

Some general points:

§6.3

$$f_{lu} \equiv \frac{m_e c}{\pi e^2} \int \sigma_{lu}(\nu) d\nu$$

↑

oscillator strength (dimensional-less)

$$A_{ul} = \frac{8\pi^2 e^2 \nu_{lu}^2}{m_e c^3} \frac{\partial f_{lu}}{\partial \nu_{lu}}$$

$$= \frac{0.667 \text{ cm}^2 \text{ s}^{-1}}{\lambda_{lu}^2} \frac{\partial f_{lu}}{\partial \nu_{lu}}$$

unit of  $A_{ul} = \text{second}^{-1}$

ex:

$A_{21}$ (Ly $\alpha$ )	0.4162
$A_{31}$ (Ly $\beta$ )	0.079
$A_{41}$ (Ly $\delta$ )	0.029
$A_{32}$ (H $\alpha$ )	0.6408
$A_{42}$ (H $\beta$ )	0.1193

Lyman- $\alpha$ : large cross-section  
most number of atoms ( $H^I$ )

$$\tau_\nu = \frac{\pi e^2}{m_e c} f_{lu} \frac{N}{l} \Phi_\nu \left[ 1 - \frac{N_{ul}/g_u}{N_e/g_e} \right]$$

$f_{lu}$  ... oscillator strength

$\Phi_\nu$  ... velocity profile

Example: Lyman- $\alpha$

$\tau_0$  ... optical depth in the center of line

$$\tau_0 = 0.76 \left( \frac{N_H}{10^{13} \text{ cm}^{-2}} \right) \left( \frac{f_{lu}}{0.4164} \right) \left( \frac{\lambda_{lu}}{1216 \text{ \AA}} \right) \frac{10 \text{ km s}^{-1}}{b}$$

$$\tau_\nu = \tau_0 \exp\left(-\frac{u^2}{b^2}\right)$$

$$b = \sqrt{2} \sigma_\nu$$

↑ RMS velocity

$$\tau_0 = 0.76 \left( \frac{N_H}{10^{13} \text{ cm}^{-2}} \right) \left( \frac{f_{lu}}{0.4164} \right) \left( \frac{\lambda_{lu}}{1216 \text{ \AA}} \right) \frac{10 \text{ km/s}}{b}$$

$$\text{Ly}\alpha: A = 5 \times 10^8 \text{ s}^{-1} \quad \therefore \Delta\nu = 5 \times 10^8 \text{ Hz}$$

$$\therefore \frac{\Delta\nu}{\nu_0} \approx \frac{5 \times 10^8 \text{ Hz}}{c/1216 \text{ \AA}}$$

$$\Rightarrow \sigma_v (\text{intrinsic}) \approx 0.06 \text{ km/s}$$

$$\tau_0 \approx 1.2 \left( \frac{N_H}{10^{11} \text{ cm}^{-2}} \right)$$

$$\text{Say } \sigma = \lambda_{lu}^2 \approx 1.5 \times 10^{-10} \text{ cm}^2$$

Allowed (strong) lines have  ~~$\tau_0$~~

$$\sigma \approx \lambda^2 \text{ at line-center.}$$

Photo-electric absorption:

Hydrogen:

$$\sigma_{pi} \approx \sigma_0 \left( \frac{h\nu}{Z^2 I_H} \right)^{-3}$$

$Z$ ... atomic charge of nucleus

$I_H$ ... ionization potential of Hydrogen

$$\sigma_0 \approx 6.3 \times 10^{-18} \text{ cm}^2$$

Compare to Thompson cross-section:

"Lyman-thick"

$$\Rightarrow N_H \gtrsim 10^{18} \text{ cm}^{-2}$$

If a region is Lyman thick then

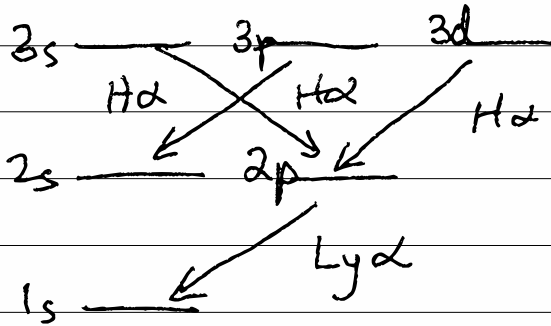
$\text{Ly}\alpha$ ,  $\text{Ly}\beta$ , ... will rattle around!

...destruction by dust!

Q:

# Hydrogens:

§14



Case A: Ly $\alpha$  escapes from regions  
ex. fast shocks, low column density  
high temperature

Recombinations to all ~~states~~ levels

$$\alpha_A = \sum_{n=1}^{\infty} \alpha_{ne}(T)$$

Recomb per unit volume per unit second  
 $\alpha_A n_e n_p$

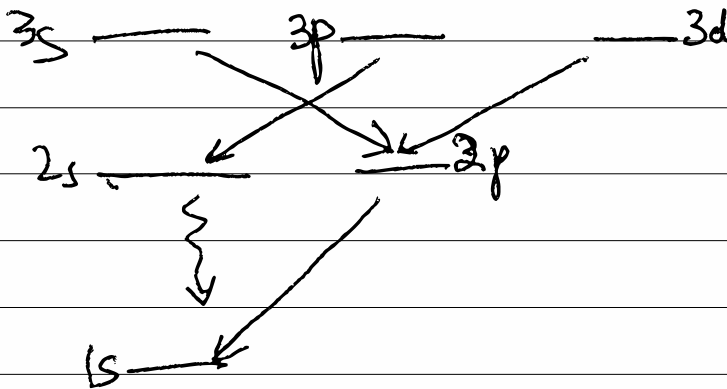
case B.

Ly $\alpha$  rattles around

"on the spot absorption"

$$\alpha_B = \sum_{n=2}^{\infty} \alpha_{ne}(T)$$

Two-Photon Decay



$$A(2s \rightarrow 1s) \approx 8 \text{ s}^{-1}$$

So can have collisions and push H into 2p state.

$$n_{e, \text{critical}} = \frac{A_{2s \rightarrow 1s}}{q_p 2s \rightarrow 2p + q_e 2s \rightarrow 2p}$$

$q_p 2s \rightarrow 2p$  -- proton collision

$q_e 2s \rightarrow 2p$ : electron collision

$$j_\nu = \frac{n_e n_p \alpha_{2s}}{\left(1 + \frac{n_e}{n_{\text{crit}}}\right)} \cdot h\nu \cdot \frac{1}{4\pi} P_\nu$$

↑  
two-photon profile

## Helium:



This is hydrogenic

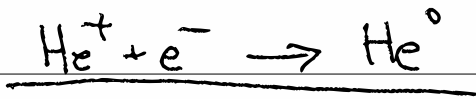
$$\text{"Ly}\alpha\text{"} = \frac{1216 \text{ \AA}}{4}$$

$$\text{Lyman edge} = \frac{912 \text{ \AA}}{4}$$

- Helium Lyman series can ionize H
- However photoelectric cross-section is much smaller than resonance line transitions
- So you see line spectrum but ~~4~~ times higher in frequency relative to H
- $2s \rightarrow 1s$  ~~is~~  $A = 527 \text{ s}^{-1}$  ( $Z^6!$ )  
 $\Rightarrow$  ionize H atoms

Locations: Planetary nebula, Wolf Rayet stars





$\alpha_{1s^2}(\text{He})$  ... recombination ~~to~~ directly to ground state

This will generate  $> 24.6 \text{ eV}$  which can ionize either  $\text{HI}$  or  $\text{HeI}$ .

$\alpha_B(\text{He})$  ... recombination to any state but  $1s^2$

$$\alpha_{1s^2} \approx 1.54 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1} \quad T = 10^4 \text{ K}$$

$$\alpha_B \approx 2.7 \times 10^{-13} \text{ cm}^3 \text{ s}^{-1} \quad T = 10^4 \text{ K}$$

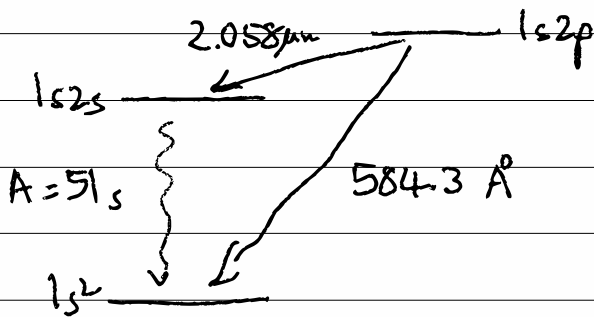
If case B holds then  $\#$  recombinations to ground state can ionize  $\text{HeI}$

$$\alpha_{\text{eff}}(\text{He}) = \alpha_B(\text{He}) + y \alpha_{1s^2}(\text{He})$$

↑  
fraction of  $\text{H}$  ionization

Recombination to levels other than  $1s^2$

25% of recombination is to singlet  
75% of recombiners is to triplet

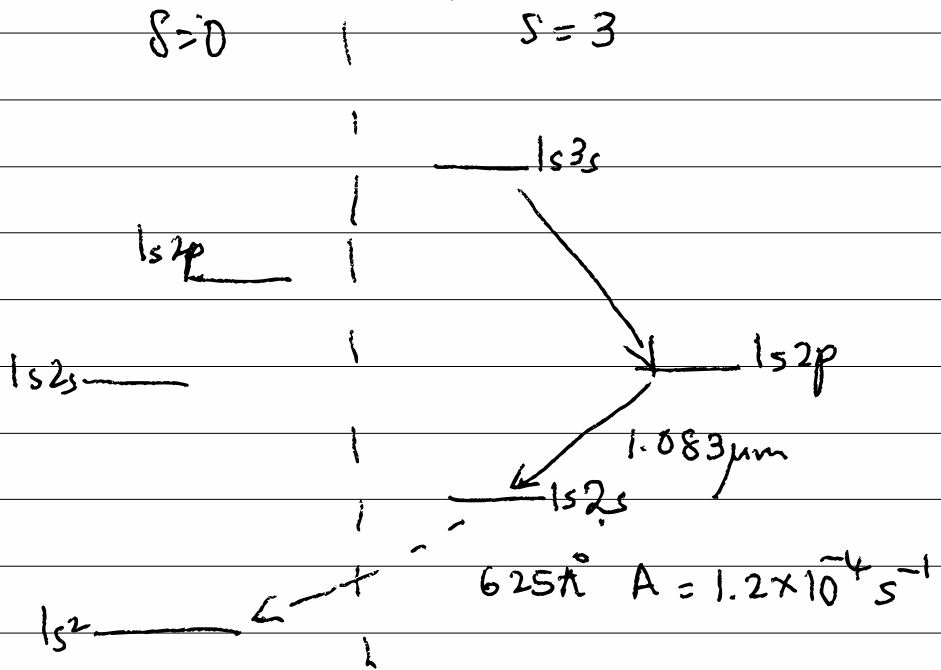


$$A(1s2p \rightarrow 1s2s) = 1.97 \times 10^6 \text{ s}^{-1}$$

$$A(1s2p \rightarrow 1s^2) = 1.8 \times 10^9 \text{ s}^{-1}$$

The  $\lambda 584.3\text{ \AA}$  is absorbed, being resonant, rapidly until it dies by hydrogen ionization, or destruction by dust, or by two photon decay

Recombination to triplet series:



The  $1s2s$   $S=3$  state is meta-stable.

The long duration allows collisional excitations to  
 $1s2p$  (triplet),  $1s3s$  (triplet)  
 $1s2p$  (singlet),  $1s2s$  (singlet),  $1s2p$  (singlet)

He I  $1.0833 \mu\text{m} \Rightarrow$  diagnostic of temperature  
 & density