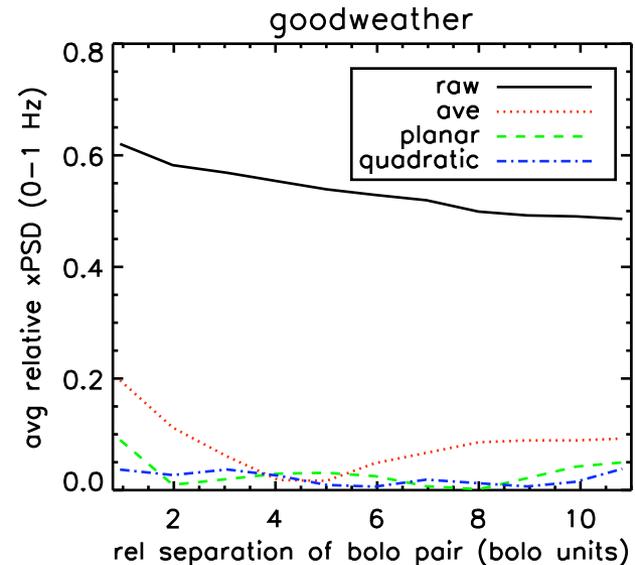
inconsistent with \sqrt{N}



(transfer functions not deconvolved)

Residual Spatial Correlations

- There is residual correlation between nearby bolometers post sky-subtraction
 - lower in better weather
 - excess correlation at sub- $(f/\#)\lambda$ separations
 - one bolo separation = $0.7 (f/\#)\lambda$
 - Need to go out to $r \sim 2 (f/\#)\lambda$ before residual correlations look flat with r
 - Effective number of pixels drops by a large factor:
degradation in $\mu\text{K}_{\text{CMB}}^2$
 \sim degradation in number of pixels

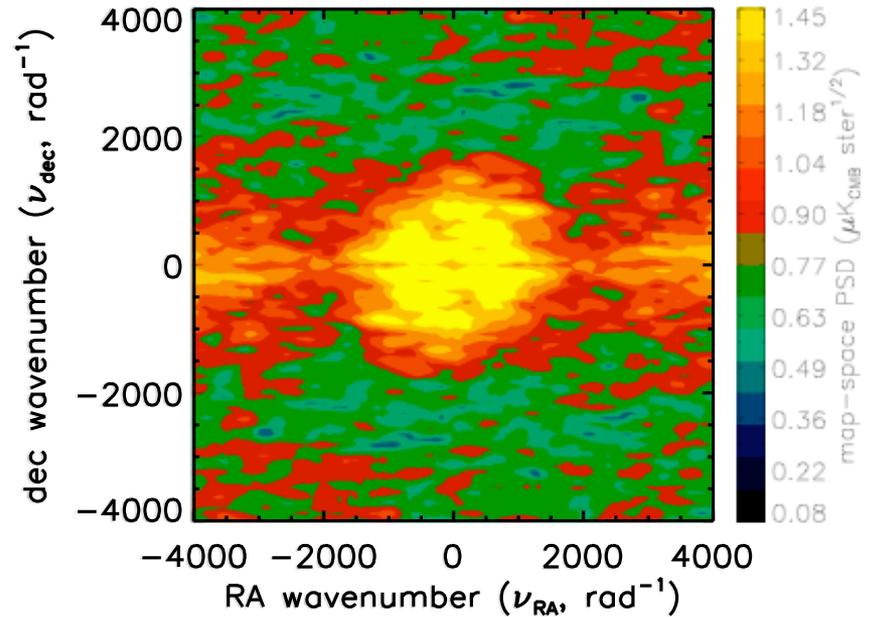
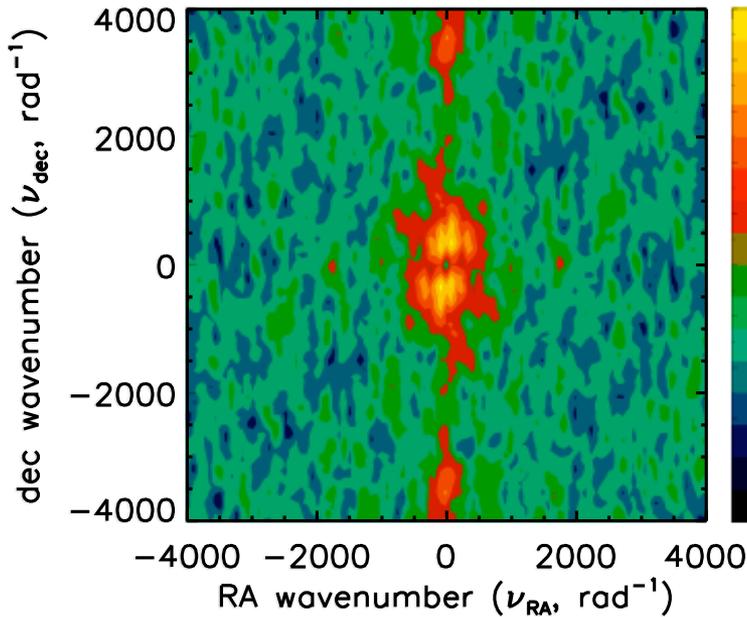


Map-Space PSDs

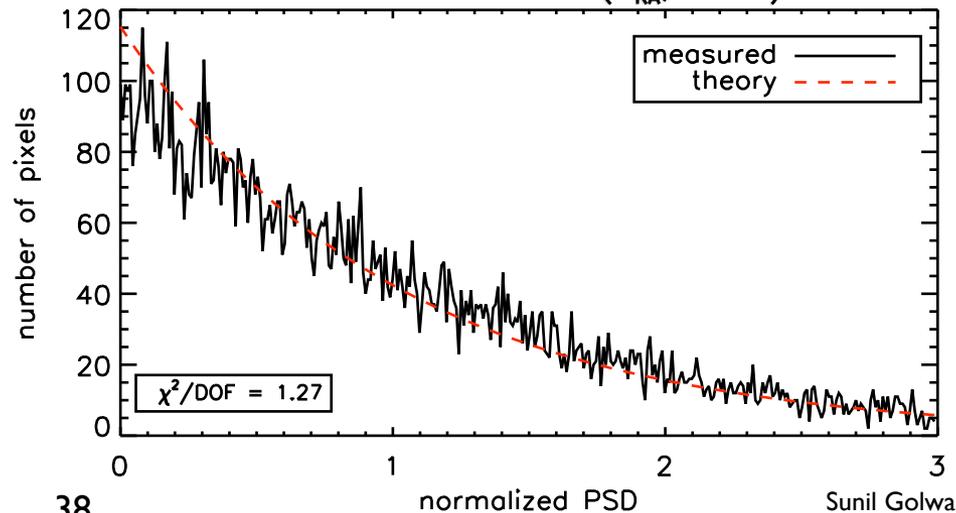
good weather
RA scan

$$\ell = 2\pi\nu$$

bad weather
dec scan



- Scan type is evident in map-space PSDs
- Variations in low-frequency noise clear
- Very gaussian except at small deviations



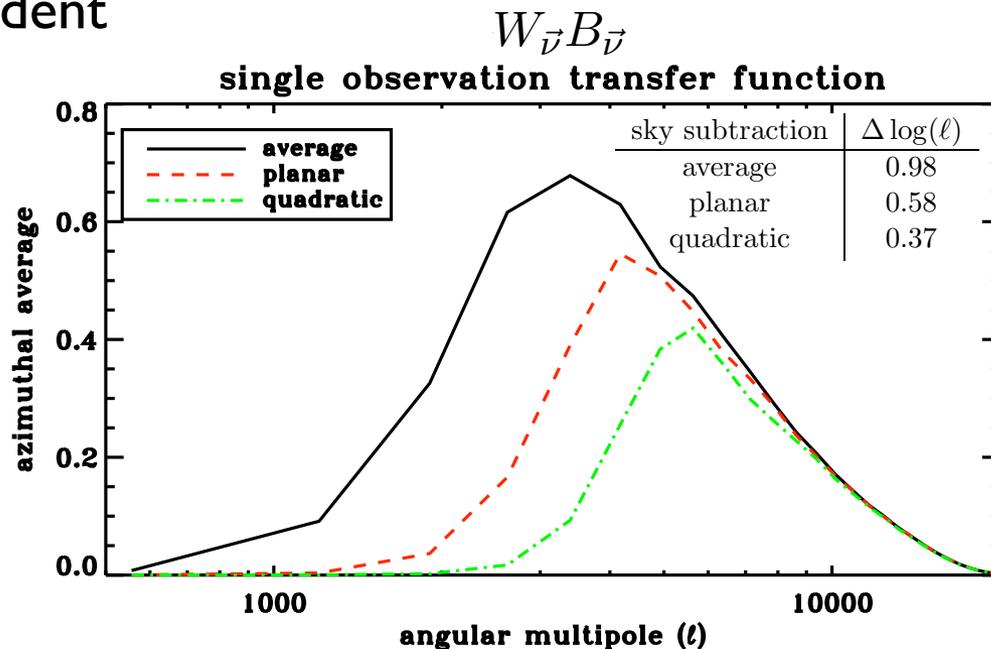
Sky Noise Removal Window Function

- Sky noise removal via correlation analysis reduces sensitivity to signal on scales $\gtrsim 8'$ FOV. $\lambda_{\theta} = 8' \Rightarrow \ell = 2700$
- More aggressive sky noise removal also removes more signal
- Measure transfer function of sky noise removal by inserting simulated CMB (flat in $\ell(\ell+1)C_{\ell}$) into timestreams and measuring attenuation at output map as a function of ℓ
- Transfer function is independent of signal amplitude at signal levels of interest
- $BW_{\text{eff}} = \Delta \log(\ell)$

$$BW_{\text{eff}} = \int_{\vec{\nu}} d\vec{\nu} S_{\vec{\nu}} W_{\vec{\nu}} B_{\vec{\nu}}$$

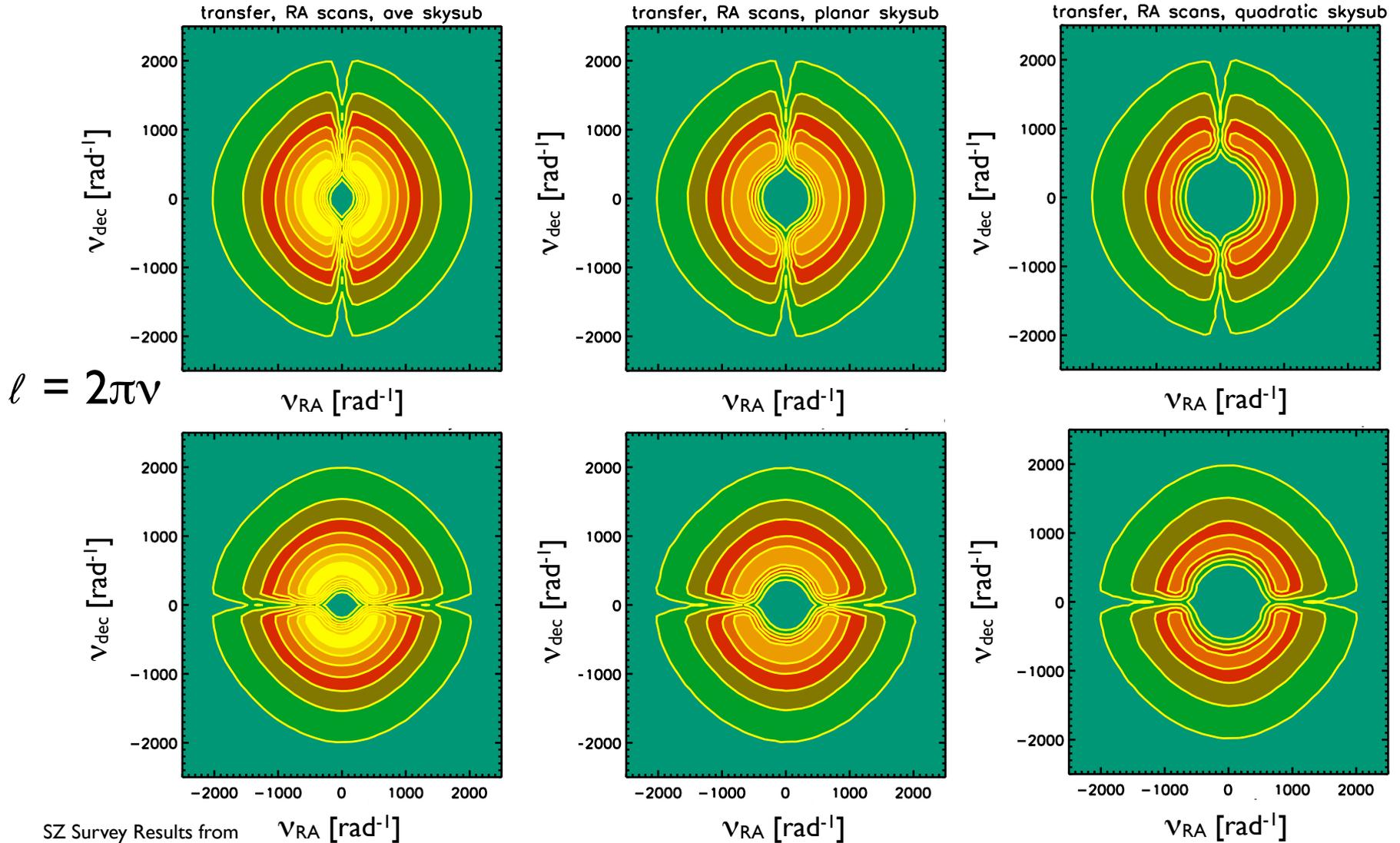
$$S_{\vec{\nu}} \propto \frac{1}{\ell(\ell+1)}$$

$$\ell = 2\pi\nu$$



Single-Observation Transfer Functions

- Asymmetry from scan pattern evident



Mapmaking

- Standard Max Likelihood mapmaking is difficult for us

$$\chi^2 = (\vec{d} - \mathbf{p}\vec{m})^T \mathbf{w} (\vec{d} - \mathbf{p}\vec{m}) \Rightarrow \vec{m}' = \mathbf{c}\mathbf{p}^T \mathbf{w}\vec{d}, \quad \mathbf{c} = (\mathbf{p}^T \mathbf{w}\mathbf{p})^{-1}$$

pointing matrix (map to timestream) map inverse of timestream noise covariance timestream max-L map estimator map-space covariance matrix (constructed from timestream covariance matrix)

- would need to include bolo-bolo correlations in timestream noise covar.
- \mathbf{c} requires inversion of $N_{\text{pix}}^2 = 16000^2$ matrix
- Simulation-based techniques have been used to deal with this
- We use hybrid method
 - Scan pattern \Rightarrow naive maps are pretty close to optimal for a single obs.
 - Stationarity of noise in each map \Rightarrow map covar. is diagonal in maps space, well describe by simple map PSD
 - Coadd observations in Fourier space with map PSD inverse var. weighting
 - Jackknives and sims used to determine transfer function and uncertainties

Pseudo-Optimal Mapmaking

- Optimizing sky noise removal
 - Optimal sky noise removal algorithm depends on the day's weather
 - Pick algorithm (ave, planar, quadratic) based on single-obs figure of merit (essentially, single-obs variance on power spectrum bandpower)

$$\text{FOM} = \sum_{\vec{\nu}} \frac{S_{\vec{\nu}}^2 W_{\vec{\nu}}^2 B_{\vec{\nu}}^2}{\mathcal{P}_{\vec{\nu}}^2} \quad S_{\vec{\nu}} \propto \frac{1}{\ell(\ell + 1)} \quad W_{\vec{\nu}} B_{\vec{\nu}} = \text{transfer function}$$

(doesn't involve the real map, just the single-obs PSD)

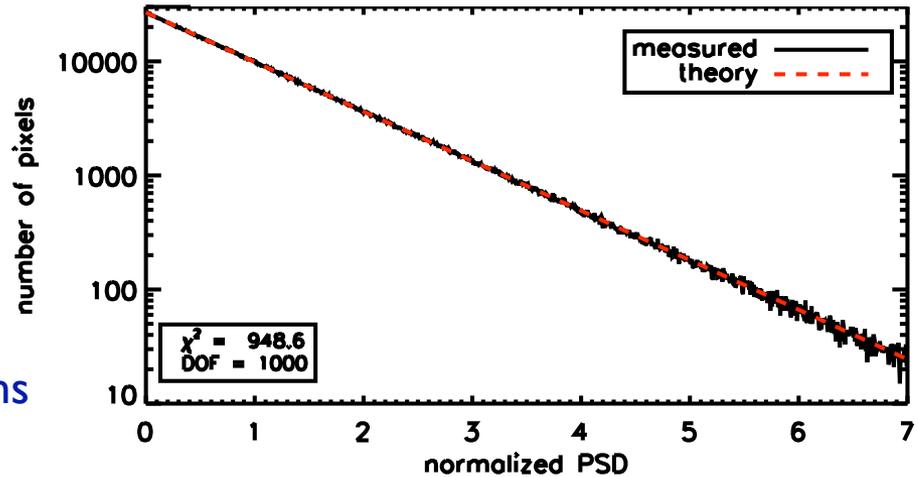
- Relative weights:

method	Fraction of obs	Fractional contribution to FOM
avg	50%	70%
planar	40%	29%
quadratic	10%	1%

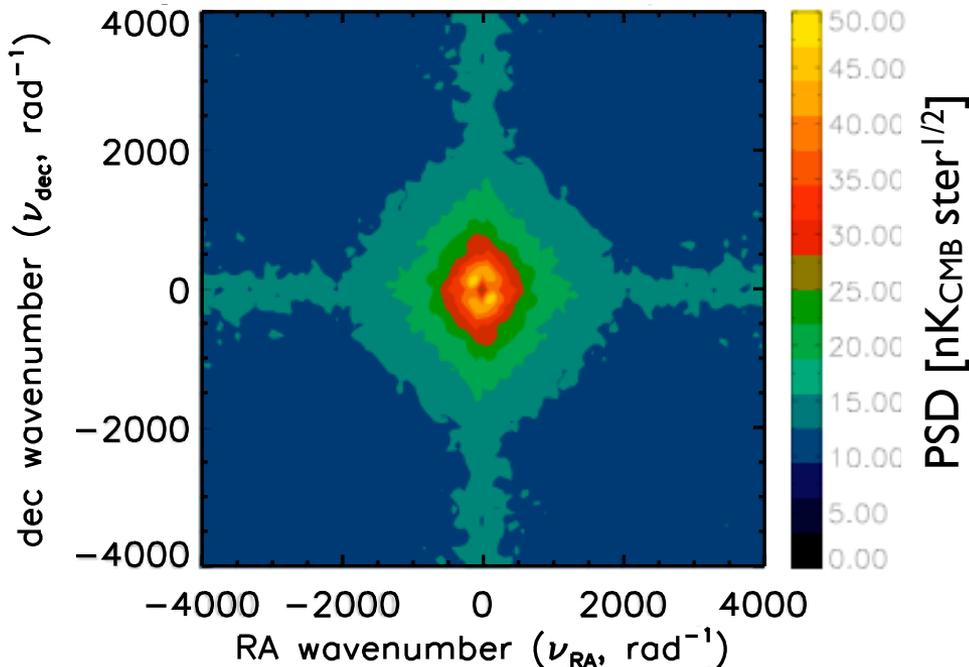
- Determine overall transfer function by weighted sum of single-obs transfer function

Coadd PSDs and Transfer Functions

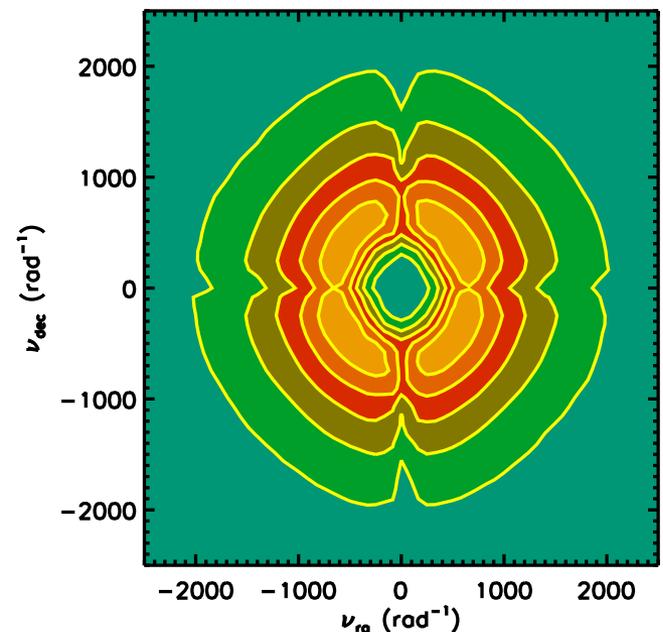
- Final PSD determined from jackknife coadds: random signs on each observation
- Final coadd PSD
 - not white
 - But beautifully gaussian: deviations present in single obs gone



Coadd PSD

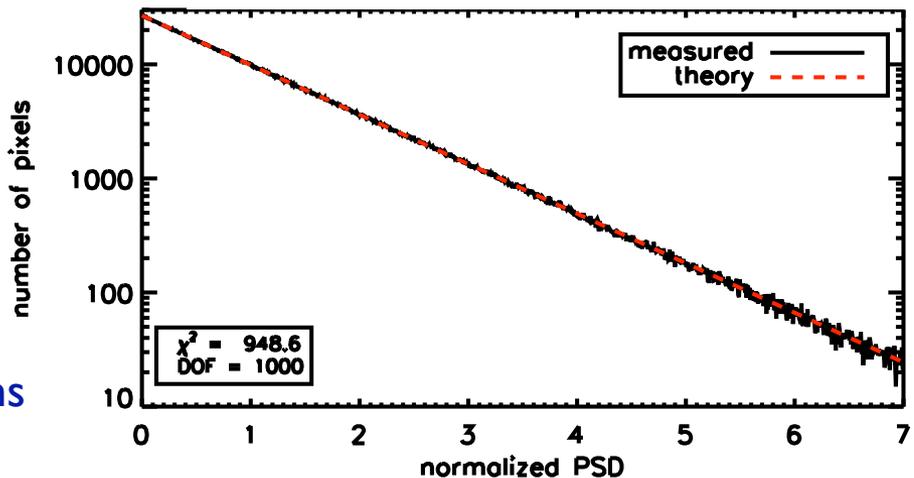


Coadd Transfer Function

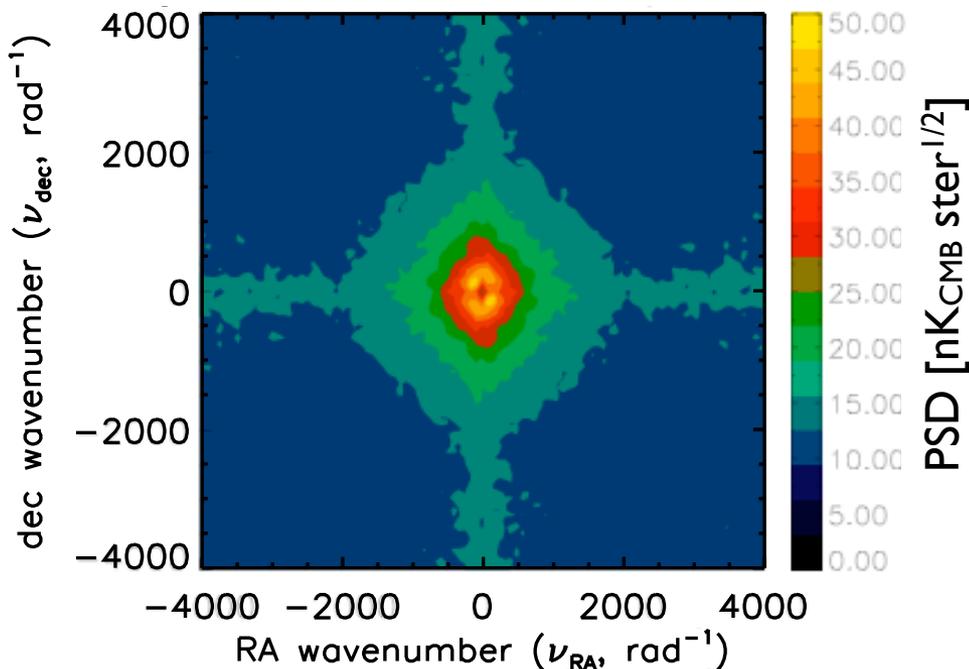


Coadd PSDs and Transfer Functions

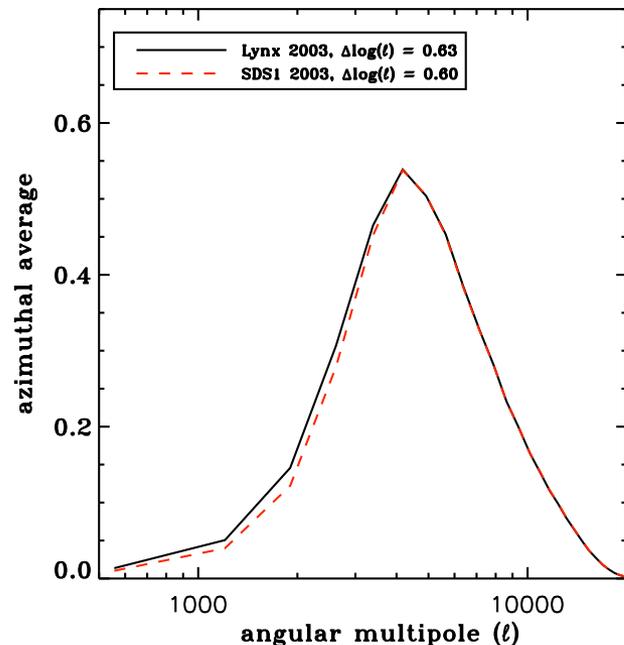
- Final PSD determined from jackknife coadds: random signs on each observation
- Final coadd PSD
 - not white
 - But beautifully gaussian: deviations present in single obs gone



Coadd PSD



Coadd Transfer Function



Total Anisotropy Power Spectrum Constraint

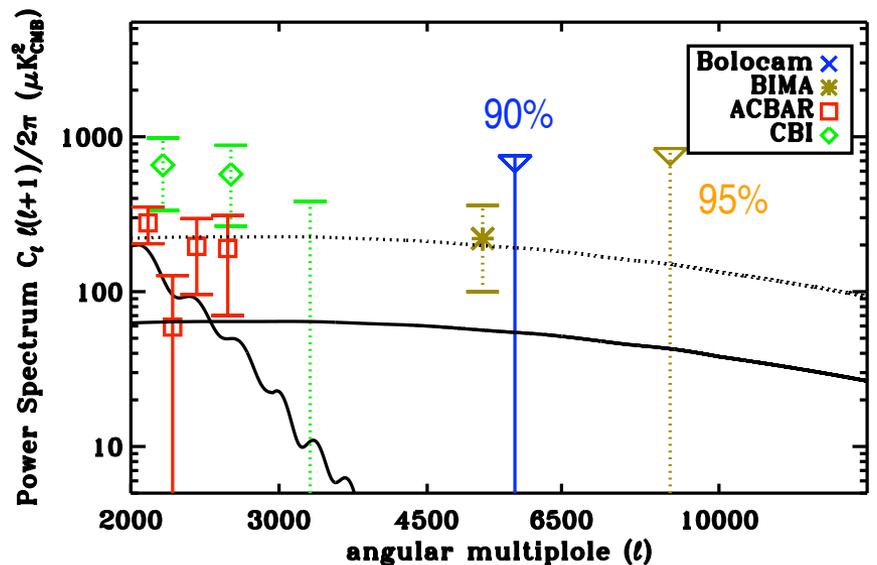
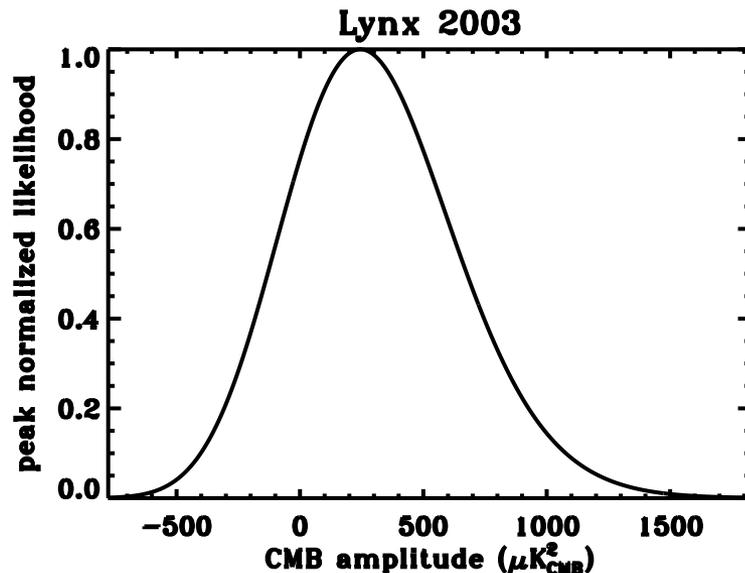
- Do we see excess noise power? Do Max L estimate of A, the amplitude of the flat bandpower anisotropy

$$\log(\mathcal{L}) = \sum_{\vec{\nu}} \left(-\log(\mathcal{P}_{\vec{\nu}} + AS_{\vec{\nu}}B_{\vec{\nu}}W_{\vec{\nu}}) - \frac{x_{\vec{\nu}}}{\mathcal{P}_{\vec{\nu}} + AS_{\vec{\nu}}B_{\vec{\nu}}W_{\vec{\nu}}} \right)$$

coadd PSD from jackknives
(noise estimate when no signal)

possible signal term

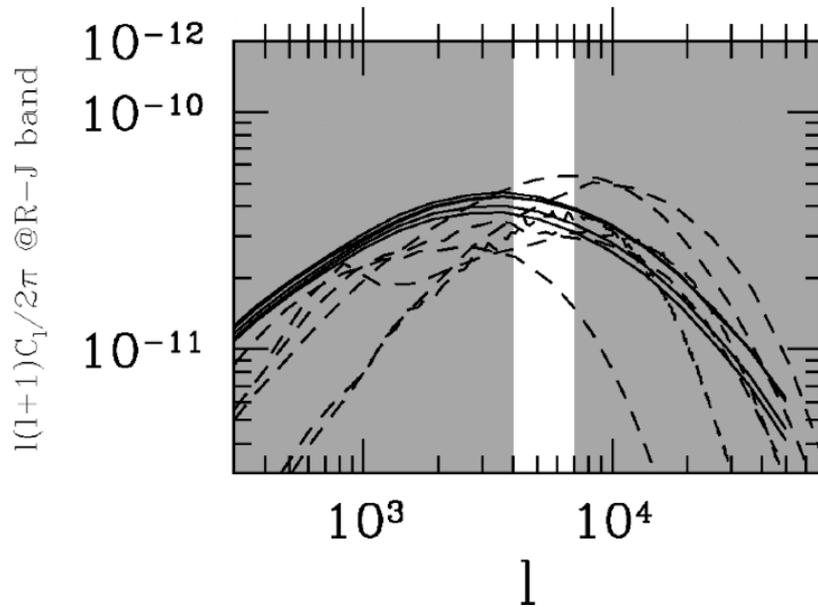
- Bayesian likelihood function width is v. approximate: correlations between Fourier modes puts in a covariance we have not included
- Use Feldman-Cousins to obtain correct frequentist confidence interval



(Kuo et al 2007 ACBAR points)

SZ Anisotropy Power Spectrum Constraint

- No detection of anisotropy power
- Want to constrain amplitude of putative SZ anisotropy power spectrum
- Complications:
 - Need to include contribution of CMB, $\sim 45 \mu\text{K}_{\text{CMB}}^2$
 - Done properly by adding expected value to P_ν
 - CMB power spectrum from Spergel et al (2007) and Kuo et al (2007)
 - Fluctuations automatically accounted for by adding random CMB realization to each jackknife noise realization.
 - What is the SZ power spectrum?
Use two models:
 - Flat bandpower like CMB
 - Komatsu and Seljak (2002) analytic spectrum
 - Other spectra in literature not very different in our ℓ range of interest



SZ Anisotropy Power Spectrum Constraint

- Constraints on amplitude of total and SZ anisotropy PS:

spectrum	flux uncertainty	68% CL interval	90% CL interval	95% CL interval	
flat-total	0	99 – 588 μK_{CMB}^2	0 – 755 μK_{CMB}^2	0 – 828 μK_{CMB}^2	
flat-SZE	0	90 – 582 μK_{CMB}^2	0 – 747 μK_{CMB}^2	0 – 830 μK_{CMB}^2	
flat-SZE	3.5% (meas)	89 – 634 μK_{CMB}^2	0 – 794 μK_{CMB}^2	0 – 876 μK_{CMB}^2	
flat-SZE	6.3% (total)	83 – 692 μK_{CMB}^2	0 – 956 μK_{CMB}^2	0 – 998 μK_{CMB}^2	
KS-SZE	0	77 – 543 μK_{CMB}^2	0 – 686 μK_{CMB}^2	0 – 766 μK_{CMB}^2	
KS-SZE	3.5% (meas)	76 – 569 μK_{CMB}^2	0 – 741 μK_{CMB}^2	0 – 834 μK_{CMB}^2	
KS-SZE	6.3% (total)	73 – 732 μK_{CMB}^2	0 – 950 μK_{CMB}^2	0 – 993 μK_{CMB}^2	

- 3 rows: no flux uncertainty, internal flux cal uncertainty, and full flux uncertainty (incl. uncertainty on external Mars model)
- SZ anisotropy scales as $\sigma_8^7(\Omega_b h)^2$
- Expected SZ anisotropy PS, using Dunkley et al (2008) cosmo params: $10 \mu\text{K}_{CMB}^2$
- Using K-S spectrum and $\Omega_b h$ from Dunkley et al (2008) and Kuo et al (2007), we set limit of $\sigma_8 < 1.55$ at 90% CL
 - $\sigma_8 = 0.80\text{-}0.85$ from primary PS + LSS, $\sigma_8 = 0.95$ from high- ℓ

What Happened?

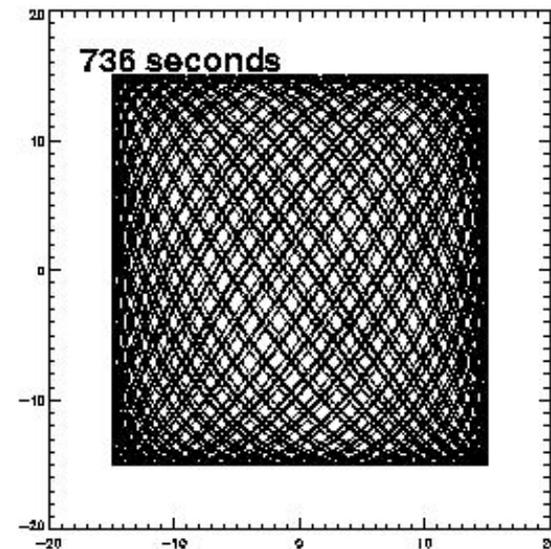
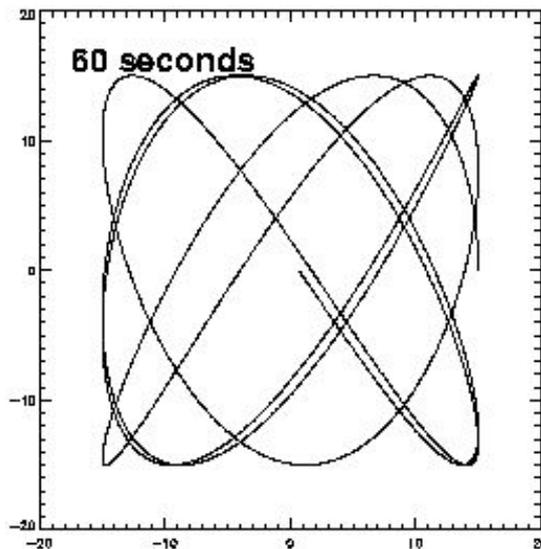
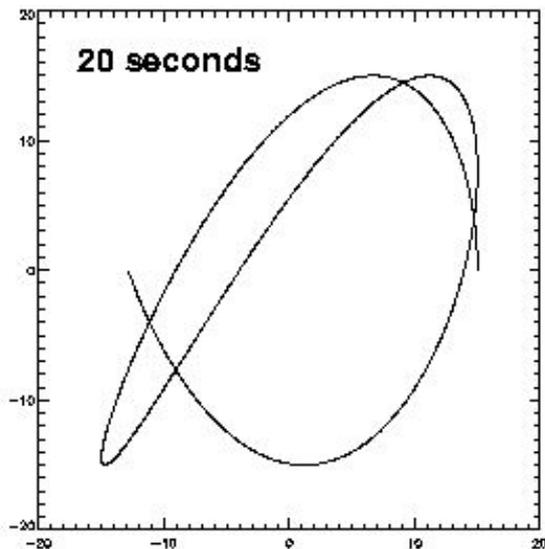
- Why did the survey end up being so unconstraining?
Sky noise, sky noise, sky noise
 - In hindsight, old SuZIE data is suggestive that spatial correlations are not simple enough to be fully removable
 - But no real measurements of atmospheric correlation function, not even from SCUBA
 - We have studied sky noise on Mauna Kea more exhaustively than anyone before (Sayers et al 2008, in prep)

Thoughts on SZ Surveys

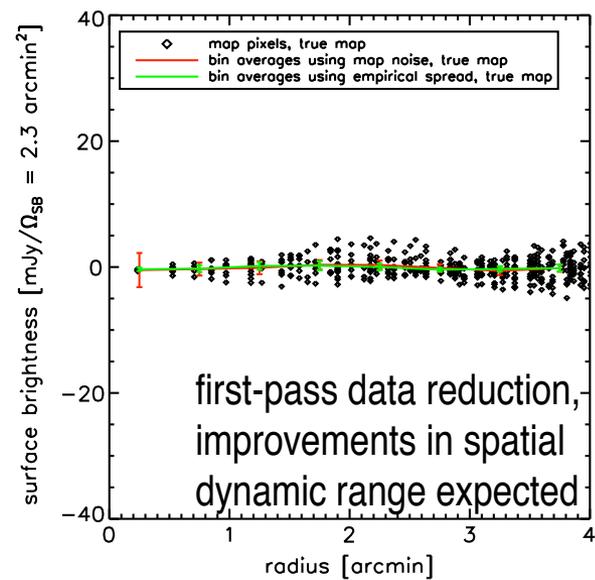
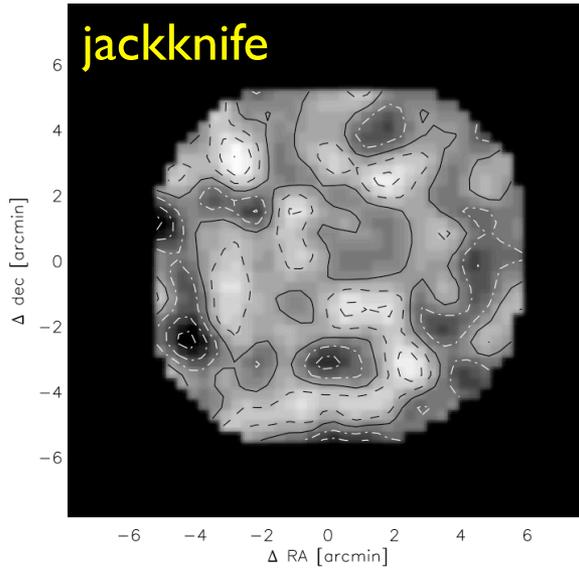
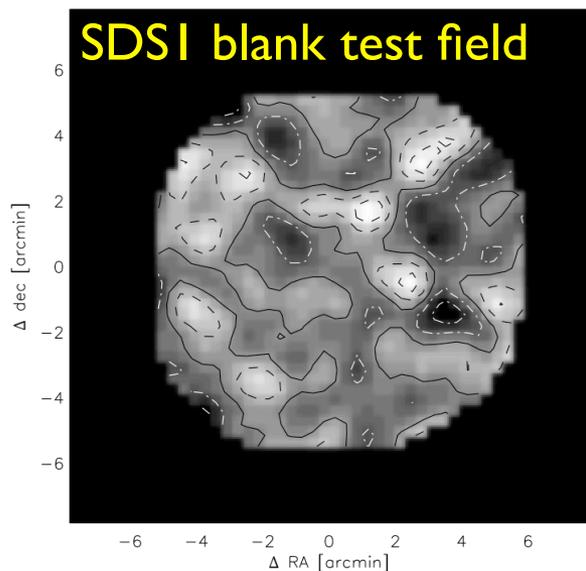
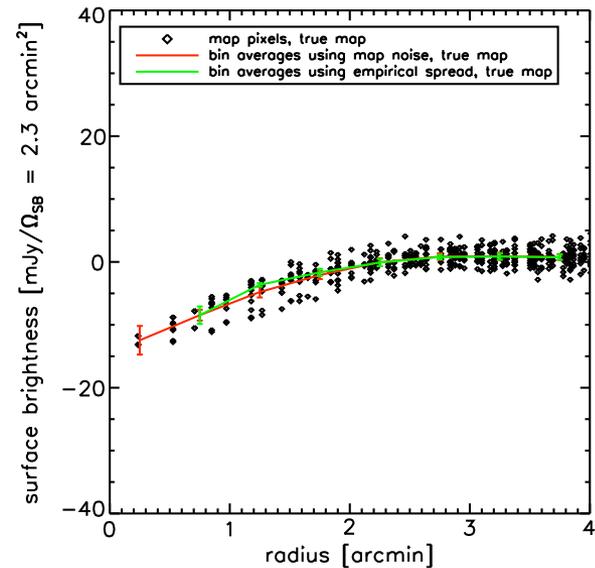
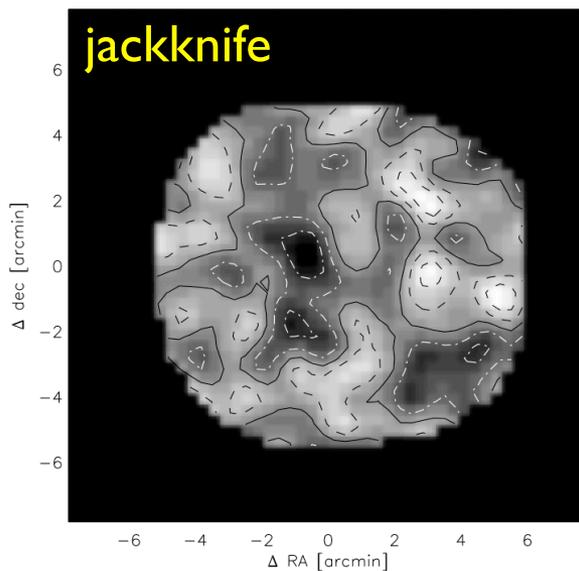
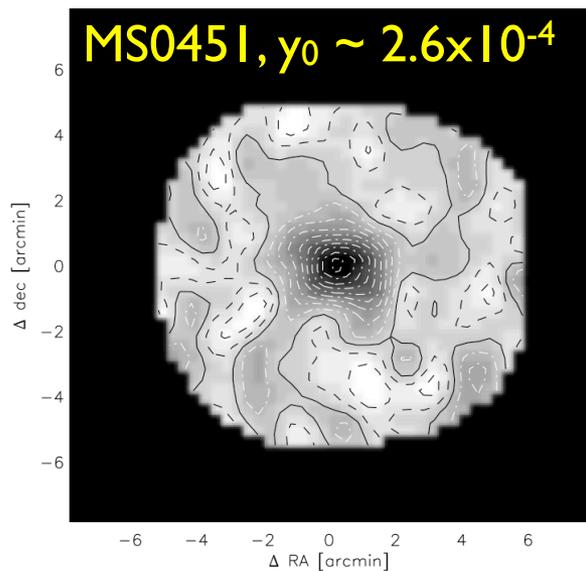
- On sky noise:
 - The sites for APEX, ACT, and SPT are better: Atacama and South Pole
 - But: even ACBAR saw sky noise at South Pole.
 - Possible sensitivity degradations
 - APEX and SPT: $2(f/\#)\lambda$ horns, so no Airy function coupling. Other degradations though:
 - imperfect correlation of atmosphere, leaving $1/f$ noise in timestreams
 - transfer function of sky noise removal will hurt sensitivity
 - ACT: $0.5(f/\#)\lambda$ bare absorber pixels; depending on how good or bad the atmosphere, may end up in same boat, with many fewer effective pixels, + above degradations

What Next?

- Bolocam
 - observing single massive clusters in raster mode has never been feasible because fields are too small: spend all the time turning around
 - we have learned how to observe single massive clusters in an efficient Lissajous scan mode, developed for SHARCII 350 μm CSO camera
 - would like to compare to OVRO/BIMA and SZA maps, resolve the discrepancies with SuZIE data



What Next?



What Next?

- MKID camera for CSO
 - New technologies enable 4-color camera with 8' FOV (750 μm - 1.3 mm, possibly extend to 2 mm)
 - *Spectral sky subtraction*: each spatial pixel observes in multiple colors, so atmosphere can be regressed out
 - SuZIE II showed that this works beautifully for 4 spatial pixels
 - But large simultaneous 4-color focal plane not feasible 7 years ago
 - No worries about spatial correlations, though need to be sure source is orthogonal to atmosphere (it is for SZ)
 - Massive SZ cluster observations in Lissajous mode
- LWCam for CCAT
 - New 25-m submm/mm telescope in Chile
 - 5 or 6-color 750 μm - 2 mm camera in planning
 - High-resolution multicolor followup of clusters discovered in large area SZ surveys, again using Lissajous mode
 - Reach SZ-confusion limit

Submm/mm MKID Camera

- Antenna coupled MKIDs, in-line bandpass filters to obtain four colors/ spatial pixel (220, 275, 350, 420 GHz)
- 8' FOV, 600 spatial pixel, on CSO 2010
- 16-pixel/2-color DemoCam fielded

