

Exotic Dark Matter Candidates

Ersen Bilgin

Department of Physics
California Institute of Technology

Ph135c, Non-Accelerator Experimental Particle Physics

Outline

- 1 Sterile Neutrinos
 - Motivation
 - The Model
- 2 Little Higgs Models
 - Motivation
 - The Model
- 3 Other Candidates
 - SUSY
 - Light Scalar Dark Matter
 - Wimpzillas
 - Universal Extra Dimensions
 - More Candidates

STERILE NEUTRINOS in ν MSM

- Asaka, Shaposhkinov and Kusenko
- Standard Model incomplete in Neutrino Sector
- Baryon asymmetry in the universe
- Origin of Dark Matter
- Adding three sterile neutrinos explains all this

STERILE NEUTRINOS in ν MSM

- Asaka, Shaposhkinov and Kusenko
- Standard Model incomplete in Neutrino Sector
- Baryon asymmetry in the universe
- Origin of Dark Matter
- Adding three sterile neutrinos explains all this

STERILE NEUTRINOS in ν MSM

- Asaka, Shaposhkinov and Kusenko
- Standard Model incomplete in Neutrino Sector
- Baryon asymmetry in the universe
- Origin of Dark Matter
- Adding three sterile neutrinos explains all this

STERILE NEUTRINOS in ν MSM

- Asaka, Shaposhkinov and Kusenko
- Standard Model incomplete in Neutrino Sector
- Baryon asymmetry in the universe
- Origin of Dark Matter
- Adding three sterile neutrinos explains all this

STERILE NEUTRINOS in ν MSM

- Asaka, Shaposhkinov and Kusenko
- Standard Model incomplete in Neutrino Sector
- Baryon asymmetry in the universe
- Origin of Dark Matter

- Adding three sterile neutrinos explains all this

INDIRECT HINTS

- **Consistent with neutrino oscillations**
- Lightest can account for cosmological dark matter
- Explain observed velocities of pulsars by the emission of light sterile neutrino in supernova explosions
- X-ray photons speed up early star formation (WMAP)
- Heavy sterile neutrinos generate asymmetries between sterile neutrinos and left-handed leptons

INDIRECT HINTS

- Consistent with neutrino oscillations
- Lightest can account for cosmological dark matter
- Explain observed velocities of pulsars by the emission of light sterile neutrino in supernova explosions
- X-ray photons speed up early star formation (WMAP)
- Heavy sterile neutrinos generate asymmetries between sterile neutrinos and left-handed leptons

INDIRECT HINTS

- Consistent with neutrino oscillations
- Lightest can account for cosmological dark matter
- Explain observed velocities of pulsars by the emission of light sterile neutrino in supernova explosions
- X-ray photons speed up early star formation (WMAP)
- Heavy sterile neutrinos generate asymmetries between sterile neutrinos and left-handed leptons

INDIRECT HINTS

- Consistent with neutrino oscillations
- Lightest can account for cosmological dark matter
- Explain observed velocities of pulsars by the emission of light sterile neutrino in supernova explosions
- X-ray photons speed up early star formation (WMAP)
- Heavy sterile neutrinos generate asymmetries between sterile neutrinos and left-handed leptons

INDIRECT HINTS

- Consistent with neutrino oscillations
- Lightest can account for cosmological dark matter
- Explain observed velocities of pulsars by the emission of light sterile neutrino in supernova explosions
- X-ray photons speed up early star formation (WMAP)
- Heavy sterile neutrinos generate asymmetries between sterile neutrinos and left-handed leptons

The Model

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - F_{\alpha I} \Phi \bar{L}_\alpha N_I - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.},$$

$$\begin{pmatrix} 0 & M^D \\ (M^D)^T & M^I \end{pmatrix}, \quad M^D = F \langle \Phi \rangle \quad (1)$$

$$m_\nu = -M^D \frac{1}{M^I} [M^D]^T$$

$$\theta^2 = \frac{1}{M_S^2} \sum_{\alpha=e\mu\tau} |M_{\alpha 1}^D|^2$$

The Model

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - F_{\alpha I} \Phi \bar{L}_\alpha N_I - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.},$$

$$\begin{pmatrix} 0 & M^D \\ (M^D)^T & M^I \end{pmatrix}, \quad M^D = F \langle \Phi \rangle \quad (1)$$

$$m_\nu = -M^D \frac{1}{M^I} [M^D]^T$$

$$\theta^2 = \frac{1}{M_S^2} \sum_{\alpha=e\mu\tau} |M_{\alpha 1}^D|^2$$

The Model

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - F_{\alpha I} \Phi \bar{L}_\alpha N_I - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.},$$

$$\begin{pmatrix} 0 & M^D \\ (M^D)^T & M^I \end{pmatrix}, \quad M^D = F \langle \Phi \rangle \quad (1)$$

$$m_\nu = -M^D \frac{1}{M^I} [M^D]^T$$

$$\theta^2 = \frac{1}{M_S^2} \sum_{\alpha=e\mu\tau} |M_{\alpha 1}^D|^2$$

The Model

$$\mathcal{L}_{\nu\text{MSM}} = \mathcal{L}_{\text{MSM}} + \bar{N}_I i \partial_\mu \gamma^\mu N_I - F_{\alpha I} \Phi \bar{L}_\alpha N_I - \frac{M_I}{2} \bar{N}_I^c N_I + \text{h.c.},$$

$$\begin{pmatrix} 0 & M^D \\ (M^D)^T & M^I \end{pmatrix}, \quad M^D = F \langle \Phi \rangle \quad (1)$$

$$m_\nu = -M^D \frac{1}{M^I} [M^D]^T$$

$$\theta^2 = \frac{1}{M_S^2} \sum_{\alpha=e\mu\tau} |M_{\alpha 1}^D|^2$$

Constraints

- $\theta < \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ to avoid overclosing the universe
- $\theta > 5 \times 10^{-4} (1 \text{ keV}/M_s)^{1/2}$ for the sterile neutrinos to be in thermal equilibrium
- Virgo cluster: $\theta < 1.6 \times 10^{-3} (1 \text{ keV}/M_s)^2$, for $1 \text{ keV} < M_s < 10 \text{ keV}$
- X-ray background : $\theta < 5.8 \times 10^{-3} (1 \text{ keV}/M_s)^{5/2}$, for $1 \text{ keV} < M_s < 100 \text{ keV}$

Constraints

- $\theta < \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ to avoid overclosing the universe
- $\theta > 5 \times 10^{-4} (1 \text{ keV}/M_s)^{1/2}$ for the sterile neutrinos to be in thermal equilibrium
- Virgo cluster: $\theta < 1.6 \times 10^{-3} (1 \text{ keV}/M_s)^2$, for $1 \text{ keV} < M_s < 10 \text{ keV}$
- X-ray background : $\theta < 5.8 \times 10^{-3} (1 \text{ keV}/M_s)^{5/2}$, for $1 \text{ keV} < M_s < 100 \text{ keV}$

Constraints

- $\theta < \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ to avoid overclosing the universe
- $\theta > 5 \times 10^{-4} (1 \text{ keV}/M_s)^{1/2}$ for the sterile neutrinos to be in thermal equilibrium
- Virgo cluster: $\theta < 1.6 \times 10^{-3} (1 \text{ keV}/M_s)^2$, for $1 \text{ keV} < M_s < 10 \text{ keV}$
- X-ray background : $\theta < 5.8 \times 10^{-3} (1 \text{ keV}/M_s)^{5/2}$, for $1 \text{ keV} < M_s < 100 \text{ keV}$

Constraints

- $\theta < \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ to avoid overclosing the universe
- $\theta > 5 \times 10^{-4} (1 \text{ keV}/M_s)^{1/2}$ for the sterile neutrinos to be in thermal equilibrium
- Virgo cluster: $\theta < 1.6 \times 10^{-3} (1 \text{ keV}/M_s)^2$, for $1 \text{ keV} < M_s < 10 \text{ keV}$
- X-ray background : $\theta < 5.8 \times 10^{-3} (1 \text{ keV}/M_s)^{5/2}$, for $1 \text{ keV} < M_s < 100 \text{ keV}$

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio I

- No sterile neutrinos at $T > 1$ GeV
 - Weak couplings between singlet fermions and fields beyond SM.
 - $\theta = \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ for the right amount of dark matter
 - This and Virgo cluster observations give $M_s < 8$ keV
 - CMB, Lyman- α , Sloan Digital Sky Survey: $M_s > 2$ keV (due to free streaming length of the sterile neutrino)

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio I

- No sterile neutrinos at $T > 1$ GeV
 - Weak couplings between singlet fermions and fields beyond SM.
 - $\theta = \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ for the right amount of dark matter
 - This and Virgo cluster observations give $M_s < 8$ keV
 - CMB, Lyman- α , Sloan Digital Sky Survey: $M_s > 2$ keV (due to free streaming length of the sterile neutrino)

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio I

- No sterile neutrinos at $T > 1$ GeV
 - Weak couplings between singlet fermions and fields beyond SM.
 - $\theta = \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ for the right amount of dark matter
 - This and Virgo cluster observations give $M_s < 8$ keV
 - CMB, Lyman- α , Sloan Digital Sky Survey: $M_s > 2$ keV (due to free streaming length of the sterile neutrino)

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio I

- No sterile neutrinos at $T > 1$ GeV
 - Weak couplings between singlet fermions and fields beyond SM.
 - $\theta = \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ for the right amount of dark matter
 - This and Virgo cluster observations give $M_s < 8$ keV
 - CMB, Lyman- α , Sloan Digital Sky Survey: $M_s > 2$ keV (due to free streaming length of the sterile neutrino)

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio I

- No sterile neutrinos at $T > 1$ GeV
 - Weak couplings between singlet fermions and fields beyond SM.
 - $\theta = \theta_{max}(M_s) = 1.3 \times 10^{-4} \left(\frac{1 \text{ keV}}{M_s} \right)^{0.8}$ for the right amount of dark matter
 - This and Virgo cluster observations give $M_s < 8$ keV
 - CMB, Lyman- α , Sloan Digital Sky Survey: $M_s > 2$ keV (due to free streaming length of the sterile neutrino)

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 - $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

SCENERIOS FOR ONE STERILE NEUTRINO

Scenerio II

- Sterile Neutrinos at equilibrium at high temperature and abundance at EW scale is the same as active neutrinos
 - Force $M_s \sim 100$ eV
 - Does not agree with Tremaine-Gunn bound:
 $M_s > M_{TG} \simeq 0.3$ keV

Scenerio III

- Sterile Neutrinos at equilibrium at high temperature but not through active sterile neutrino oscillations
 - Not studied much
 - $\theta < 2 \times 10^{-4}$ by X-ray observations not directly detectable in laboratory
 - $M_s > M_{TG} \simeq 0.3$ keV

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Introduce two heavy sterile neutrinos
- Need $M \sim \mathcal{O}(1 - 10)$ GeV to satisfy matter-antimatter asymmetry.
- also need heavy sterile neutrinos to be degenerate in mass
 - amplify CP-violating effects in sterile neutrino oscillations
- Modify bound on M_s : $M_s > 0.55$ keV
- $\theta \sim 9.2 \times 10^{-3}$

- Can be falsified by neutrino physics data
- If Mini-Boone confirmed LSND results, need to add one more neutrino
- Preliminary results say it doesn't.

- Can be falsified by neutrino physics data
- If Mini-Boone confirmed LSND results, need to add one more neutrino
- Preliminary results say it doesn't.

- Can be falsified by neutrino physics data
- If Mini-Boone confirmed LSND results, need to add one more neutrino
- Preliminary results say it doesn't.

Questions on ν MSM

- How does sterile neutrino decay into X-ray photon?
 - $N_1 \rightarrow \gamma\nu, \gamma\bar{\nu}$ (similar to ultra-violet radiation from active neutrino decays)
- What happens to sterile neutrinos below equilibrium temp
 - It drops and currently it should be $\sim \mathcal{O}(10)$ less than active neutrinos
- Explain Tremaine-Gunn Bound – Why doesn't it apply to neutralinos?
 - Liouville's Theorem: For non-interacting particles density of fluid element in phase space does not change.
 - Maximum coarse grained phase space density can only decrease.
 - Apply this to isothermal gas spheres (neutrinos)

Questions on ν MSM

- How does sterile neutrino decay into X-ray photon?
 - $N_1 \rightarrow \gamma\nu, \gamma\bar{\nu}$ (similar to ultra-violet radiation from active neutrino decays)
- What happens to sterile neutrinos below equilibrium temp
 - It drops and currently it should be $\sim \mathcal{O}(10)$ less than active neutrinos
- Explain Tremaine-Gunn Bound – Why doesn't it apply to neutralinos?
 - Liouville's Theorem: For non-interacting particles density of fluid element in phase space does not change.
 - Maximum coarse grained phase space density can only decrease.
 - Apply this to isothermal gas spheres (neutrinos)

Questions on ν MSM

- How does sterile neutrino decay into X-ray photon?
 - $N_1 \rightarrow \gamma\nu, \gamma\bar{\nu}$ (similar to ultra-violet radiation from active neutrino decays)
- What happens to sterile neutrinos below equilibrium temp
 - It drops and currently it should be $\sim \mathcal{O}(10)$ less than active neutrinos
- Explain Tremaine-Gunn Bound – Why doesn't it apply to neutralinos?
 - Liouville's Theorem: For non-interacting particles density of fluid element in phase space does not change.
 - Maximum coarse grained phase space density can only decrease.
 - Apply this to isothermal gas spheres (neutrinos)

LITTLE HIGGS MODELS

- **Alternative to SUSY**
- Higgs mass quadratically divergent.
- To stabilize its mass need new physics at ~ 1 TeV
- Precision electroweak measurements give no evidence of new physics up to $\gtrsim 5 - 7$ GeV
- Not Natural.

LITTLE HIGGS MODELS

- Alternative to SUSY
- Higgs mass quadratically divergent.
- To stabilize its mass need new physics at ~ 1 TeV
- Precision electroweak measurements give no evidence of new physics up to $\gtrsim 5 - 7$ GeV
- Not Natural.

LITTLE HIGGS MODELS

- Alternative to SUSY
- Higgs mass quadratically divergent.
- To stabilize its mass need new physics at ~ 1 TeV
- Precision electroweak measurements give no evidence of new physics up to $\gtrsim 5 - 7$ GeV
- Not Natural.

LITTLE HIGGS MODELS

- Alternative to SUSY
- Higgs mass quadratically divergent.
- To stabilize its mass need new physics at ~ 1 TeV
- Precision electroweak measurements give no evidence of new physics up to $\gtrsim 5 - 7$ GeV
- Not Natural.

LITTLE HIGGS MODELS

- Alternative to SUSY
- Higgs mass quadratically divergent.
- To stabilize its mass need new physics at ~ 1 TeV
- Precision electroweak measurements give no evidence of new physics up to $\gtrsim 5 - 7$ GeV
- Not Natural.

CHENG - LOW MODEL

- 5 – 7 GeV bound assumes new fields couple at tree level
- Only need quantum loop diagrams to cancel quadratic divergences
- No bound if interaction vertices involve Higgs and two or more new TeV particles
- Introduce new symmetry “ *T-Parity* ” acting only on new particles
- Avoids Higgs interacting with one new TeV particle and relaxes the constraints.

CHENG - LOW MODEL

- 5 – 7 GeV bound assumes new fields couple at tree level
- Only need quantum loop diagrams to cancel quadratic divergences
- No bound if interaction vertices involve Higgs and two or more new TeV particles
- Introduce new symmetry “ *T-Parity* ” acting only on new particles
- Avoids Higgs interacting with one new TeV particle and relaxes the constraints.

CHENG - LOW MODEL

- 5 – 7 GeV bound assumes new fields couple at tree level
- Only need quantum loop diagrams to cancel quadratic divergences
- No bound if interaction vertices involve Higgs and two or more new TeV particles
- Introduce new symmetry “ *T-Parity* ” acting only on new particles
- Avoids Higgs interacting with one new TeV particle and relaxes the constraints.

CHENG - LOW MODEL

- 5 – 7 GeV bound assumes new fields couple at tree level
- Only need quantum loop diagrams to cancel quadratic divergences
- No bound if interaction vertices involve Higgs and two or more new TeV particles
- Introduce new symmetry “ *T-Parity* ” acting only on new particles
- Avoids Higgs interacting with one new TeV particle and relaxes the constraints.

CHENG - LOW MODEL

- 5 – 7 GeV bound assumes new fields couple at tree level
- Only need quantum loop diagrams to cancel quadratic divergences
- No bound if interaction vertices involve Higgs and two or more new TeV particles
- Introduce new symmetry “ *T-Parity* ” acting only on new particles
- Avoids Higgs interacting with one new TeV particle and relaxes the constraints.

T-PARITY'S PHENOMENOLOGICAL CONSEQUENCES

- Lightest particle that transforms under this symmetry (LTP) is stable
- If charged leaves tracks, if neutral results in missing energy in colliders
- Similar to R-parity conserving SUSY and KK-parity conserving UEDs.
- Spin is different from LSP, and easier to detect than KK excitations.

T-PARITY'S PHENOMENOLOGICAL CONSEQUENCES

- Lightest particle that transforms under this symmetry (LTP) is stable
- If charged leaves tracks, if neutral results in missing energy in colliders
- Similar to R-parity conserving SUSY and KK-parity conserving UEDs.
- Spin is different from LSP, and easier to detect than KK excitations.

T-PARITY'S PHENOMENOLOGICAL CONSEQUENCES

- Lightest particle that transforms under this symmetry (LTP) is stable
- If charged leaves tracks, if neutral results in missing energy in colliders
- Similar to R-parity conserving SUSY and KK-parity conserving UEDs.
- Spin is different from LSP, and easier to detect than KK excitations.

T-PARITY'S PHENOMENOLOGICAL CONSEQUENCES

- Lightest particle that transforms under this symmetry (LTP) is stable
- If charged leaves tracks, if neutral results in missing energy in colliders
- Similar to R-parity conserving SUSY and KK-parity conserving UEDs.
- Spin is different from LSP, and easier to detect than KK excitations.

Lightest T-odd Particle (LTP)

- neutral B' gauge boson with mass 600 GeV - 1.2 TeV gives the right relic density for dark matter.
 - Annihilation of B' into electron-positron pairs not suppressed.
 - Can be detected at the anti-matter detector on the International Space Station as a peak in the positron energy distribution.
- $SU(2)_W$ singlets and neutral components of some scalar fields are also DM candidates.

Lightest T-odd Particle (LTP)

- neutral B' gauge boson with mass 600 GeV - 1.2 TeV gives the right relic density for dark matter.
 - Annihilation of B' into electron-positron pairs not suppressed.
 - Can be detected at the anti-matter detector on the International Space Station as a peak in the positron energy distribution.
- $SU(2)_W$ singlets and neutral components of some scalar fields are also DM candidates.

Lightest T-odd Particle (LTP)

- neutral B' gauge boson with mass 600 GeV - 1.2 TeV gives the right relic density for dark matter.
 - Annihilation of B' into electron-positron pairs not suppressed.
 - Can be detected at the anti-matter detector on the International Space Station as a peak in the positron energy distribution.
- $SU(2)_W$ singlets and neutral components of some scalar fields are also DM candidates.

Lightest T-odd Particle (LTP)

- neutral B' gauge boson with mass 600 GeV - 1.2 TeV gives the right relic density for dark matter.
 - Annihilation of B' into electron-positron pairs not suppressed.
 - Can be detected at the anti-matter detector on the International Space Station as a peak in the positron energy distribution.
- $SU(2)_W$ singlets and neutral components of some scalar fields are also DM candidates.

Supersymmetric Candidates

- **Sneutrinos: superpartners of SM neutrinos.**
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- **Sneutrinos: superpartners of SM neutrinos.**
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- Sneutrinos: superpartners of SM neutrinos.
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- Sneutrinos: superpartners of SM neutrinos.
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- Sneutrinos: superpartners of SM neutrinos.
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- Sneutrinos: superpartners of SM neutrinos.
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Supersymmetric Candidates

- Sneutrinos: superpartners of SM neutrinos.
 - mass $\sim 550 - 2300$ GeV
 - scattering cross section larger than the upper limits from experiments.
- Gravitinos: Superpartners of Gravitons
 - Lightest SUSY particle and stable in some SUSY scenarios
 - Hard to detect since they only interact through gravitation
- Axinos: Superpartners of Axions
 - Phenomenology similar to gravitinos.
 - Can be cold, warm or hot DM candidate.

Light Scalar Dark Matter

- For fermionic dark matter with standard Fermi interactions, mass of WIMPs \lesssim GeV (Lee and Weinberg)
- Other types of particles (scalar dark matter): mass 1-100 MeV is possible
- 511 keV gamma-ray line from the INTEGRAL satellite from the galactic bulge could be scalar dark matter annihilating into positrons which annihilate to give out the gamma ray line.
- Recently axinos or sterile neutrinos suggested to cause the 511 keV emission.

Light Scalar Dark Matter

- For fermionic dark matter with standard Fermi interactions, mass of WIMPs \lesssim GeV (Lee and Weinberg)
- Other types of particles (scalar dark matter): mass 1-100 MeV is possible
- 511 keV gamma-ray line from the INTEGRAL satellite from the galactic bulge could be scalar dark matter annihilating into positrons which annihilate to give out the gamma ray line.
- Recently axinos or sterile neutrinos suggested to cause the 511 keV emission.

Light Scalar Dark Matter

- For fermionic dark matter with standard Fermi interactions, mass of WIMPs \lesssim GeV (Lee and Weinberg)
- Other types of particles (scalar dark matter): mass 1-100 MeV is possible
- 511 keV gamma-ray line from the INTEGRAL satellite from the galactic bulge could be scalar dark matter annihilating into positrons which annihilate to give out the gamma ray line.
- Recently axinos or sterile neutrinos suggested to cause the 511 keV emission.

Light Scalar Dark Matter

- For fermionic dark matter with standard Fermi interactions, mass of WIMPs \lesssim GeV (Lee and Weinberg)
- Other types of particles (scalar dark matter): mass 1-100 MeV is possible
- 511 keV gamma-ray line from the INTEGRAL satellite from the galactic bulge could be scalar dark matter annihilating into positrons which annihilate to give out the gamma ray line.
- Recently axinos or sterile neutrinos suggested to cause the 511 keV emission.

WIMPZILLAS

- **Unitarity bound** gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

WIMPZILLAS

- *Unitarity bound* gives maximum annihilation cross-section as a function of mass
 - Use WMAP constraint on $\Omega_{DM}h^2$ to get $m_{DM} \lesssim 34$ TeV
- Wimpzillas were NOT in thermal equilibrium during freeze-out
- They have mass $> 10^{10}$ GeV
- created from *gravitational production* at the end of inflation
- Motivation: Cosmic rays above GZK cutoff (5×10^{19} eV)
 - these rays interact at resonance with CMB photons so universe should be opaque to them

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Universal Extra Dimensions

- Conservation of Momentum in higher dimensional space
- Conservation of KK number in compactified space
- Kaluza-Klein Particle
 - studied since 1984
 - Lightest Kaluza-Klein Particle (LKP) has mass ~ 400 to 1200 GeV
 - See the Extra Dimensions talk on May 27.

Even more candidates

- Q-Balls, mirror particles, CHARGed Massive Particles (CHAMPs), self interacting dark matter, D-matter, cryptons, superweakly interacting dark matter, brane world dark matter, heavy fourth generation neutrinos, etc.

Summary

- Don't know much about dark matter
- Different assumptions lead to different dark matter candidates
- Need more experimental data to rule out or verify specific candidates

Summary

- Don't know much about dark matter
- Different assumptions lead to different dark matter candidates
- Need more experimental data to rule out or verify specific candidates

Summary

- Don't know much about dark matter
- Different assumptions lead to different dark matter candidates
- Need more experimental data to rule out or verify specific candidates

References I



G. Bertone, D. Hopper, J. Silk

Particle Dark Matter: Evidence Candidates and Constraints

arXiv:hep-ph/0404175v2



T. Asaka, M. Shaposhkinov, A. Kusenko

Opening a new window for warm dark matter

arXiv:hep-ph/0602150v2



H. Cheng, I. Low

TeV Symmetry and the Little Hierarchy Problem

arXiv:hep-ph/0308199v2