### Ay 127, Spring 2013: Cosmological Tests and

the Contents of the Universe

- Tests for the expansion of the universe
- Classical and modern cosmological tests of global geometry and dynamics
  - Except for the CMB fluc's, to be covered in more detail later
- The "concordance cosmology"
- The nature of different density components (and how do we measure them)
  - Except that we'll leave a more detailed discussion of the nature of the non-baryonic DM and DE for another lecture

#### Tests for the Expansion of the Universe

- Tolman surface brightness (SB) test
  - In a stationary, Euclidean universe SB = const.
  - In an expanding universe,  $SB \sim (1+z)^{-4}$
  - In a "tired light" model,  $SB \sim (1+z)^{-1}$
- Time dilation of Supernova light curves
  - Time stretches by a factor of (1+v/c) = (1+z)
- The match between the energy density and  $T^4$  for the blackbody and the CMBR
  - For a blackbody, energy density  $u \sim T^4$
  - In an expanding universe, for photons, energy density is  $u \sim (1+z)^4$ , and since  $T \sim 1/\lambda \sim (1+z)$ ,  $u \sim T^4$

#### The Tolman Test

Surface brightness is flux per unit solid angle:  $B = \frac{f}{d\omega}$ 

This is the same as the luminosity per unit area, at some distance D. In cosmology,

$$B = \frac{L}{D_L^2} \frac{D_A^2}{dl^2}$$

In a stationary, Euclidean case,  $D = D_L = D_A$ , so the distances cancel, and SB = const. But in an expanding universe,  $D_L = D(1+z)$ , and  $D_A = D/(1+z)$ , so:

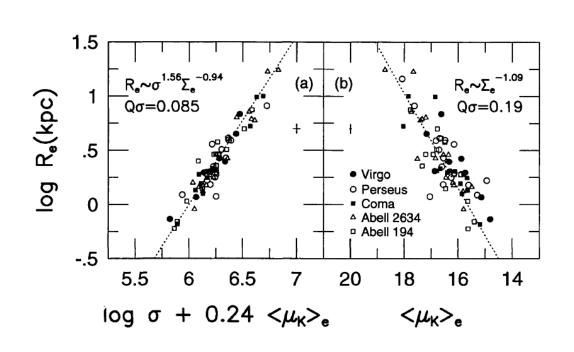
 $B = \frac{L}{dl^2} \frac{D_A^2}{D_A^2} = \frac{L}{dl^2} (1+z)^{-4}$ 

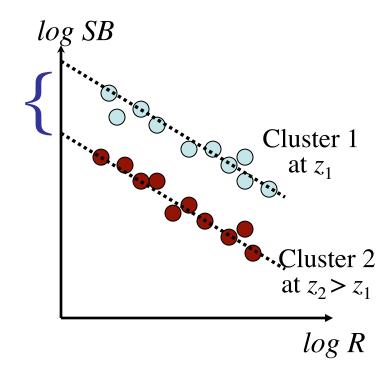
Note that this is independent of cosmology!

#### **Performing the The Tolman Test**

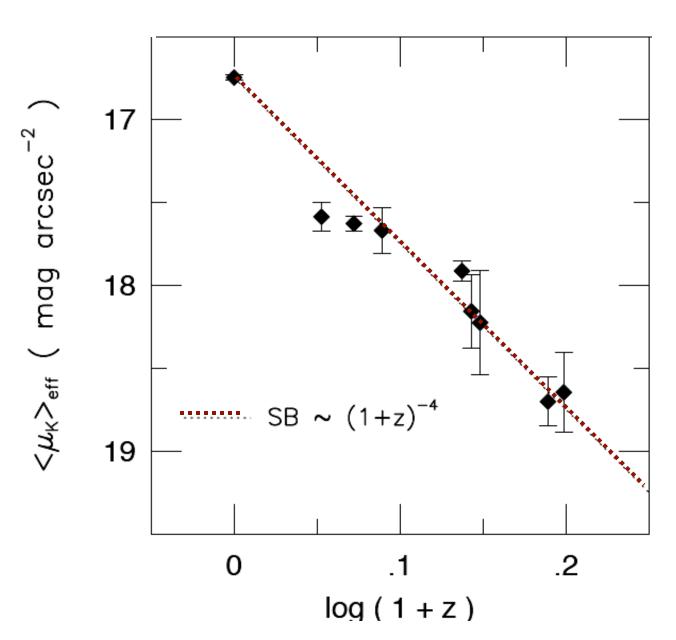
We need a standard (constant) unit of surface brightness = luminosity/area, to observe at a range of redshifts (a "standard fuzz"?)

A good choice is the intercept of surface brightness scaling relations for elliptical galaxies in clusters





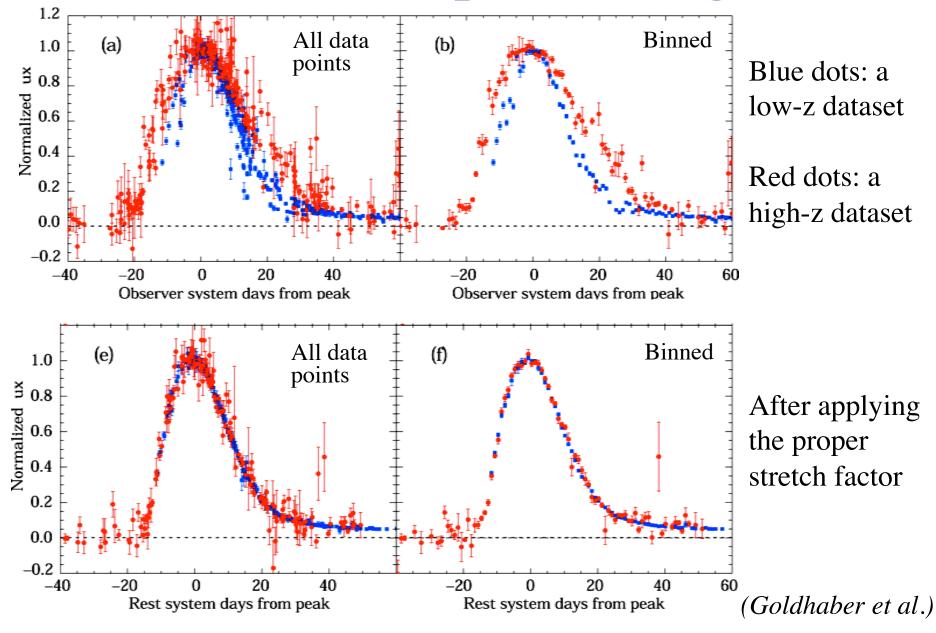
#### The Tolman Test Results



Surface brightness intercept of the Fundamental Plane correlation, for elliptical galaxies in clusters out to z ~ 0.6. It assumes a reasonable galaxy evolution model correction.

(from Pahre et al.)

#### Time Dilation of Supernova Lightcurves

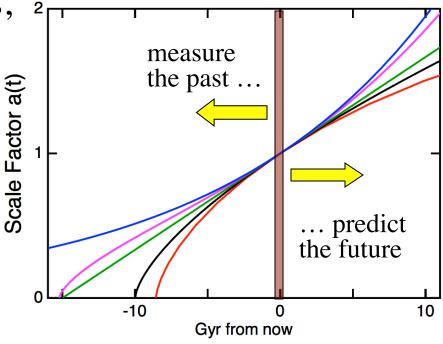


#### Cosmological Tests: The Why and How

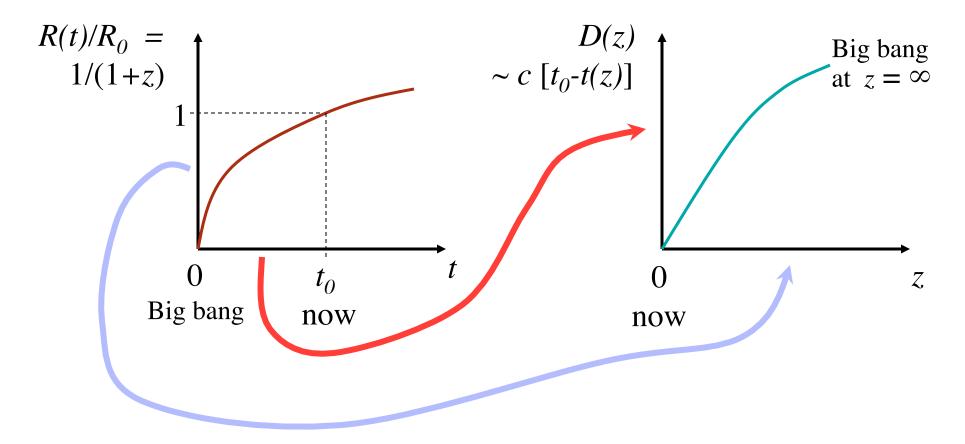
- The goal is to determine the global geometry and the dynamics of the universe, and its ultimate fate
- The basic method is to somehow map the history of the expansion, and compare it with model predictions
- A model (or a family of models) is assumed, e.g., the Friedmann-Lemaitre models, 2

typically defined by a set of parameters, e.g.,  $H_0$ ,  $\Omega_{0,m}$ ,  $\Omega_{0,\Lambda}$ ,  $q_0$ , etc.

 Model equations are integrated, and compared with the observations

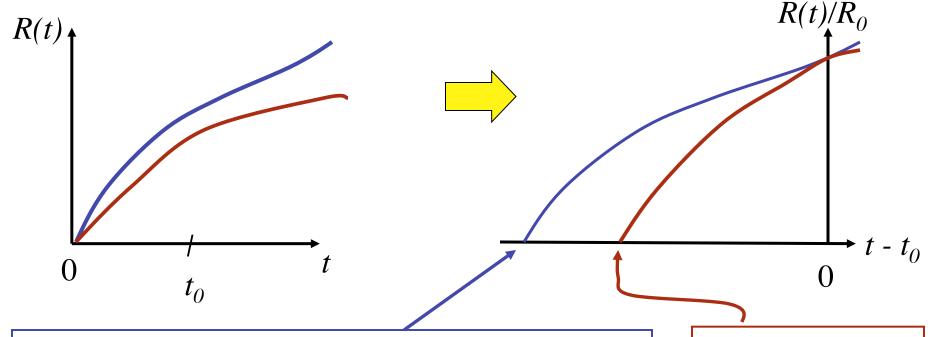


#### The Basis of Cosmological Tests



All cosmological tests essentially consist of comparing some measure of (relative) distance (or look-back time) to redshift. Absolute distance scaling is given by the  $H_0$ .

## Cosmological Tests: Expected Generic Behavior of Various Models



Models with a lower density and/or positive  $\Lambda$  expand faster, are thus larger, older today, have more volume and thus higher source counts, at a given z sources are further away and thus appear fainter and smaller

Models with a higher density and lower  $\Lambda$  behave exactly the opposite

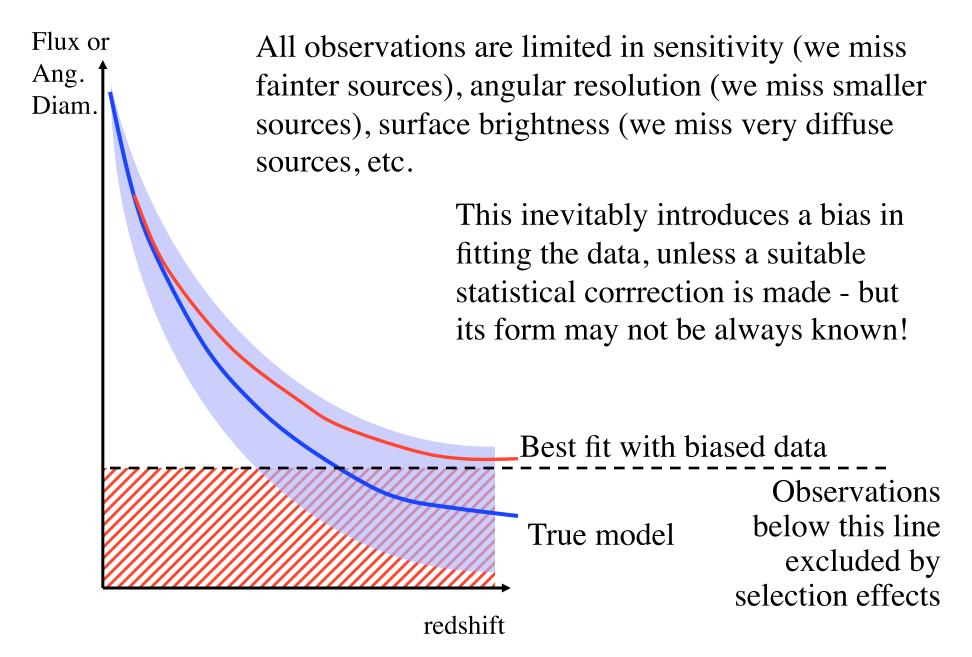
### The Types of Cosmological Tests

- The Hubble diagram: flux (or magnitude) as a proxy for the luminosity distance, vs. redshift requires "standard candles"
- Angular diameter as a proxy for the angular distance, vs. redshift requires "standard rulers"
- Source counts as a function of redshift or flux (or magnitude), probing the evolution of a volume element requires a population of sources with a constant comoving density "standard populations"
- Indirect tests of age vs. redshift, usually highly model-dependent "standard clocks"
- Local dynamical measurements of the mass density,  $\Omega_{m0}$
- If you measure  $H_0$  and  $t_0$  independently, you can constrain a combination of  $\Omega_{m0}$  and  $\Omega_{\Lambda}$

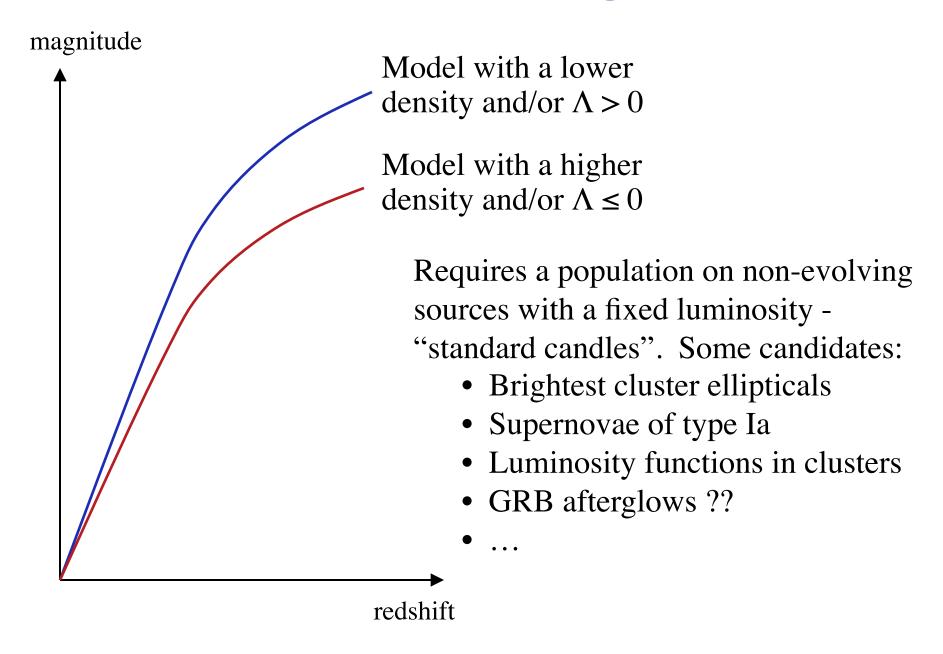
#### Cosmological Tests: A Brief History

- A program of "classical" cosmological tests (Hubble diagram, angular diameter test, source counts) was initiated by Hubble, and carried out at Palomar and elsewhere by Sandage and others, from 1950s through 1970s
- Galaxies, clusters of galaxies, and radio sources were used as standard candles, rulers, or populations. Unfortunately, all are subject to strong and poorly constrained *evolutionary effects*, which tend to dominate over the cosmology this foiled most of the attempted tests, and became obvious by 1980's
- In the late 1990's, Supernova Ia Hubble diagram, and especially measurements of CMBR fluctuations power spectra (essentially an angular diameter test) completely redefined the subject
- The cosmological parameters are now known with a remarkable precision a few percent; this is the era of "*precision cosmology*"

#### **Selection Effects and Biases**



#### The Hubble Diagram

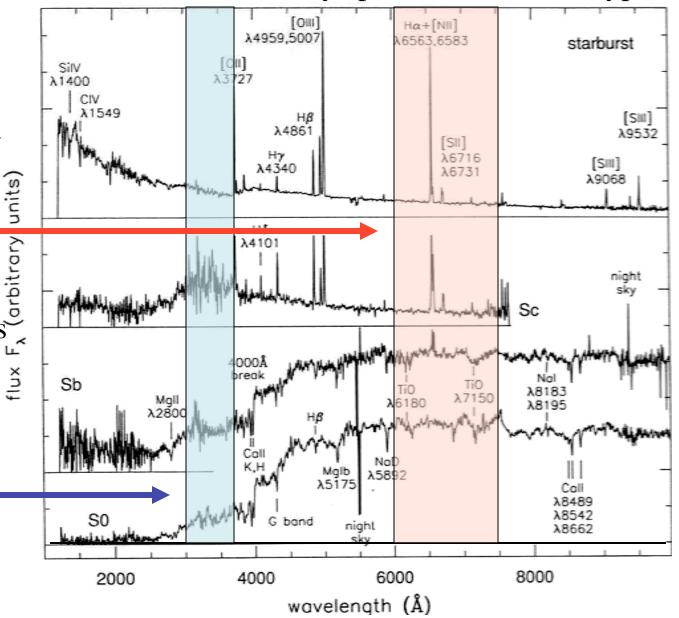


#### The K-Correction

#### Galaxy spectra of different types

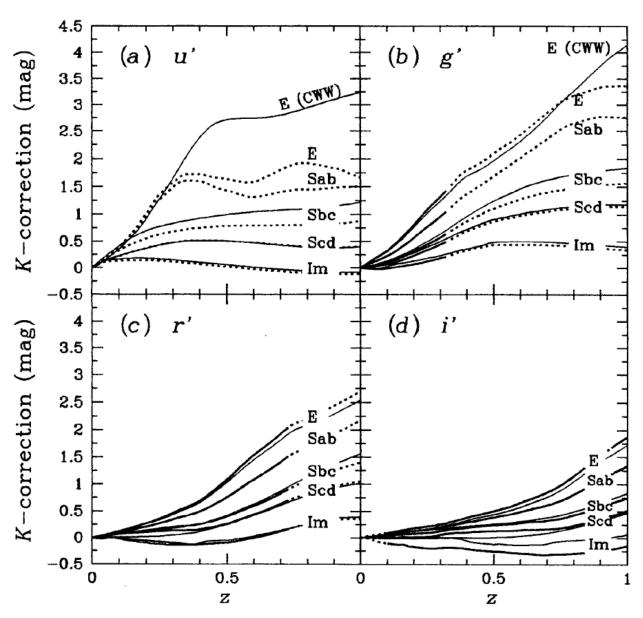
Photometric measurements are always obtained in some bandpass fixed in the observer's frame, e.g., the *U,B,V,R*...

But in a redshifted galaxy, this bandpass now samples some other (bluer in the galaxy's restframe) region of the spectrum, and it is also (1+z) times narrower



#### The K-Correction

Thus, we integrate the spectrum over the bandpass in the observed



frame, and in the galaxy's restframe, take a ratio, express it in magnitudes, and that is the

#### **K-correction**

It has to be done for all different types of galaxy spectra, as it depends on the star formation rates, and it varies with bandpass

(Fukugita *et al*.)

#### The Hubble Diagram: Early Work

• Mostly done at Palomar by Sandage and collaborators, and by Gunn and collaborators, using brightest cluster ellipticals, with corrections for cluster richness etc.

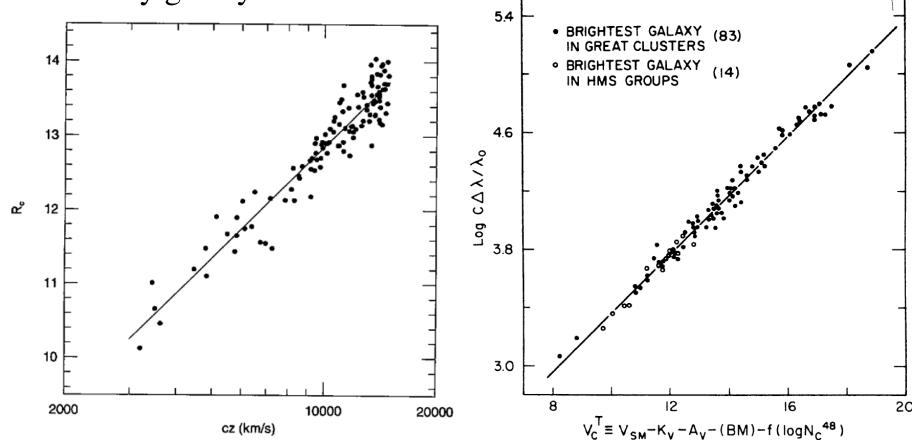
PAST PROPER TIME SAMPLED

• Foiled by galaxy evolution!

PAST PROPER TIME SAMPLED

(UNIT- I HUBBLE TIME)

O.OI O.OZ O.OS O.IO O.ZO O.3O



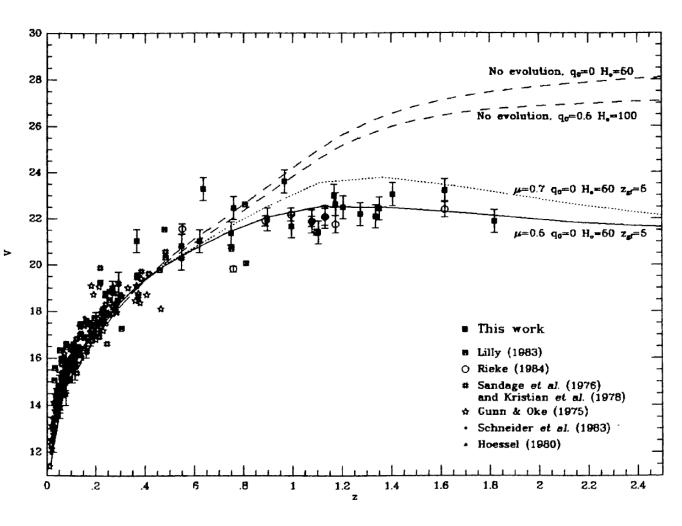
#### **Effects of Galaxy Evolution**

 Alas, galaxies were generally brighter in the past, since there was more star formation, and young, luminous, massive stars have

short lifetimes

• This tends to overwhelm the cosmological effects, especially in the bluer bands

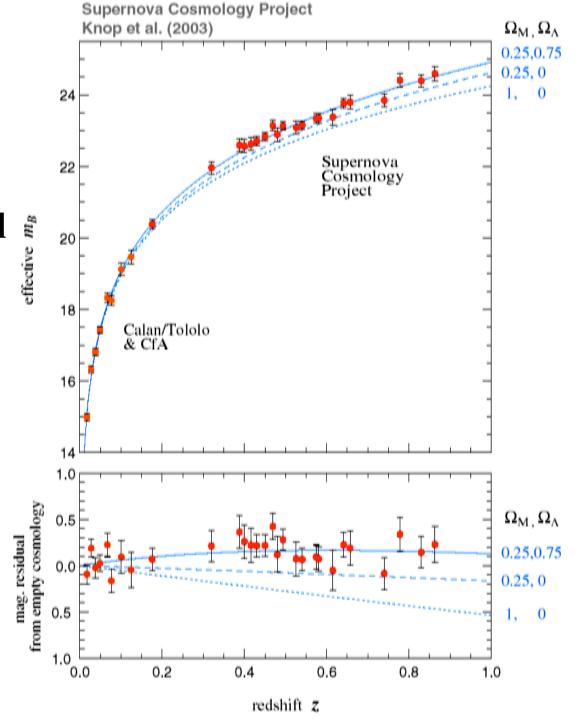
The Hubble diagram → for powerful radio galaxies (Djorgovski et al.)



#### The Supernova Ia Hubble Diagram

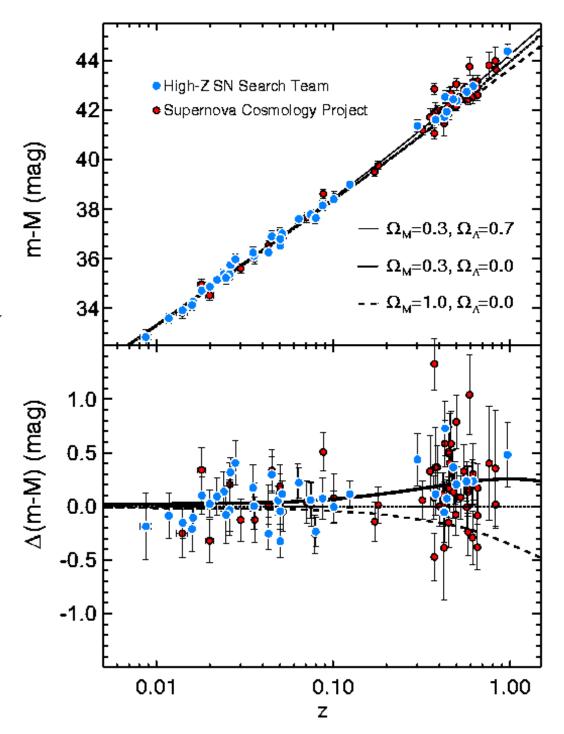
- The field was reborn with the advent of the SN Ia Hubble diagram, following the standardization of their peak brightness using light curve shapes
- There are still some unknowns:
  - Explosions not fully understood; many possible models:
     Chandrasekhar-mass models, deflagrations vs. detonations
  - Progenitor systems not known: white dwarfs yes, but double degenerate vs. single degenerate binaries ...
- SN Ia are not really standard candles ...
  - There are large variations in light curve shapes, colors, spectral evolution, and some clear outliers; possible differences in physical parameters, e.g, Ni mass
- But they *are* good distance indicators, after the empirical correction for light curve shapes
- Do they evolve (e.g., due to metallicity)? Maybe...

This yielded the evidence for an accelerating universe and the positive cosmological constant, independently and simultaneously by two groups: The Supernova Cosmology Project at LBL (Perlmutter et al.), and ...

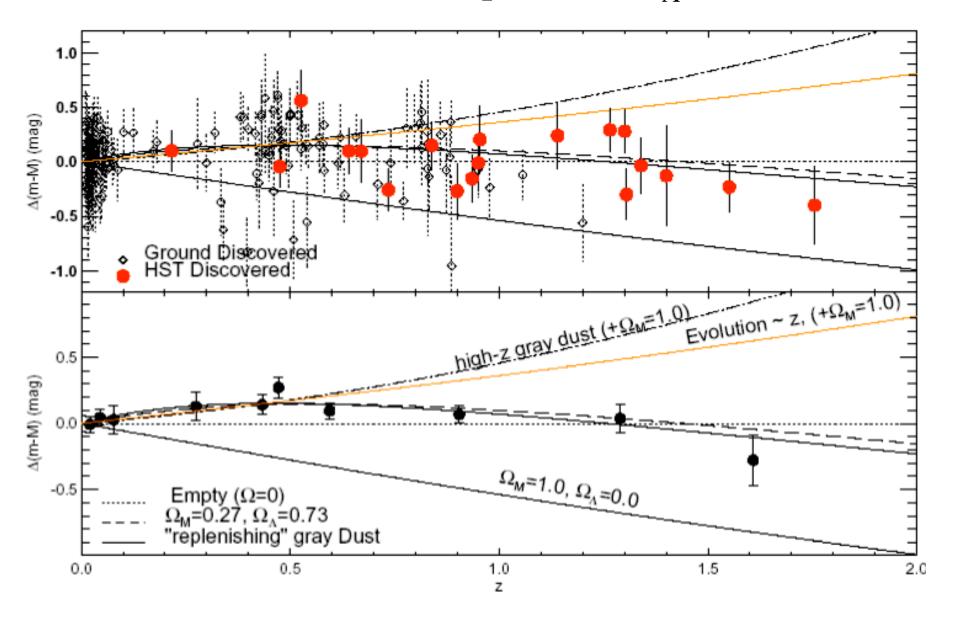


... and by the High-Z
Supernova Team
(B. Schmidt, A. Riess, et al.)

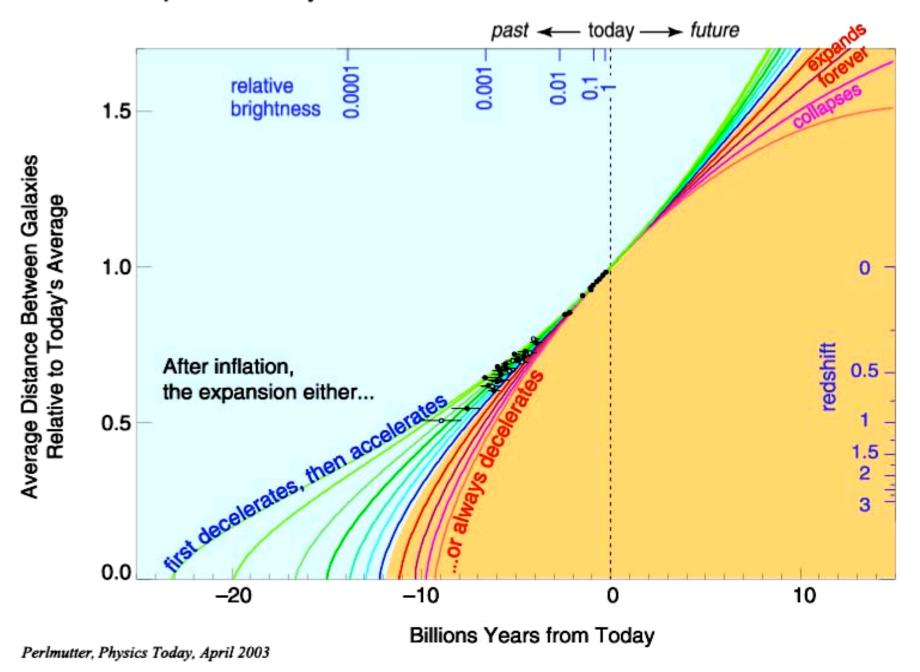
Both teams found very similar results ...



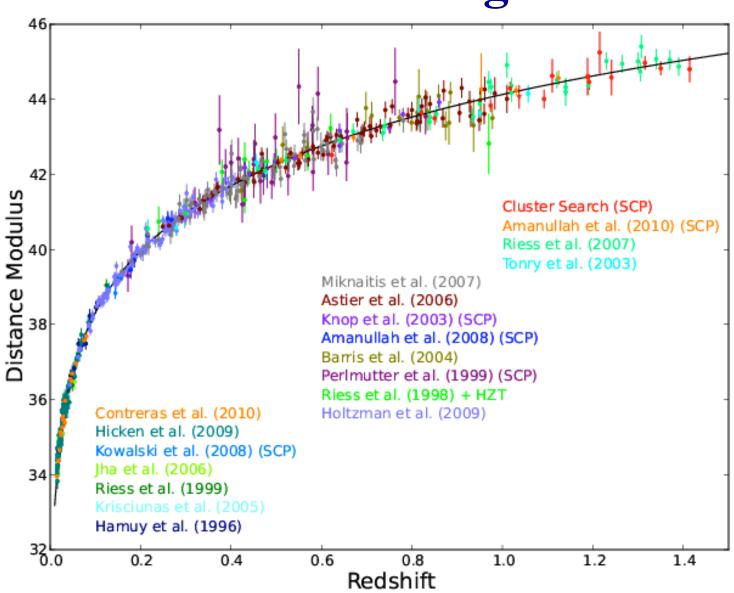
#### Current evidence points to $\Omega_{\Lambda} \sim 0.7$

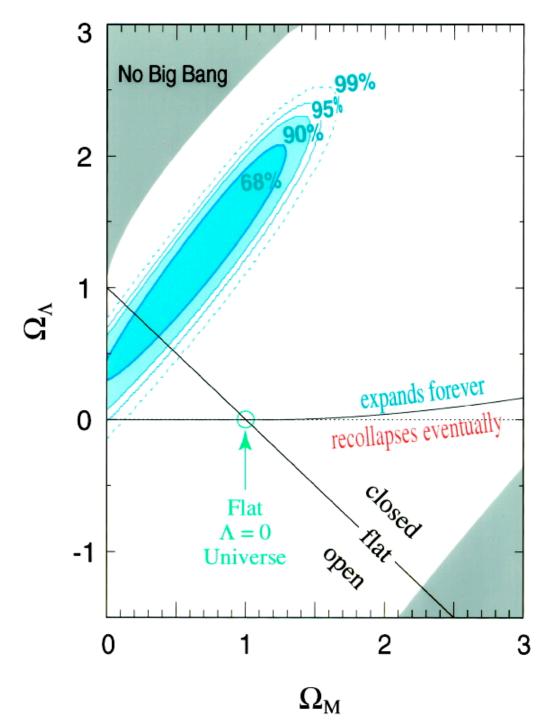


#### **Expansion History of the Universe**



# A Modern Version of the SN Hubble Diagram



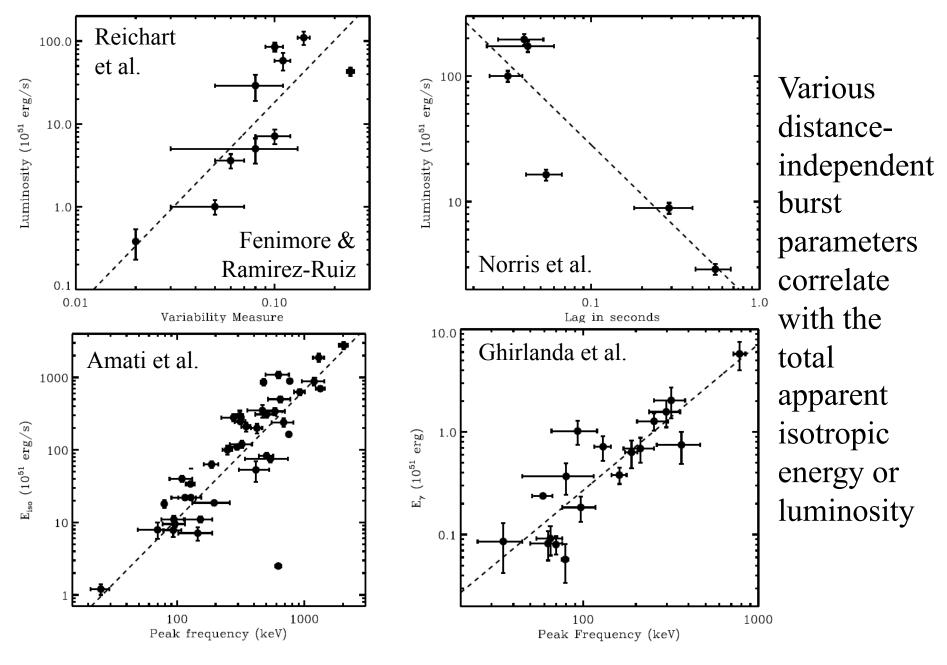


SN measurements on their own actually define an allowed region in the plane of  $[\Omega_m, \Omega_{\Lambda}]$ 

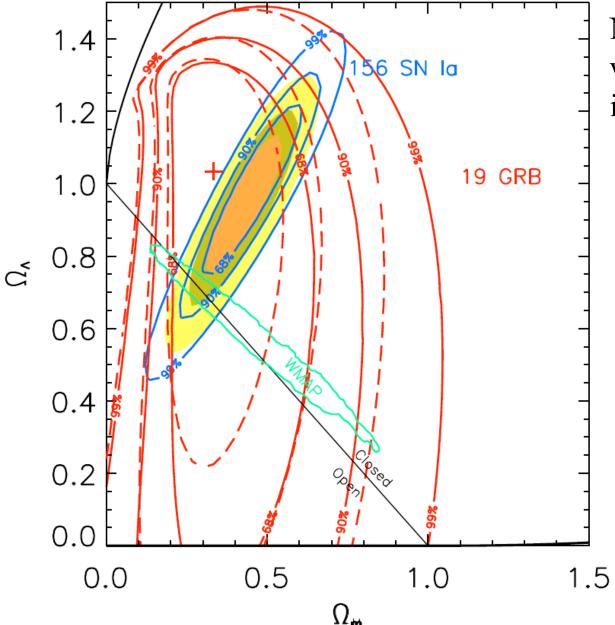
Example of **degeneracy**: distinct universes produce identical results for this cosmological test

We need some additional, constraints (e.g., flatness) to pin down the actual value of  $\Omega_{\Lambda}$ 

#### GRBs as Standard Candles?



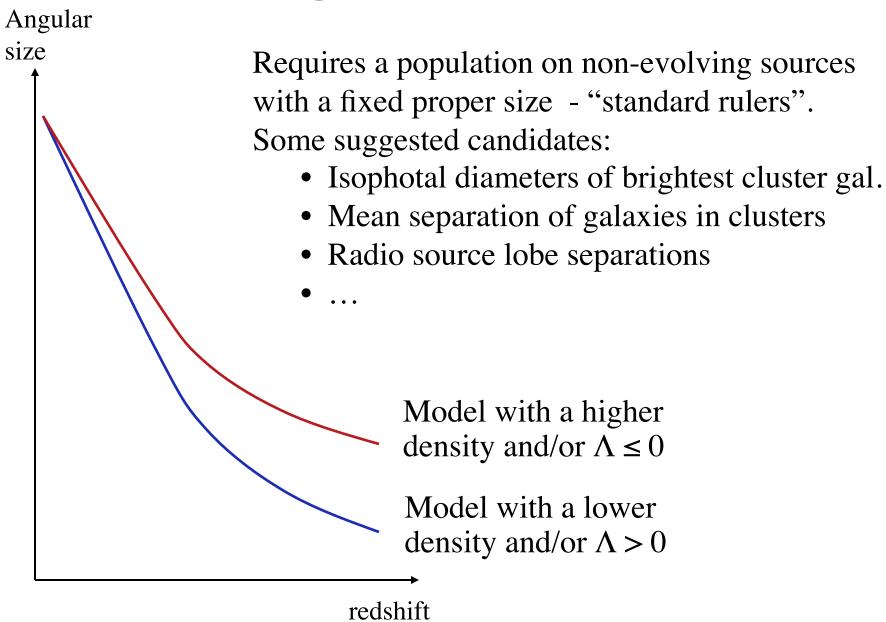
#### GRBs as Standard Candles?



Not quite competitive with SNe yet, but there is a promise...

(Figure from Lazzati et al.)

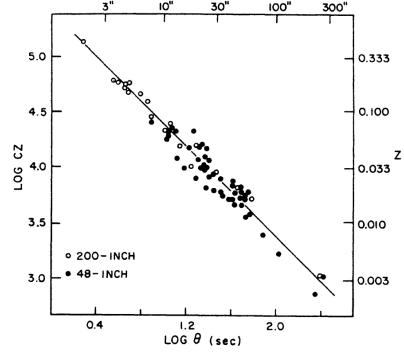
#### The Angular Diameter Test

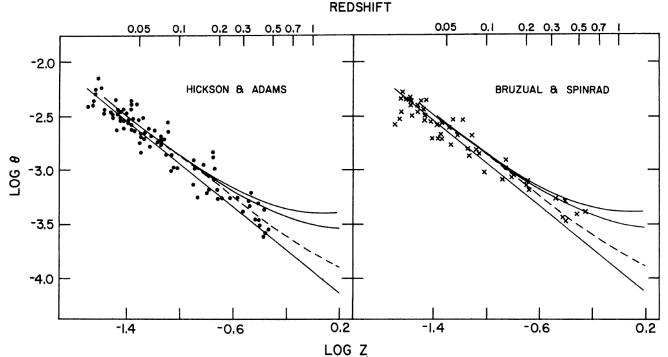


#### The Angular Diameter Test: Some Early Examples

Brightest cluster ellipticals →

Clusters of galaxies

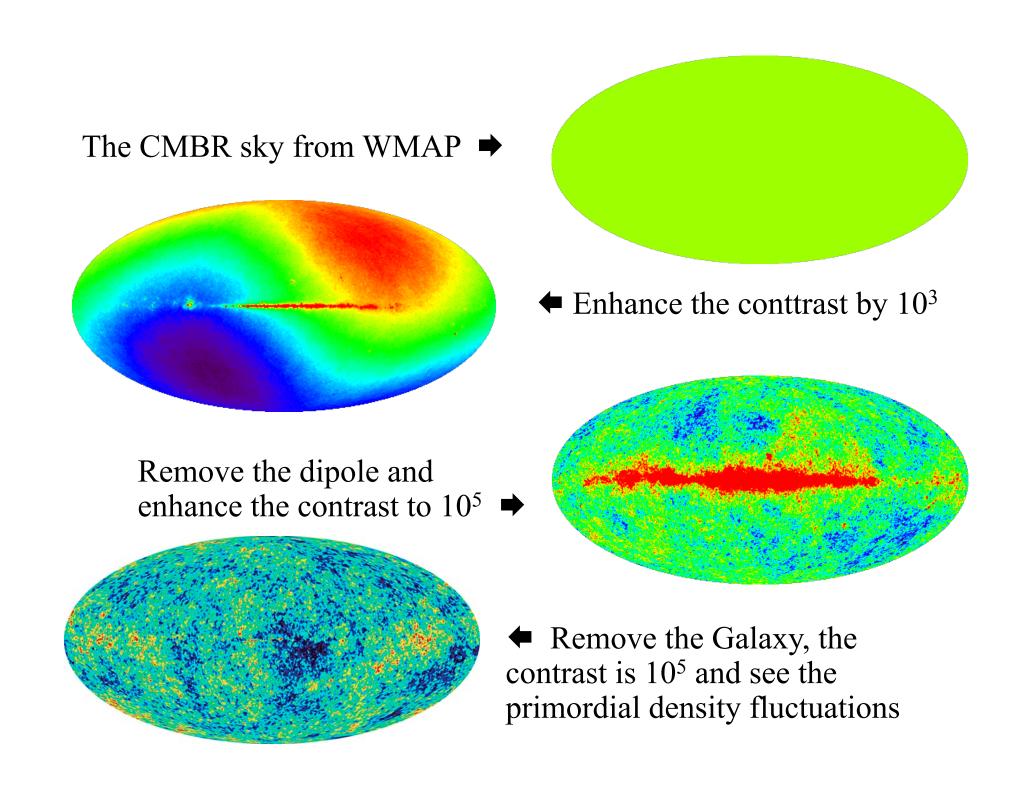




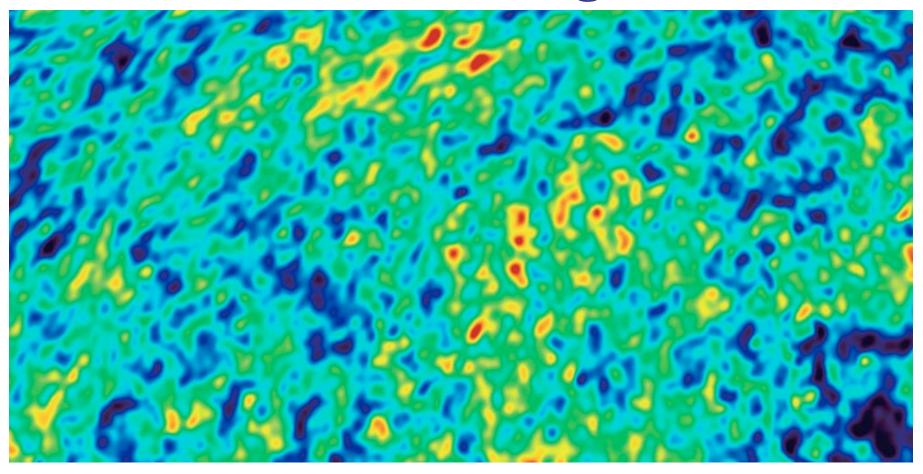
Again,
evolution
overwhelms
the
cosmological
effects ...

# The Modern Angular Diameter Test: CMBR Fluctuations

- Uses the size of the particle horizon at the time of the recombination (the release of the CMBR) as a standard ruler
- This governs the largest wavelength of the sound waves produced in the universe then, due to the infall of baryons into the large-scale density fluctuations
- These sound waves cause small fluctuations in the temperature of the CMB ( $\Delta T/T \sim 10^{-5} 10^{-6}$ ) at the appropriate angular scales ( $\sim$  a degree and less)
- They are measured as the angular power spectra of temperature fluctuations of the CMBR



### A characteristic Fluctuation Scale Exists of ~ 1 degree



This corresponds to the size of the particle horizon at the decoupling, and thus to the longest sound wavelength which can be present

## The results look like this:

Angular Scale 90°  $0.2^{\circ}$  $0.5^{\circ}$ 6000 TT Cross Power Spectrum 5000 - CDM All Data 4000 CBI  $(l+1)C_l/2\pi (\mu K^2)$ ACBAR 3000 2000 1000

100

200

Multipole moment (1)

400

800

1400

WMAP, angular power spectrum, Bennett et al. 2003

Observed position of the first peak is at:

0

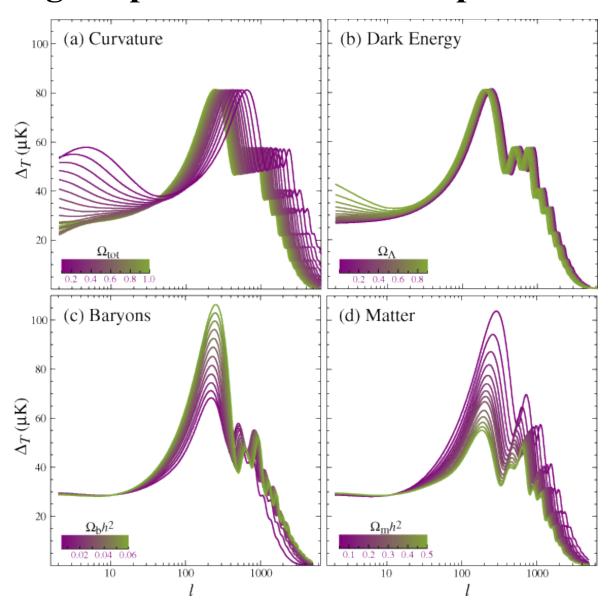
$$l = 220 \quad \Longrightarrow \quad \Omega_{total} = 1.02 \pm 0.02$$

40

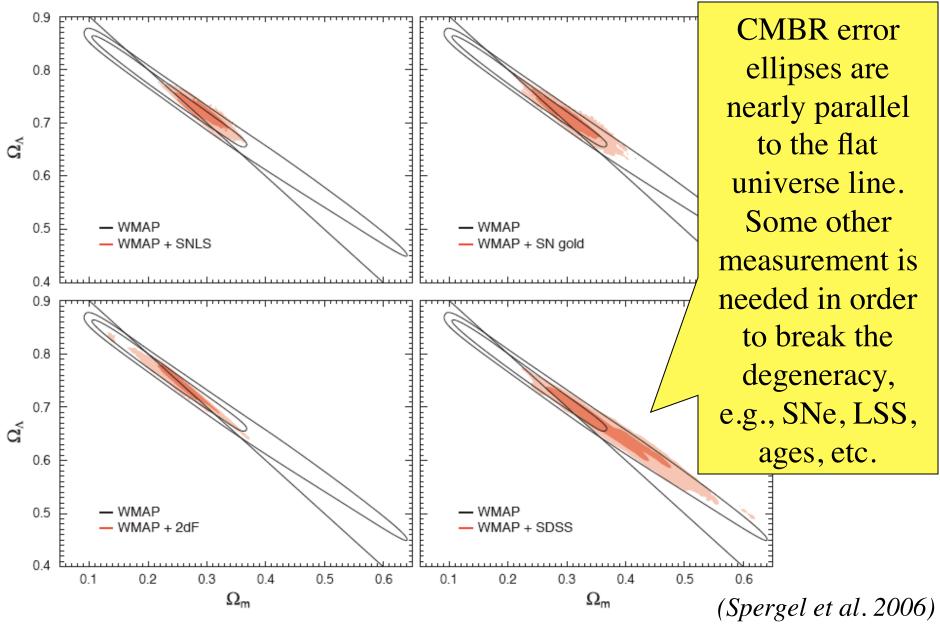
i.e., the Universe is flat (or very close to being flat)

10

## Positions and amplitudes of peaks depend on a variety of cosmological parameters in a complex fashion

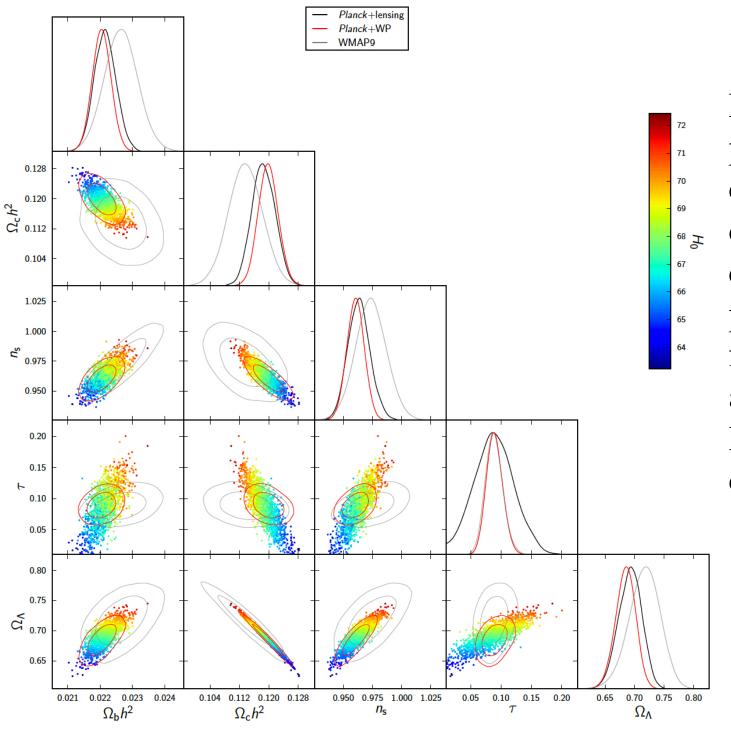


### **CMBR Parameter Degeneracy**



#### **Estimating Cosmological Parameters**

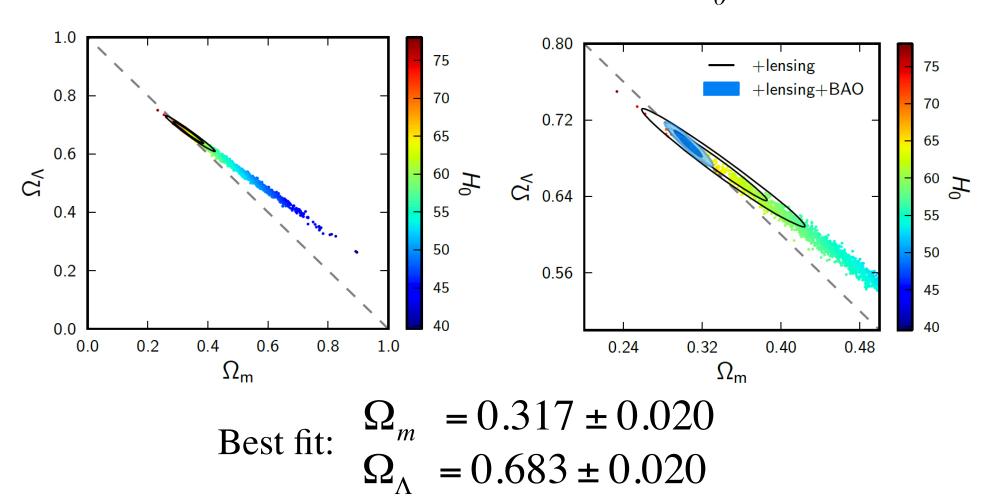
- Many observables depend on complicated combinations of individual cosmological parameters; this is especially true for the analysis of CMB experiments
- Thus, one really gets probability contours or distributions in a multi-dimensional parameter space, which can then be projected on any given parameter axis
- Generally this entails a very laborious and computationally intensive parameter estimation
- It helps if one can declare some of the parameters to be fixed *a priori*, on the basis of our knowledge or prejudices, e.g., "We'll assume that the univese is flat", or "we'll assume the value of  $H_0$  from the HST Key Project", etc.



Examples of probability distributions of the various cosmological parameters, from a joint analysis of Planck and other data

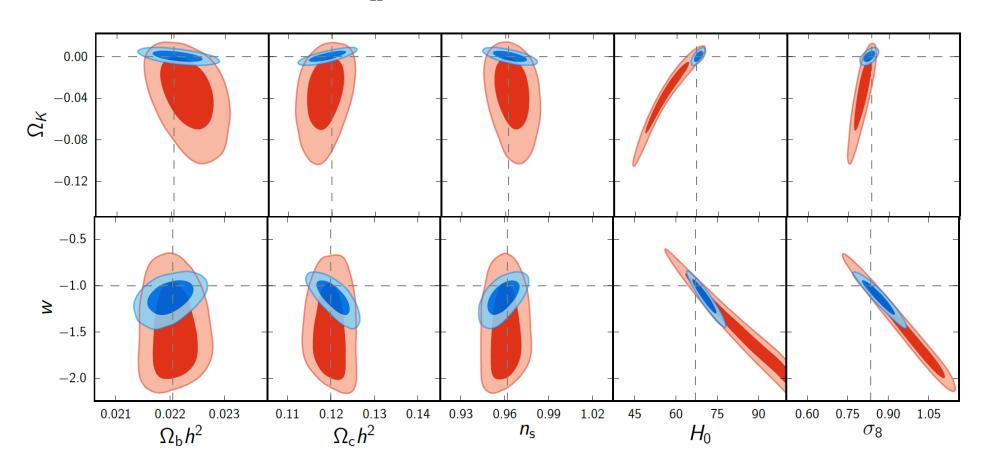
# Some Planck Results (2013)

Matter density and vacuum energy (cosmological constant), for different values of the  $H_0$ 



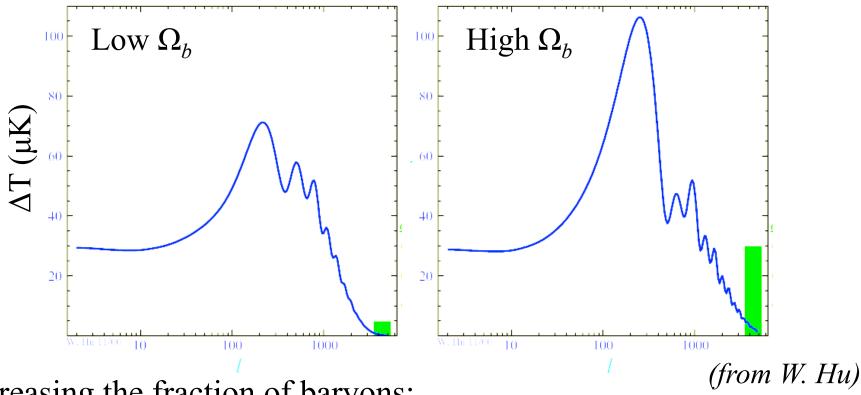
## Some Planck Results (2013)

### Curvature $\Omega_k$ and the EOS parameter w



**Fig. 21.** 68% and 95% confidence regions on one-parameter extensions of the base ΛCDM model for *Planck*+WP (red) and *Planck*+WP+BAO (blue). Horizontal dashed lines correspond to the fixed base model parameter value, and vertical dashed lines show the mean posterior value in the base model for *Planck*+WP.

## **Baryon Content of the Universe**



Increasing the fraction of baryons:

- Increases the amplitude of the Doppler peaks
- Changes the *relative* strength of the peaks odd peaks become stronger relative to the even peaks (compressions/rarefactions)

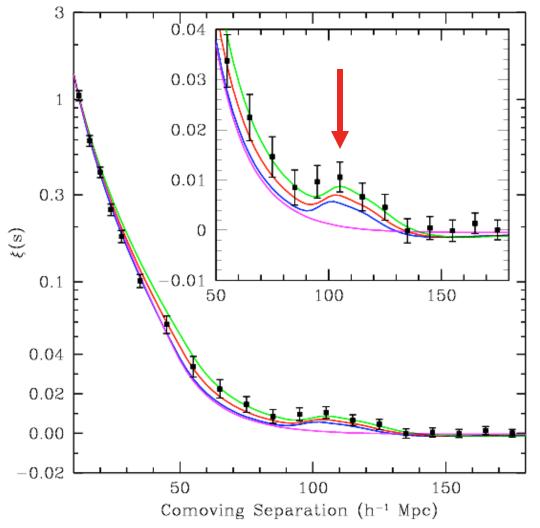
*Planck* results:  $\Omega_b h^2 = 0.022068 \pm 0.00033$ 

### Planck results, 2013

	Planck		Planck+lensing		Planck+WP	
Parameter	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits
$\Omega_{\rm b}h^2$	0.022068	$0.02207 \pm 0.00033$	0.022242	$0.02217 \pm 0.00033$	0.022032	$0.02205 \pm 0.00028$
$\Omega_{\rm c} h^2$	0.12029	$0.1196 \pm 0.0031$	0.11805	$0.1186 \pm 0.0031$	0.12038	$0.1199 \pm 0.0027$
$100\theta_{\mathrm{MC}}$	1.04122	$1.04132 \pm 0.00068$	1.04150	$1.04141 \pm 0.00067$	1.04119	$1.04131 \pm 0.00063$
τ	0.0925	$0.097 \pm 0.038$	0.0949	$0.089 \pm 0.032$	0.0925	$0.089^{+0.012}_{-0.014}$
n <sub>s</sub>	0.9624	$0.9616 \pm 0.0094$	0.9675	$0.9635 \pm 0.0094$	0.9619	$0.9603 \pm 0.0073$
$ln(10^{10}A_s)\;.\;.\;.\;.\;.\;.$	3.098	$3.103 \pm 0.072$	3.098	$3.085 \pm 0.057$	3.0980	$3.089^{+0.024}_{-0.027}$
$\overline{\Omega_{\Lambda} \ldots \ldots \ldots \ldots}$	0.6825	$0.686 \pm 0.020$	0.6964	$0.693 \pm 0.019$	0.6817	$0.685^{+0.018}_{-0.016}$
$\Omega_m \ldots \ldots \ldots$	0.3175	$0.314\pm0.020$	0.3036	$0.307 \pm 0.019$	0.3183	$0.315^{+0.016}_{-0.018}$
$\sigma_8$	0.8344	$0.834 \pm 0.027$	0.8285	$0.823 \pm 0.018$	0.8347	$0.829 \pm 0.012$
$z_{re}$	11.35	$11.4^{+4.0}_{-2.8}$	11.45	$10.8^{+3.1}_{-2.5}$	11.37	$11.1 \pm 1.1$
$H_0$	67.11	$67.4 \pm 1.4$	68.14	$67.9 \pm 1.5$	67.04	$67.3 \pm 1.2$
$10^9 A_s$	2.215	$2.23\pm0.16$	2.215	$2.19^{+0.12}_{-0.14}$	2.215	$2.196^{+0.051}_{-0.060}$
$\Omega_m \hbar^2 \ \dots \dots \dots$	0.14300	$0.1423 \pm 0.0029$	0.14094	$0.1414 \pm 0.0029$	0.14305	$0.1426 \pm 0.0025$
$\Omega_m h^3  .  .  .  .  .  .  .  .  .  $	0.09597	$0.09590 \pm 0.00059$	0.09603	$0.09593 \pm 0.00058$	0.09591	$0.09589 \pm 0.00057$
<i>Y</i> <sub>P</sub>	0.247710	$0.24771 \pm 0.00014$	0.247785	$0.24775 \pm 0.00014$	0.247695	$0.24770 \pm 0.00012$
Age/Gyr	13.819	$13.813 \pm 0.058$	13.784	$13.796 \pm 0.058$	13.8242	$13.817 \pm 0.048$
Ζ*	1090.43	$1090.37 \pm 0.65$	1090.01	$1090.16 \pm 0.65$	1090.48	$1090.43 \pm 0.54$

# **Baryon Acoustic Oscillations (BAO)**

Eisenstein et al. 2005 (using SDSS red galaxies); also seen by the 2dF redshift survey



The 1st Doppler peak seen in the CMBR imprints a preferred scale for clustering of galaxies.

Detection of this feature in galaxy clustering at  $z \sim 0.3$  gives us another instance of a "standard ruler" for an angular diameter test, at redshifts z < 1100

Future redshift surveys can do much better yet

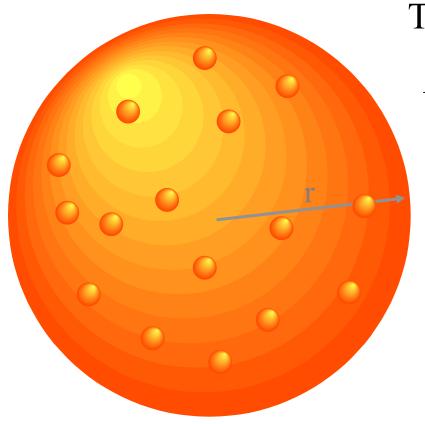
### **The Number Counts**

- Essentially a volume vs. redshift test in disguise; use luminosity distance as a proxy for redshifts
- If one can measure lots of reshifts (expensive!), one could also do a more direct test of source counts per unit comoving volume, as a f(z)
- Usually assume that the comoving number density of sources being counted is non-evolving (aha!)
- In radio astronomy, done as a source counts as a function of limiting flux; in optical-IR astronomy, as galaxy counts as a f(magnitude)
- Nowadays, the evolution effect, flux limits, etc., are included in modeling predicted counts, which are then compared with the observations

### **Euclidean Number Counts**

Assume a class of objects with luminosities L, which down to some limiting flux f are visible out to a distance r.

Then, the observed number N is:



$$\begin{array}{ccc}
N \propto V & & & \\
V \propto r^3 & & \\
& \Rightarrow N \propto r^3
\end{array}$$

Since the flux f follows the inverse square law,

$$r^3 \propto f^{-\frac{3}{2}}$$

Thus we have:

$$N \propto f^{-3/2}$$

### **Euclidean Number Counts**

We can generalize this to multiple populations of sources, e.g., sources with different intrinsic luminosities. They all behave in the same way:

$$N = N_{0,1} f^{-\frac{3}{2}} + N_{0,2} f^{-\frac{3}{2}} + \cdots$$

So again: 
$$N = f^{-3/2} \sum N_{0,i}$$

To get the *differential counts* (e.g., per unit magnitude):

$$\frac{dN}{df} \propto -\frac{3}{2} f^{-\frac{5}{2}}$$

Since

$$d\ln N = -\frac{3}{2}d\ln f$$

we get: 
$$\frac{d \ln N}{d \ln f} = -\frac{3}{2}$$

# Cosmological Number Counts

In relativistic cosmological models, the volume element is generally:

$$dV = \frac{R^3 r^2 dr d\varphi}{\left(1 - kr^2\right)^{/2}}$$

So the count of sources out to some distance  $r_0$  is:

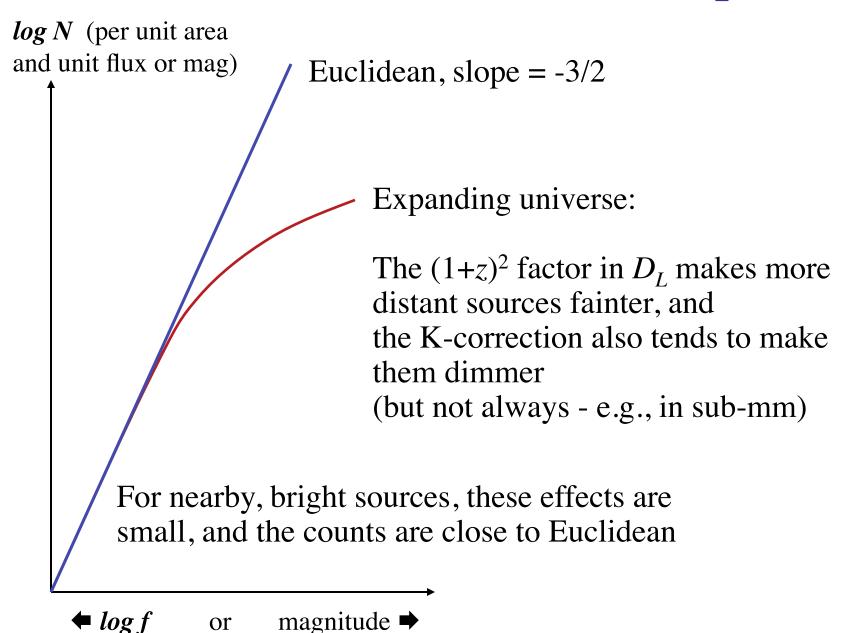
$$N = \int ndV = \int n_0 R_0^3 \frac{r^2 dr d\varphi}{(1 - kr^2)^{1/2}} = 4\pi n_0 R_0^3 \int_0^{r_0} \frac{r^2 dr}{(1 - kr^2)^{1/2}}$$

Since their fluxes are:  $f = L / (4 \pi D_L^2)$ 

 $\rightarrow$  Both N and f depend on cosmology!

As it turns out, all matter-dominated, P = 0 models have  $\frac{d \ln N}{d \ln f} > -\frac{3}{2}$ 

### Source Counts: The Effect of Expansion



## **Source Counts: The Effect of Cosmology**

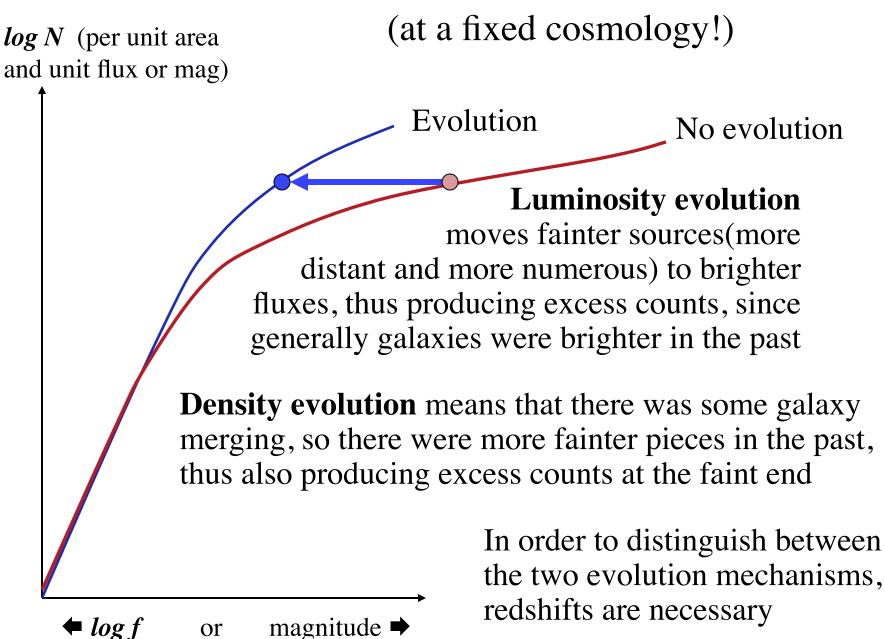
log N (per unit area
and unit flux or mag)
(with no evolution!)

Model with a lower density and/or  $\Lambda > 0$  has more volume and thus more sources to count

Model with a higher density and/or  $\Lambda \leq 0$  has a smaller volume and thus fewer sources to count

For nearby, bright sources, these effects are small, and the counts are close to Euclidean

### **Source Counts: The Effect of Evolution**

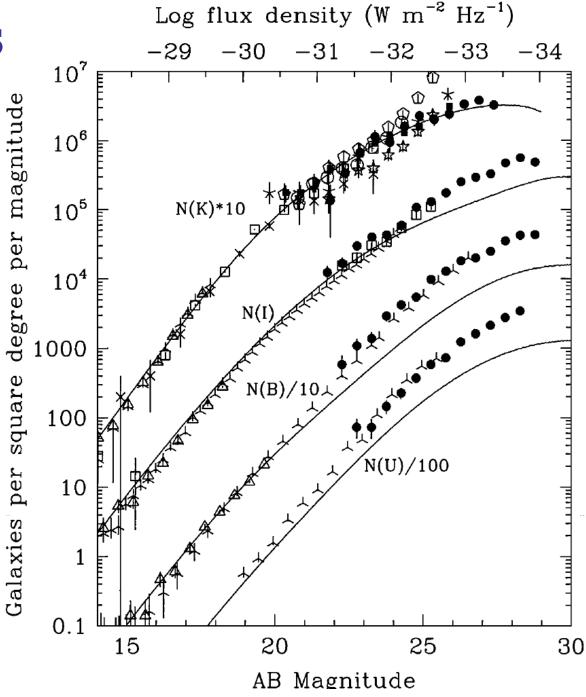


# **Galaxy Counts** in Practice

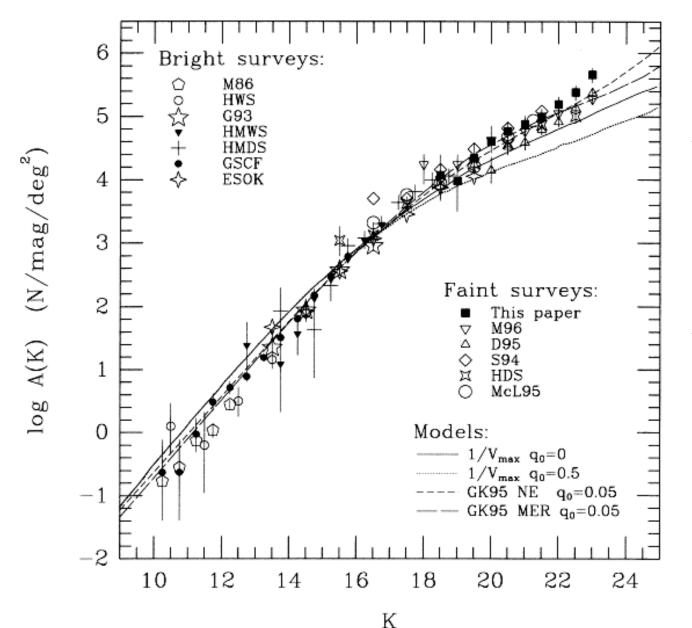
The deepest galaxy counts to date come from HST deep and ultra-deep observations, reaching down to ~ 29<sup>th</sup> mag

All show excess over the no-evolution models, and more in the bluer bands

The extrapolated total count is  $\sim 10^{11}$  galaxies over the entire sky



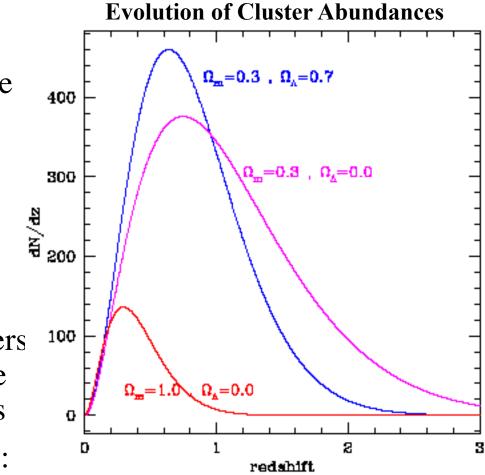
# **Galaxy Counts in Practice**



These effects are less prominent, but still present in the near-IR bands, where the effects of unobscured star formation should be less strong, as the light is dominated by the older, slowly evolving red giants

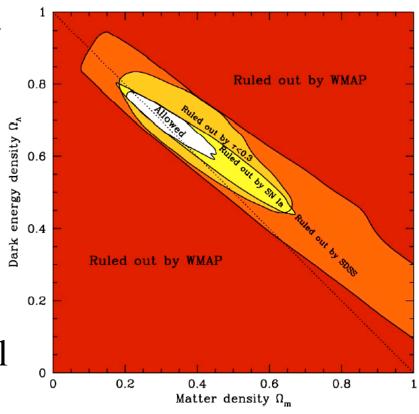
# **Abundance of Rich Galaxy Clusters**

- Given the number density of nearby clusters, we can calculate how many distant clusters we expect to see
- In a high density universe, clusters are just forming now, and we don't expect to find any distant ones
- In a low density universe, clusters began forming long ago, and we expect to find many distant ones
- Evolution of cluster abundances:
  - Structures grow more slowly in a low density universe, so we expect to see less evolution when we probe to large distances
  - Expected number in survey grows because volume probed within a particular spot on the sky increases rapidly with distance



# The Dark Energy

- The **dominant component** of the observed matter/energy density:  $\Omega_{0,DE} \approx 0.7$
- Causes the accelerated expansion of the universe
- May affect the growth of density perturbations
- Effective only at cosmological distances



- Its physical nature is as yet *unknown*; this may be the biggest outstanding problem in physics today
- *Cosmological constant* is just one special case; a more general possibility is called *quintessence*

## Cosmological Constant or Quintessence?

- Cosmological constant: energy density constant in time and spatially uniform
  - Corresponds to the energy density of the physical vacuum
  - A coincidence problem: why is  $\Omega_{\Lambda} \sim \Omega_{\rm m}$  just now?
- Quintessence: time dependent and possibly spatially inhomogeneous; e.g. scalar field rolling down a potential
- Both can be described in the equation of state formalism:

$$P = w \rho$$

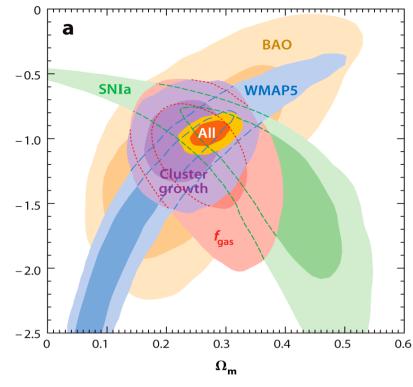
$$\rho \sim R^{-3(w+1)}$$

Cosmological constant: w = const. = -1,  $\rho = const.$ 

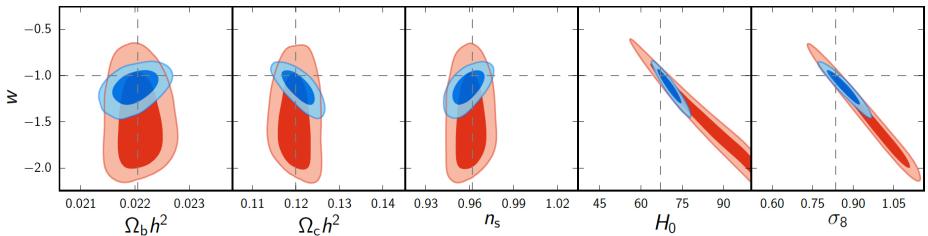
Quintessence: w can have other values and change in time

### Observational Constraints on w

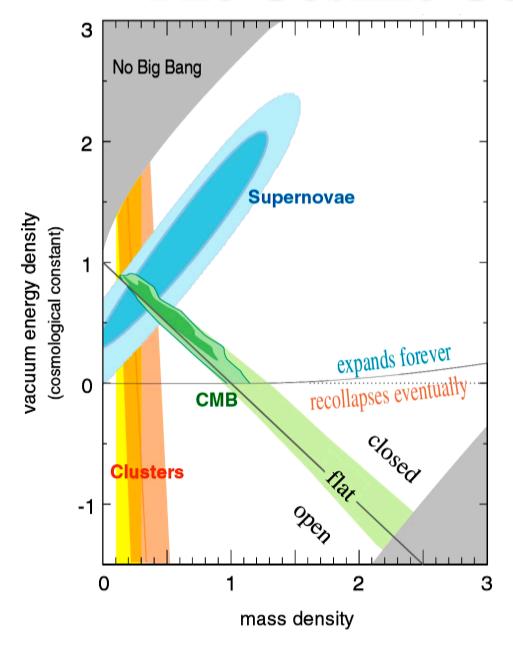
Strongly favor values of  $w \sim -1$ , i.e., cosmological constant. Some models can be excluded, but there is still room for  $\rho_{vac} \neq const.$  models



Planck + WMAP (red) + BAO (blue)



### The Cosmic Concordance



### Supernovae alone

**⇒** Accelerating expansion

 $\Rightarrow \Lambda > 0$ 

#### CMB alone

**⇒** Flat universe

$$\Rightarrow \Lambda > 0$$

Any two of SN, CMB, LSS

⇒ Dark energy ~70%

Also in agreement with the age estimates (globular clusters, nucleocosmochronology, white dwarfs)

# Today's Best Guess Universe

#### Age:

$$t_0 = 13.82 \pm 0.05 \text{ Gyr}$$

#### **Hubble constant:**

$$H_0 = 69 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

#### **Density of ordinary matter:**

$$\Omega_{baryon} = 0.04$$

#### Density of all forms of matter:

$$\Omega_{matter} = 0.31$$

#### **Cosmological constant:**

$$\Omega_{\Lambda} = 0.69$$

Best fit CMB model - consistent with ages of oldest stars

CMB + HST Key Project to measure Cepheid distances

CMB + comparison of nucleosynthesis with Lyman-a forest deuterium measurement

Cluster dark matter estimate CMB power spectrum

Supernova data, CMB evidence for a flat universe plus a low matter density

# The Component Densities

at  $z \sim 0$ , in critical density units, assuming  $h \approx 0.7$ 

Total matter/energy density:  $\Omega_{0,tot} \approx 1.00$  From CMB, and consistent with SNe, LSS

Matter density:  $\Omega_{0,m} \approx 0.31$  From local dynamics and LSS, and consistent with SNe, CMB

Baryon density:  $\Omega_{0,b} \approx 0.045$  From cosmic nucleosynthesis, and independently from CMB

**Luminous baryon density:**  $\Omega_{0,lum} \approx 0.005$  From the census of luminous matter (stars, gas)

Since:  $\Omega_{0,tot} > \Omega_{0,m} > \Omega_{0,b} > \Omega_{0,lum}$ There is baryonic dark matter

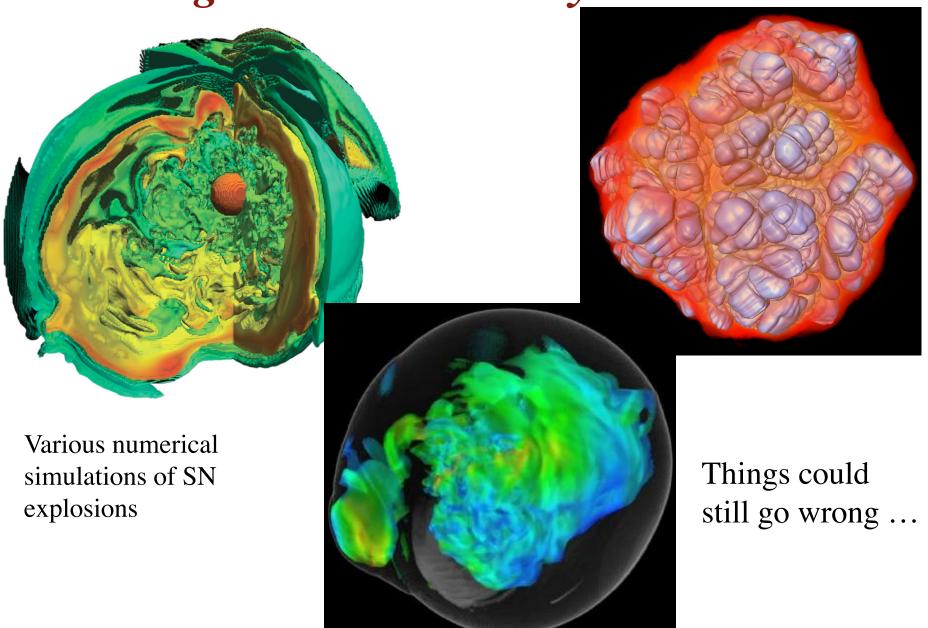
There is dark energy

# **Cosmological Tests Summary**

- Tests of the global geometry and dynamics: correlate redshifts (~ scale factors) with some relative measure of distance (~ look back time); could use:
  - "standard candles" (for luminosity distances; e.g., SNe)
  - "standard rulers" (for angular diameter dist's; e.g., CMBR fluc's)
  - "standard abundances" (for volume-redshift test; e.g., rich clusters)
- Get matter density from local dynamics or LSS
- Combine with constraints from the  $H_0$ , ages
- There are often parameter couplings and degeneracies, especially with the CMB alone
- Multiple approaches provide cross-checks, break degeneracies
- Concordance cosmology is now fairly well established

# **Appendix: Supplementary Slides**

### Warning! SNe are a Messy Phenomenon!

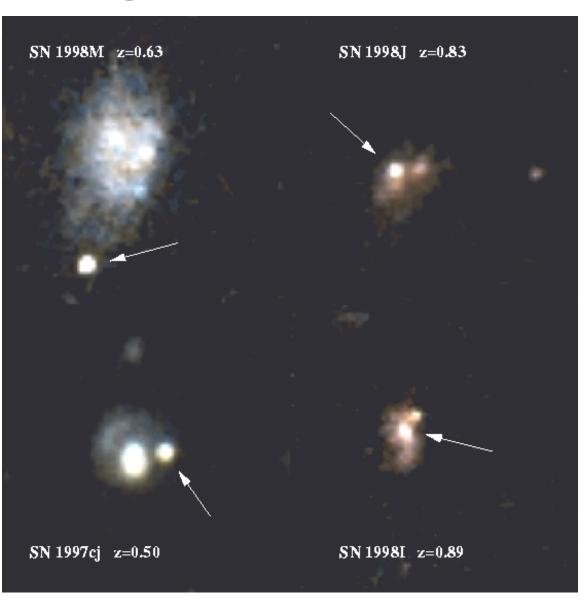


# **Examples of High-Redshift SNe**

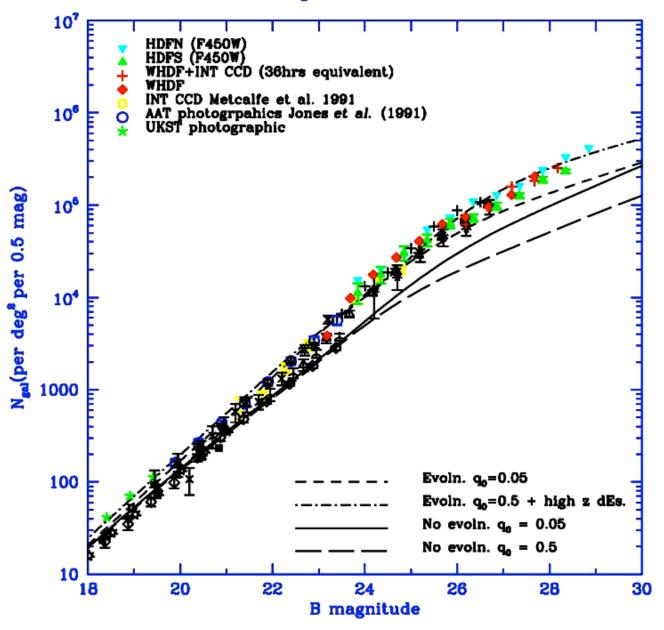
HST observations of SNe in distant galaxies (*Riess et al.*)

Note: you need to ...

- Detect them
- Measure the light curves
- Do the K-corrections
- Get the redshifts



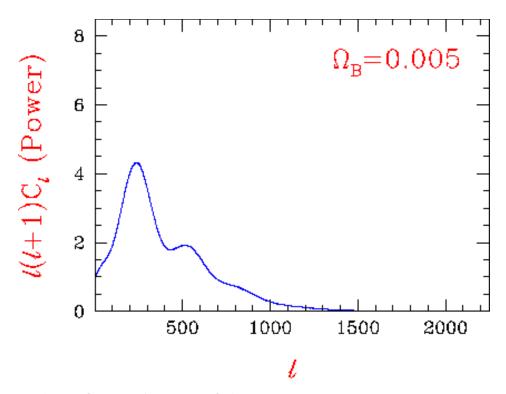
# **Galaxy Counts in Practice**



Observed counts demand some evolution, and favor larger volume (i.e., low  $\Omega_{\rm m}, \Omega_{\Lambda} > 0$ ) cosmological models

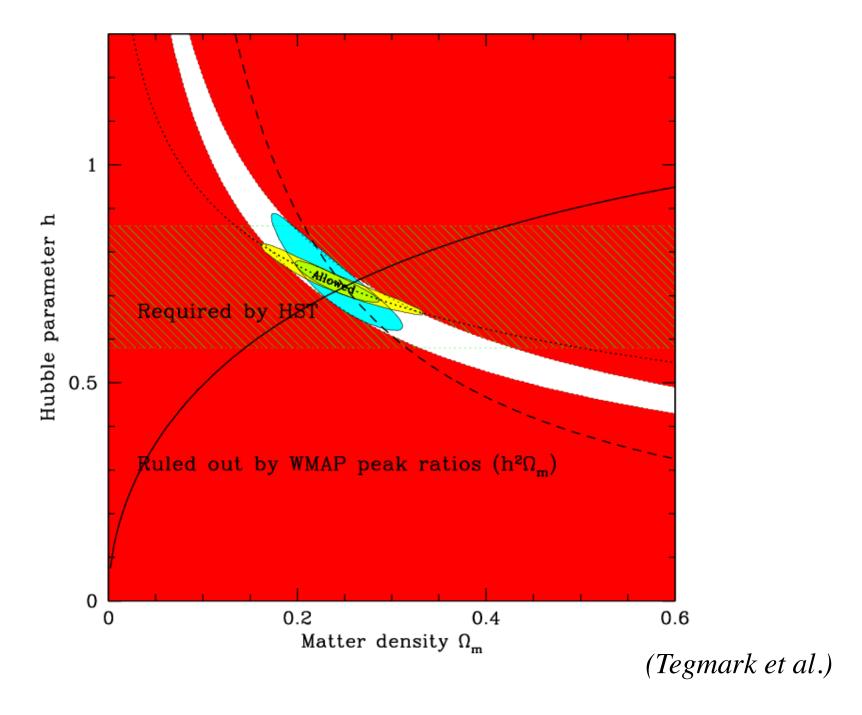
We expect the evolution effects to be stronger in the bluer bands, since they probe UV continua of massive, luminous, short-lived stars

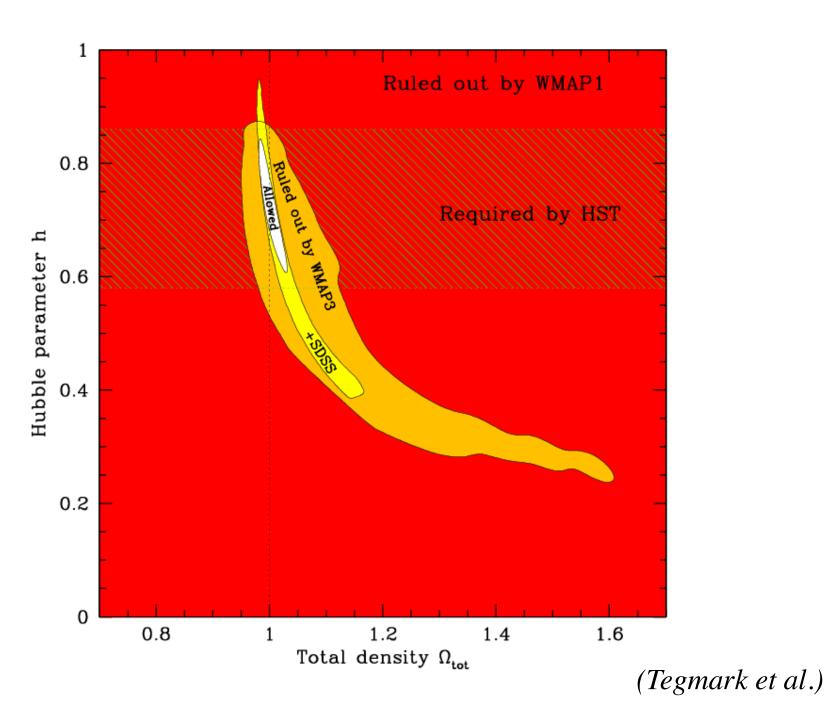
### **Baryon content of the Universe**

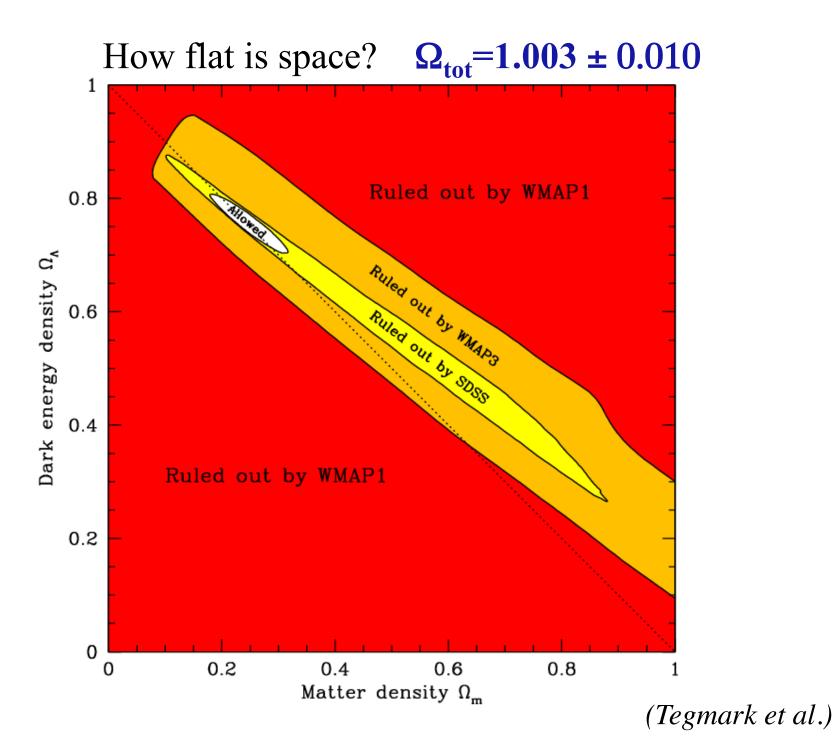


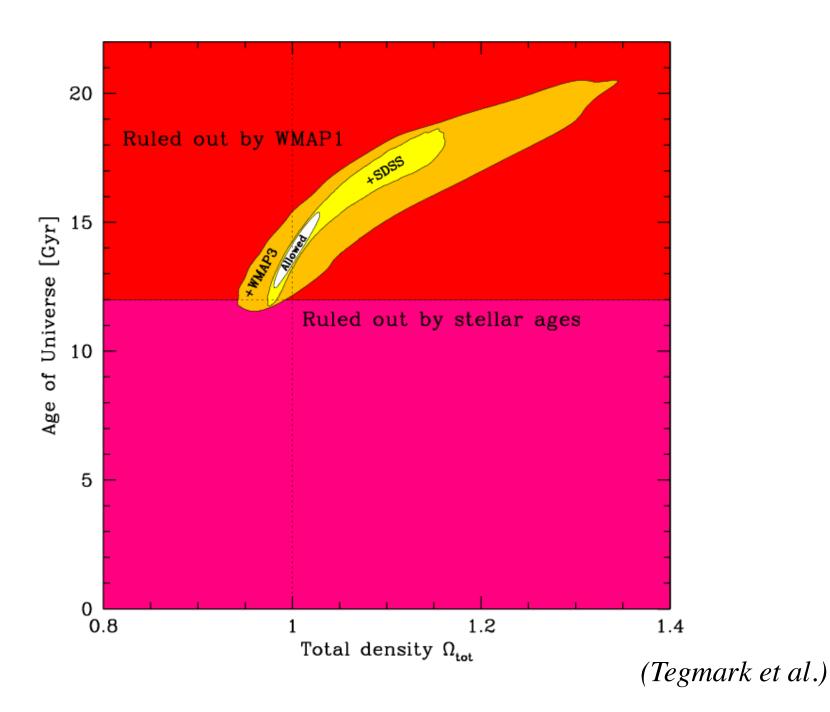
Increasing the fraction of baryons:

- Increases the amplitude of the Doppler peaks
- Changes the *relative* strength of the peaks odd peaks (due to compressions) become stronger relative to the even peaks (due to rarefactions)

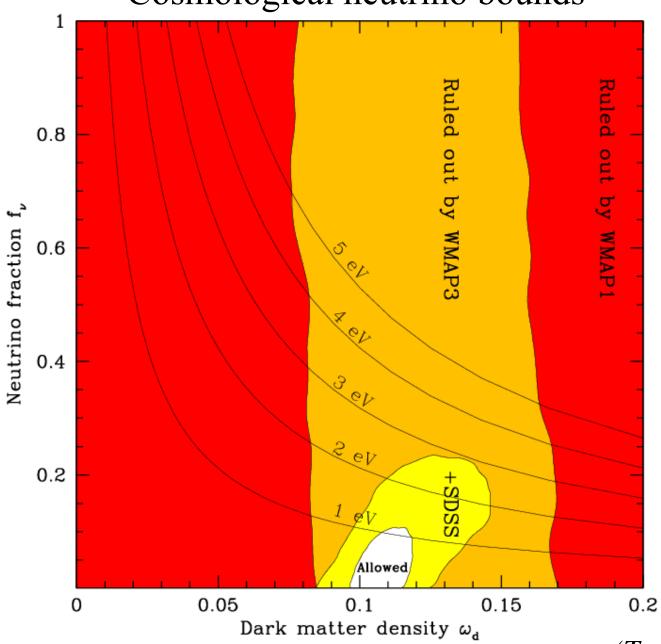








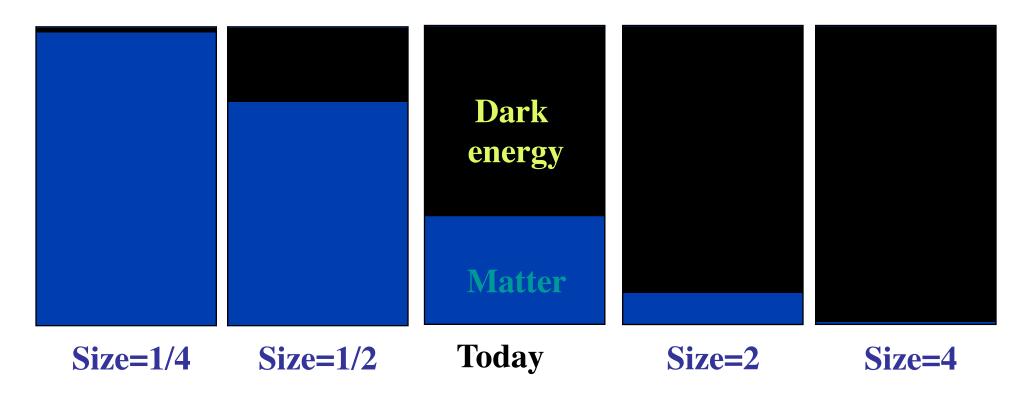
### Cosmological neutrino bounds



(Tegmark et al.)

### The Cosmic Coincidence Problem

If the dark energy is really due to a cosmological constant, its density does not change in time, whereas the matter density does - and they just happen to be comparable today! Seems un-natural ...

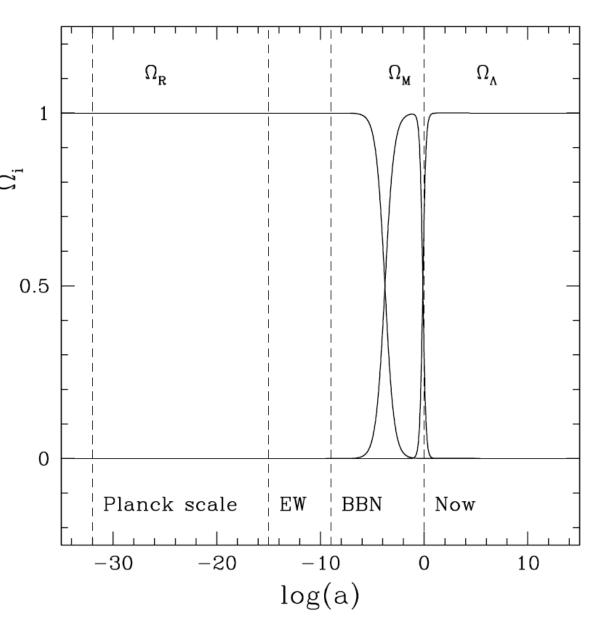


## The Cosmological Coincidence Problem

The time dependence of the density parameter of various mass/energy density components:

We seem to live in a special era, when the vacuum energy density is just starting to dominate the dynamics of the universe ...

... However, this is entirely an artifact of using the log X axis...



### Joining the CMBR, SN, and LSS Results

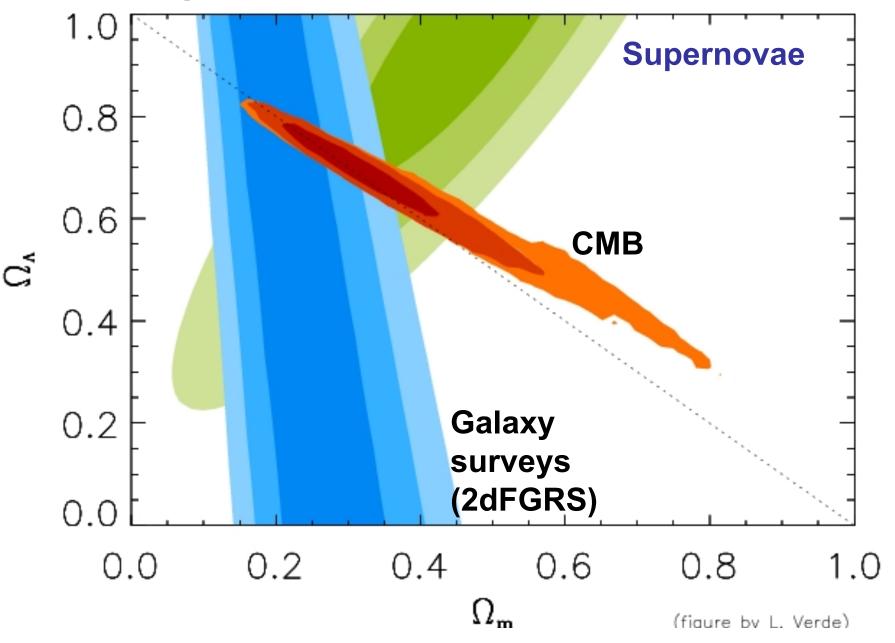


TABLE 2
DERIVED COSMOLOGICAL PARAMETERS

Parameter	Mean (68% Confidence Range)
Amplitude of galaxy fluctuations, $\sigma_8$	$0.9 \pm 0.1$
Characteristic amplitude of velocity fluctuations, $\sigma_8 \Omega_m^{0.6}$	$0.44 \pm 0.10$
Baryon density/critical density, $\Omega_b$	$0.047 \pm 0.006$
Matter density/critical density, $\Omega_m$	$0.29 \pm 0.07$
Age of the universe, $t_0$	$13.4 \pm 0.3  \mathrm{Gyr}$
Redshift of reionization, $z_r$	$17\pm5$
Redshift at decoupling, z <sub>dec</sub>	$1088^{+1}_{-2}$
Age of the universe at decoupling, $t_{\text{dec}}$	$372 \pm 14 \mathrm{kyr}$
Thickness of surface of last scatter, $\Delta z_{\rm dec}$	$194 \pm 2$
Thickness of surface of last scatter, $\Delta t_{\rm dec}$	$115 \pm 5 \mathrm{kyr}$
Redshift at matter/radiation equality, $z_{eq}$	$3454^{+385}_{-392}$
Sound horizon at decoupling, r <sub>s</sub>	$144 \pm 4 \mathrm{Mpc}$
Angular diameter distance to the decoupling surface, $d_A$	$13.7 \pm 0.5  \mathrm{Gpc}$
Acoustic angular scale, ${}^{b}\ell_{A}$	$299\pm 2$
Current density of baryons, n <sub>b</sub>	$(2.7 \pm 0.1) \times 10^{-7}$ cm $^{-3}$
Baryon/photon ratio, $\eta$	$(6.5^{+0.4}_{-0.3}) \times 10^{-10}$

Note.—Fit to the WMAP data only.

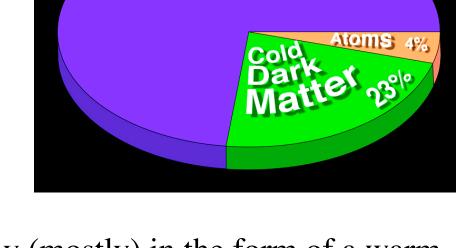
WMAP results, Spergel et al. 2003

<sup>&</sup>lt;sup>a</sup> Assumes ionization fraction,  $x_e = 1$ .

b  $l_A = \pi d_C/r_s$ .

# **Contents of the Universe: Summary**

- $\Omega_0 = 1.00 \pm 0.02$
- $\Omega_m \approx 0.27 \pm 20\%$ 
  - Ω<sub>b</sub> ≈ 0.045 ± 10 %○ Includes Ω<sub>visible</sub> ≈ 0.005
  - $\Omega_{non-b} \approx 0.22$   $\circ \text{ Includes } \Omega_{\nu} < 0.005$
  - $-\Omega_{CMBR} \approx 0.0001$
- $\Omega_{de} \approx 0.73 \pm 10\%$



- The baryonic DM is probably (mostly) in the form of a warm gas ( $\sim 10^5 10^6$  K), associated with galaxies and groups
- The non-baryonic DM may have more than one component, aside from the neutrinos; their nature is as yet unknown, but plausible candidates exist (wimps, axions)
- The physical nature of the DE is currently completely unknown

### This is Not Exactly New ...

Nature Vol. 257 October 9 1975

### An accelerating Universe

James E. Gunn\*

Hale Observatories, California Institute of Technology, Carnegie Institution of Washington, Pasadena, California 91125

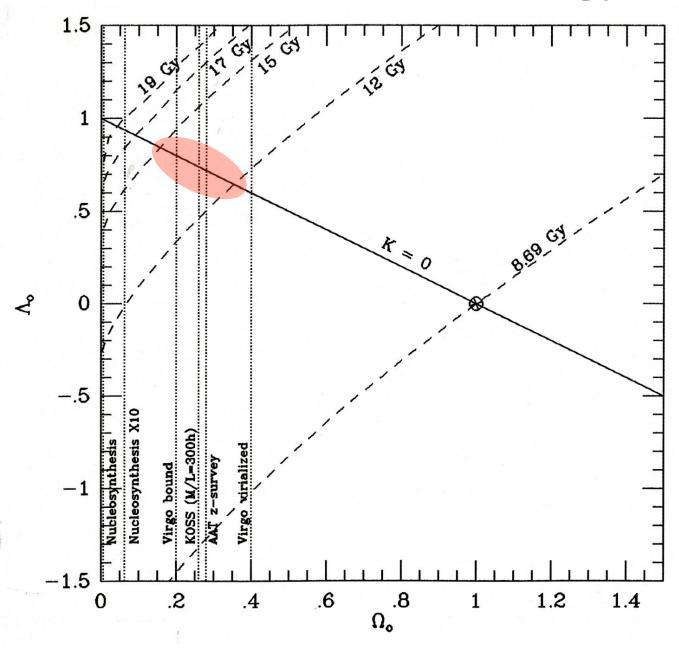
Beatrice M. Tinsley\*†

### B. Tinsley, Nature Vol. 273 18 May 1978 Accelerating Universe revisited

They were driven to this conclusion by the combination of data on the Hubble constant, ages of globular clusters, Hubble diagram, and density measurements ... just like today

For the next 20 years, cosmological constant was invoked mainly as a means to solve the apparent conflict between the ages of globular clusters and chemical elements, and the age of the universe derived from the  $H_0$  and density parameter

# Concordance Cosmology, Circa 1985



 $H_o = 75 \text{ km/s/Mpc}$ 

Globular cluster ages, dynamical measurements of matter density, and  $H_0$ , all consistent with the newly fashionable, flat (k=0) inflationary universe

(Djorgovski 1985, unpublished)